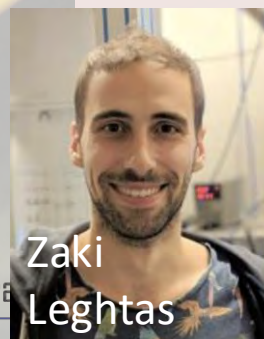
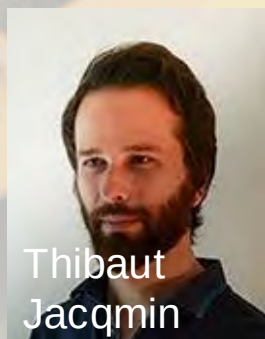


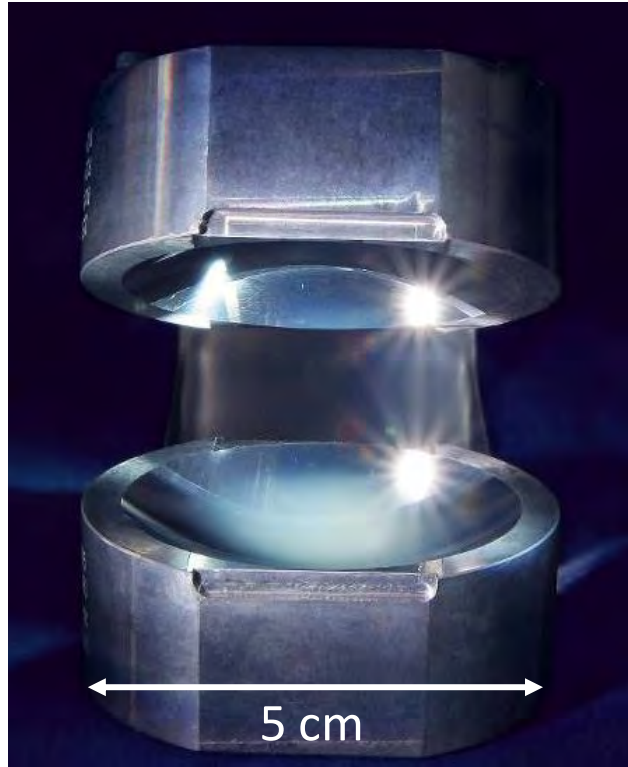
# ***Probing the quantum motion of a MHz mechanical resonator with a resonant rf-fluxonium***

S. Deléglise, K. Gerashchenko, R. Rousseau, L. Balembois, H. Patange, P. Manset, A. Riva, H. Le Sueur, P.-F. Cohadon, T. Briant, A. Heidmann, T. Jacqmin

Collab: Z. Leghtas (LPENS), E. Flurin (CEA), A. Sarlette (Inria)



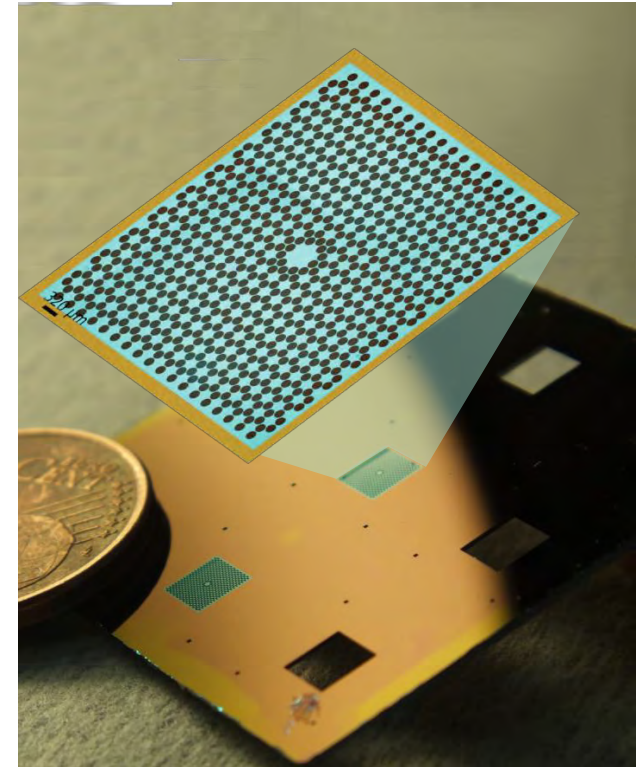
from Cavity Quantum Electro Dynamics ...



## Photon box

- $\omega_{cav}/2\pi = 50 \text{ GHz}$
- $T_1 = 130 \text{ ms}$
- $Q = 4 \cdot 10^{10}$
- $n_{th} = 0.05 \text{ @ } 800 \text{ mK}$

...to circuit Quantum Acousto Dynamics

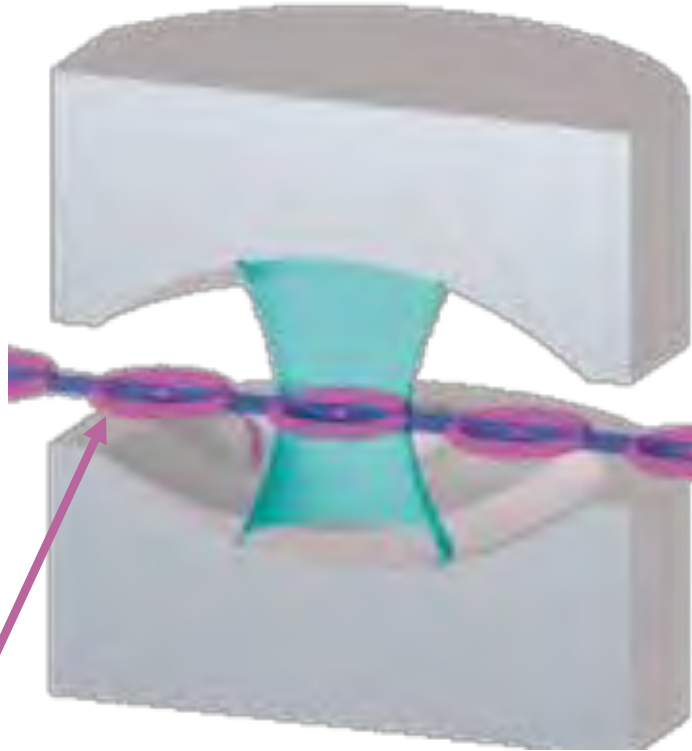


## Phonon box

- $\omega_m/2\pi = 1 \text{ MHz}$
- $T_1 = 2 \text{ minutes}$
- $Q = 10^9$
- $n_{th} = 50 \text{ @ } 10 \text{ mK}$

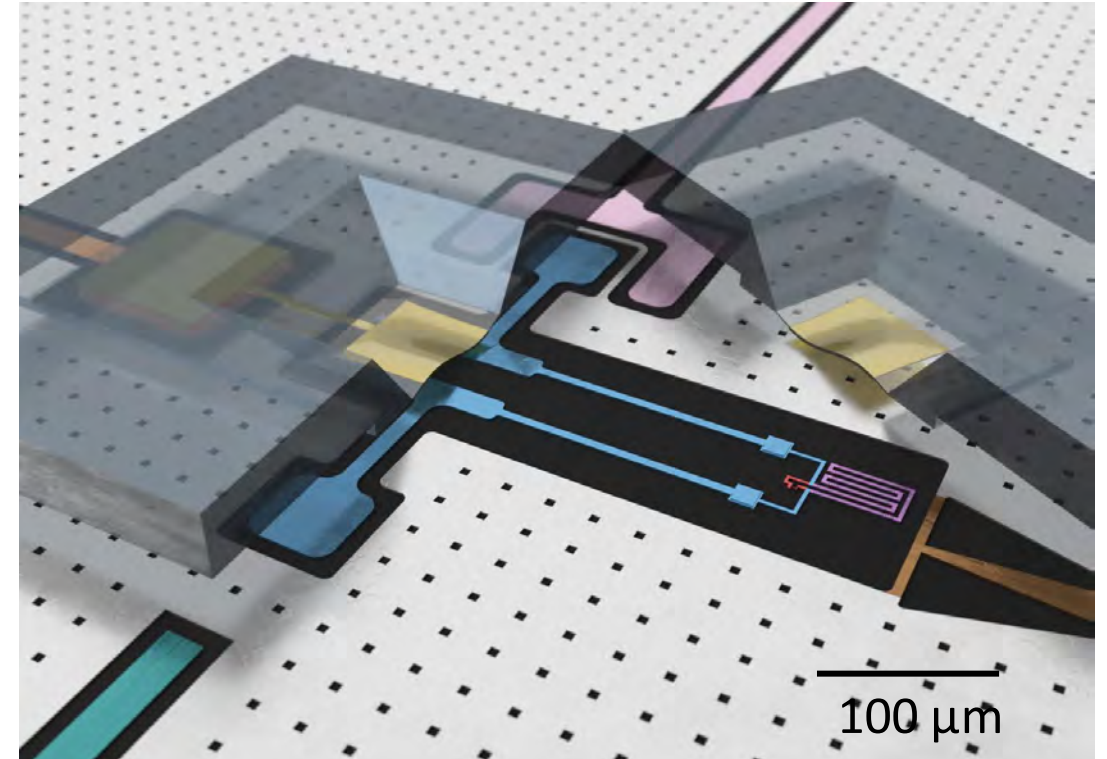


from Cavity Quantum Electro Dynamics ...

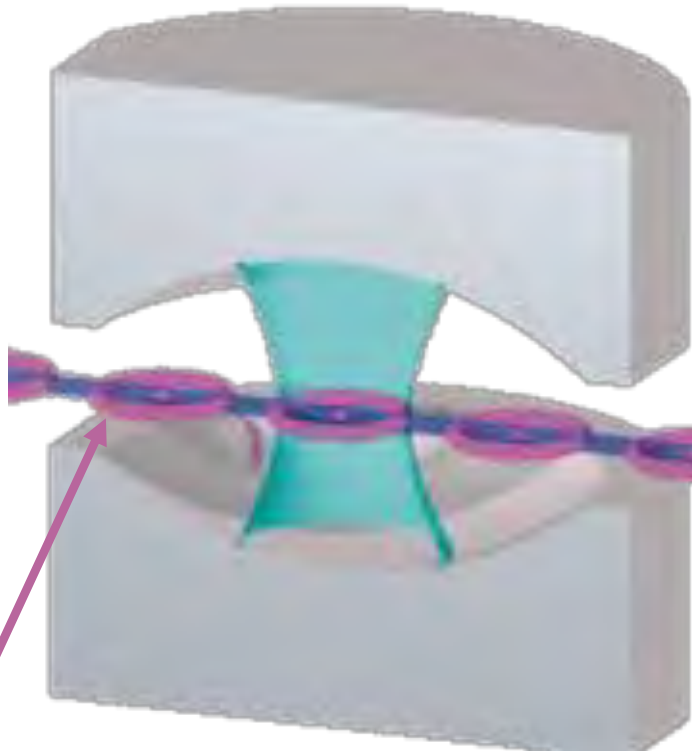


Individual Rydberg atoms

...to circuit Quantum Acousto Dynamics



from Cavity Quantum Electro Dynamics ...

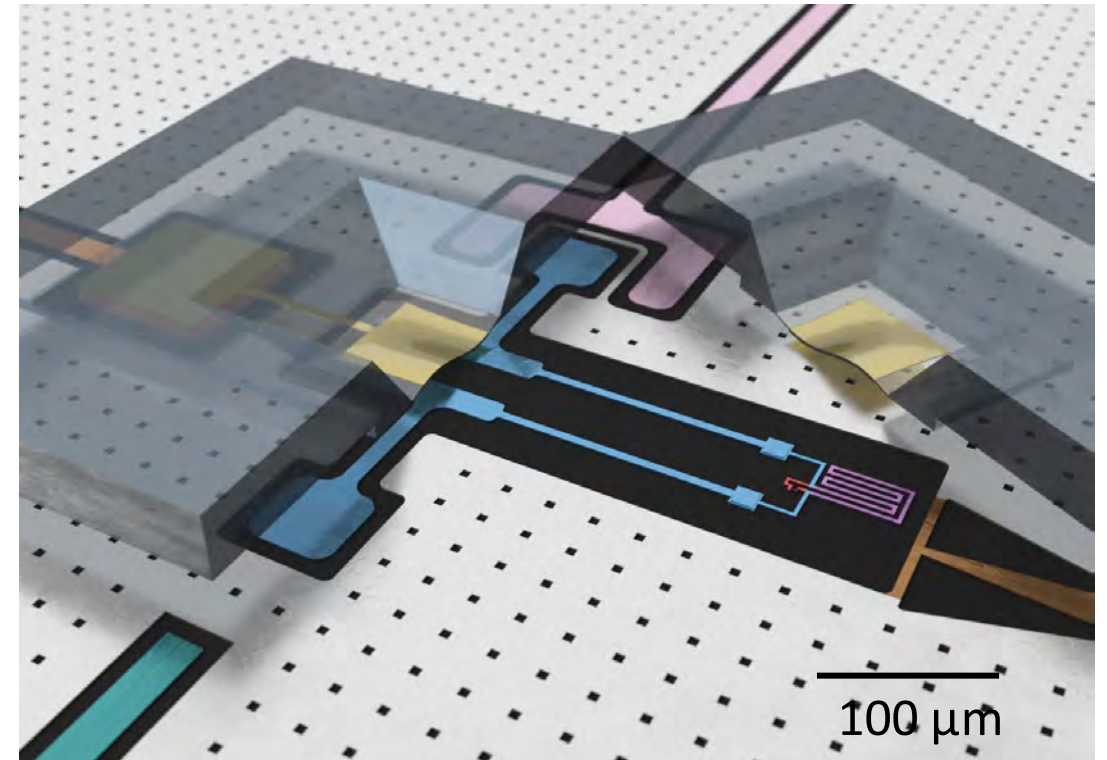


Individual Rydberg atoms

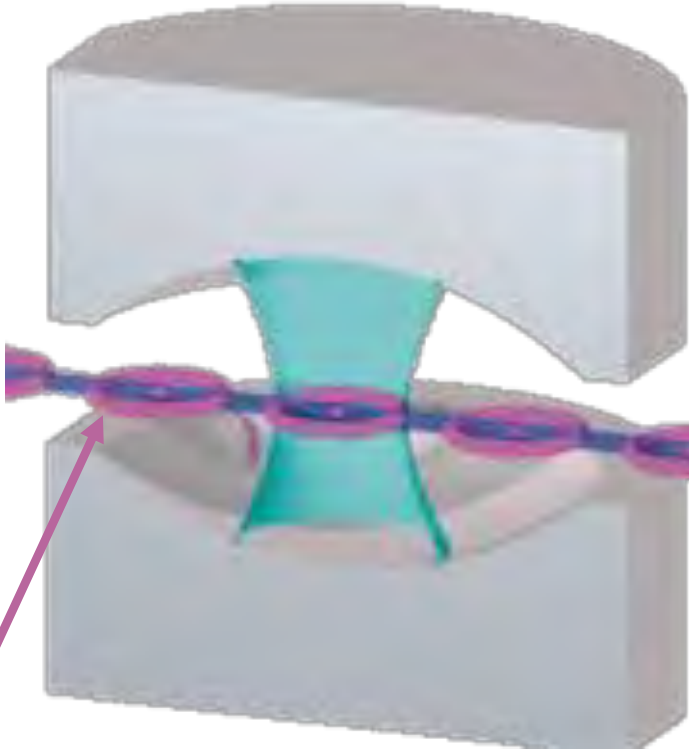
~ 120 atoms per  $T_1$

- Strong coupling regime

...to circuit Quantum Acousto Dynamics



from Cavity Quantum Electro Dynamics ...

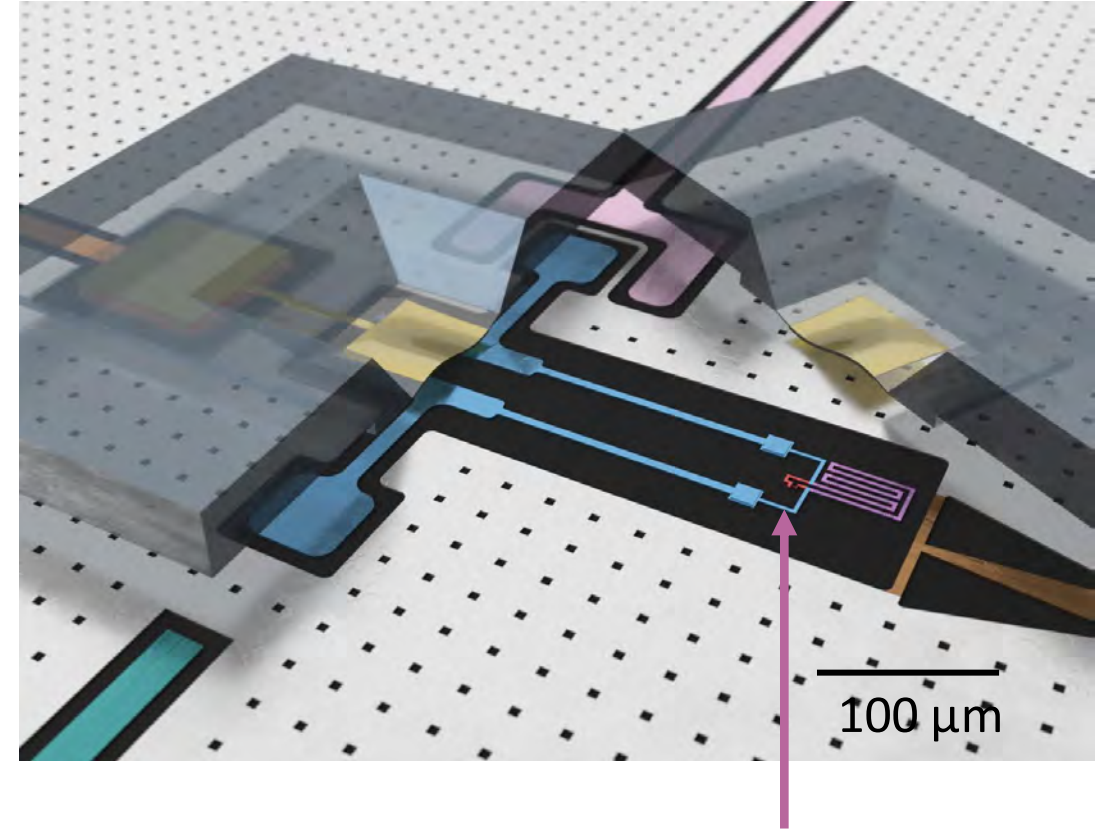


Individual Rydberg atoms

~ 120 atoms per  $T_1$

- Strong coupling regime

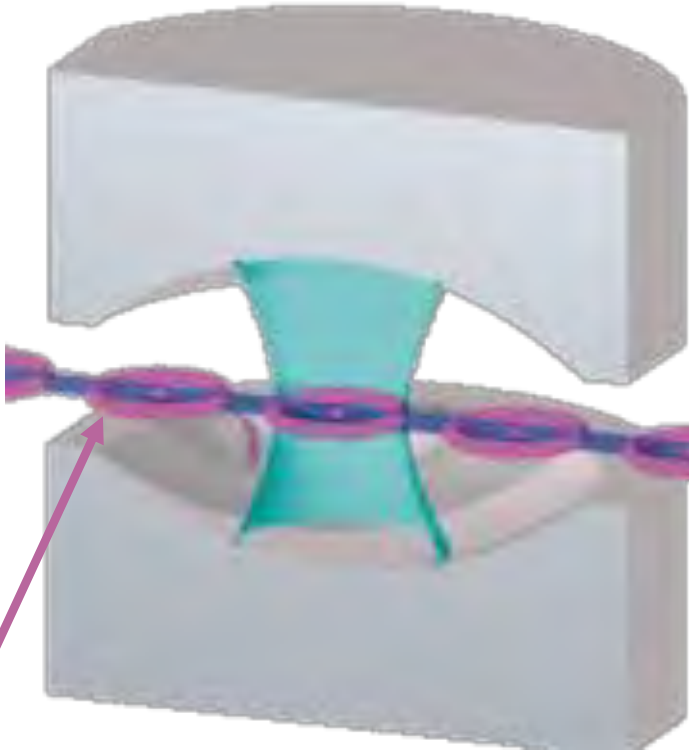
...to circuit Quantum Acousto Dynamics



Artificial atom (Fluxonium qubit)



from Cavity Quantum Electro Dynamics ...

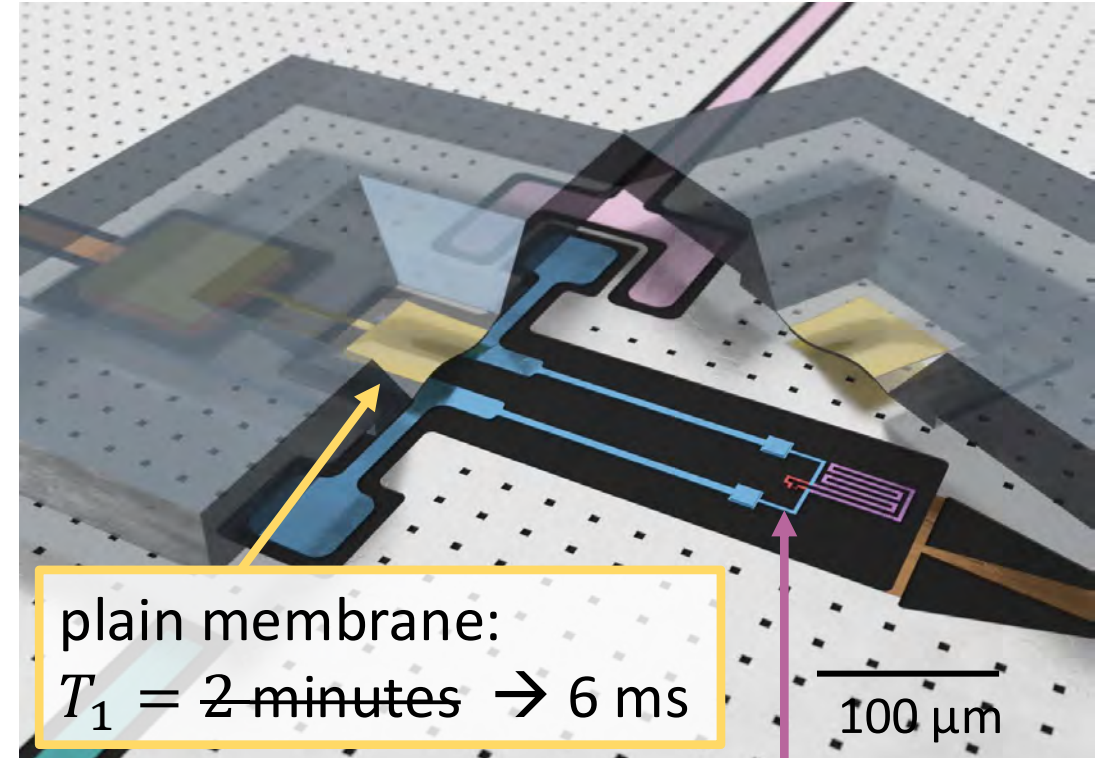


Individual Rydberg atoms

~ 120 atoms per  $T_1$

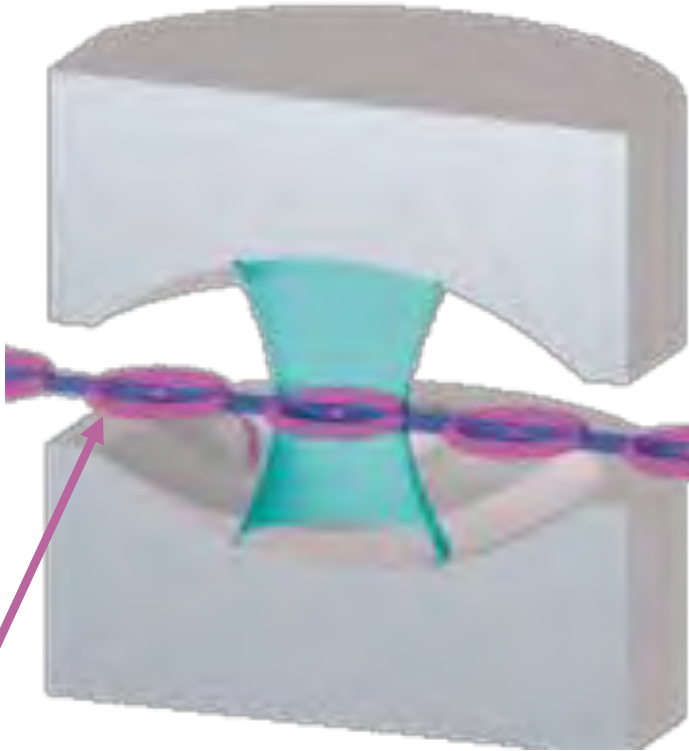
- Strong coupling regime

...to circuit Quantum Acousto Dynamics



Artificial atom (Fluxonium qubit)

from Cavity Quantum Electro Dynamics ...

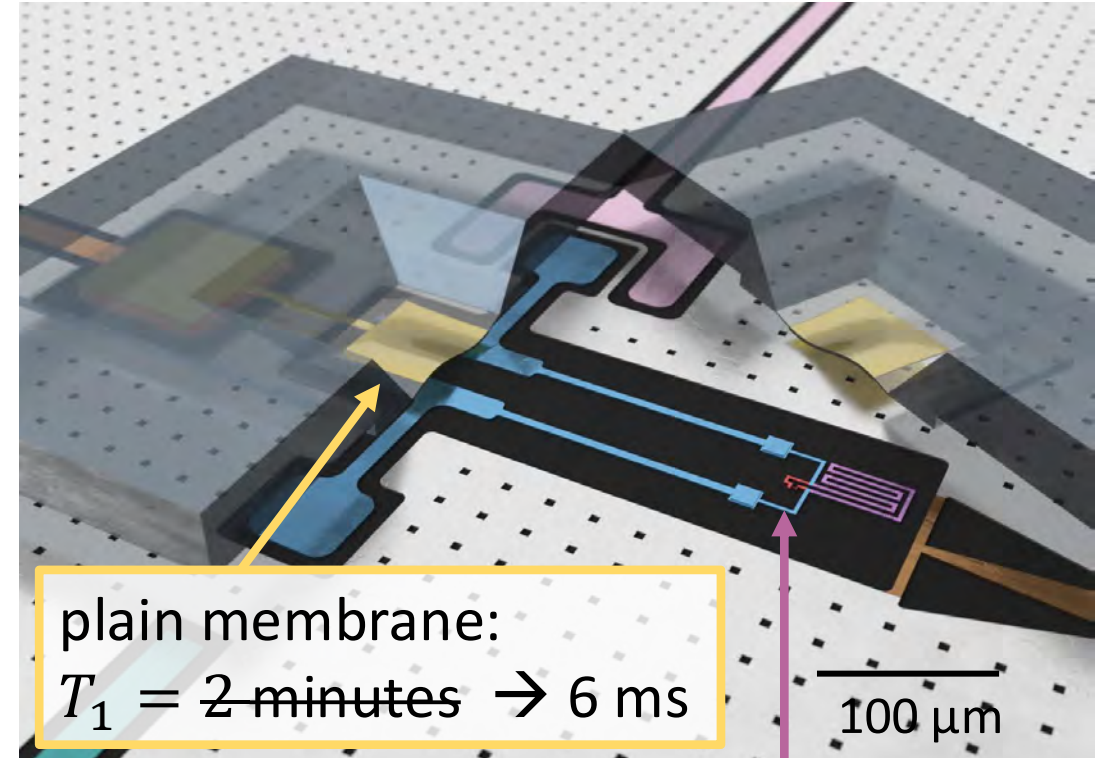


Individual Rydberg atoms

~ 120 atoms per  $T_1$

- Strong coupling regime

...to circuit Quantum Acousto Dynamics



plain membrane:  
 $T_1 = 2 \text{ minutes} \rightarrow 6 \text{ ms}$

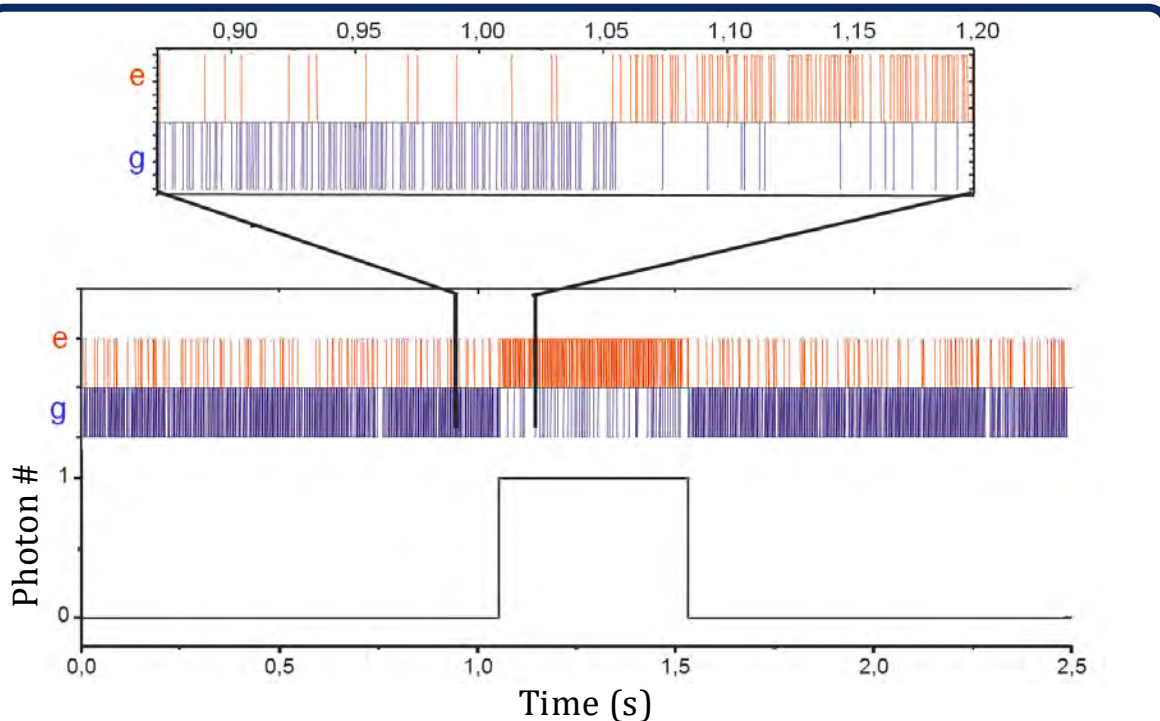
Artificial atom (Fluxonium qubit)

~ 400 qubit/membrane interactions per  $T_1$

- Weak coupling regime

from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state

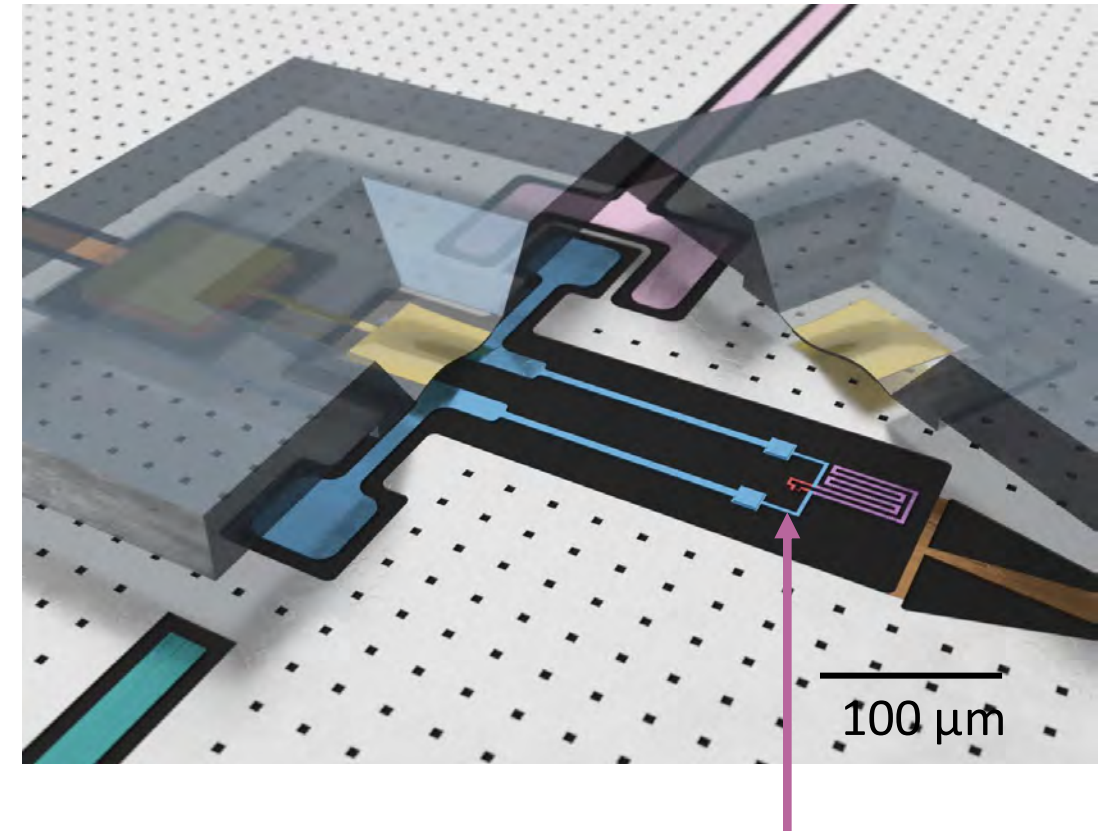


Birth and death of a photon in a cavity

Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

- Strong coupling regime

...to circuit Quantum Acousto Dynamics



Artificial atom (Fluxonium qubit)

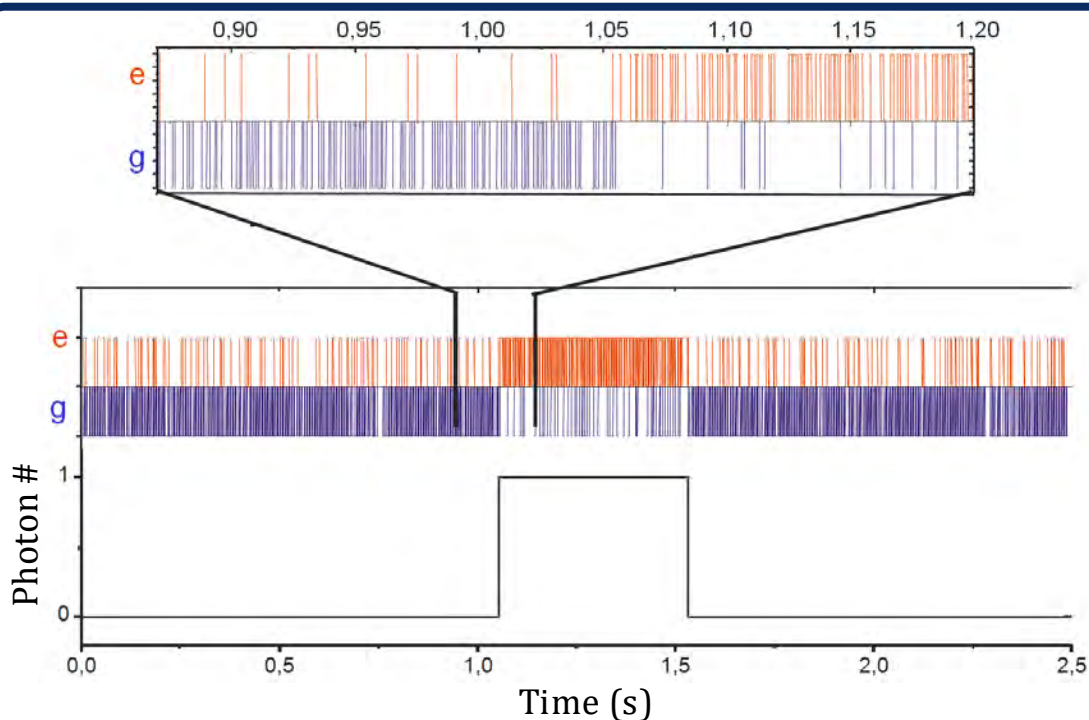
~ 400 qubit/membrane interactions per  $T_1$

- Weak coupling regime



## from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state



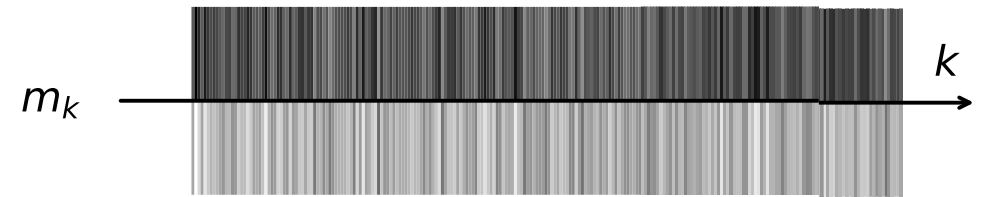
Birth and death of a photon in a cavity

Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

- Strong coupling regime

## ...to circuit Quantum Acousto Dynamics

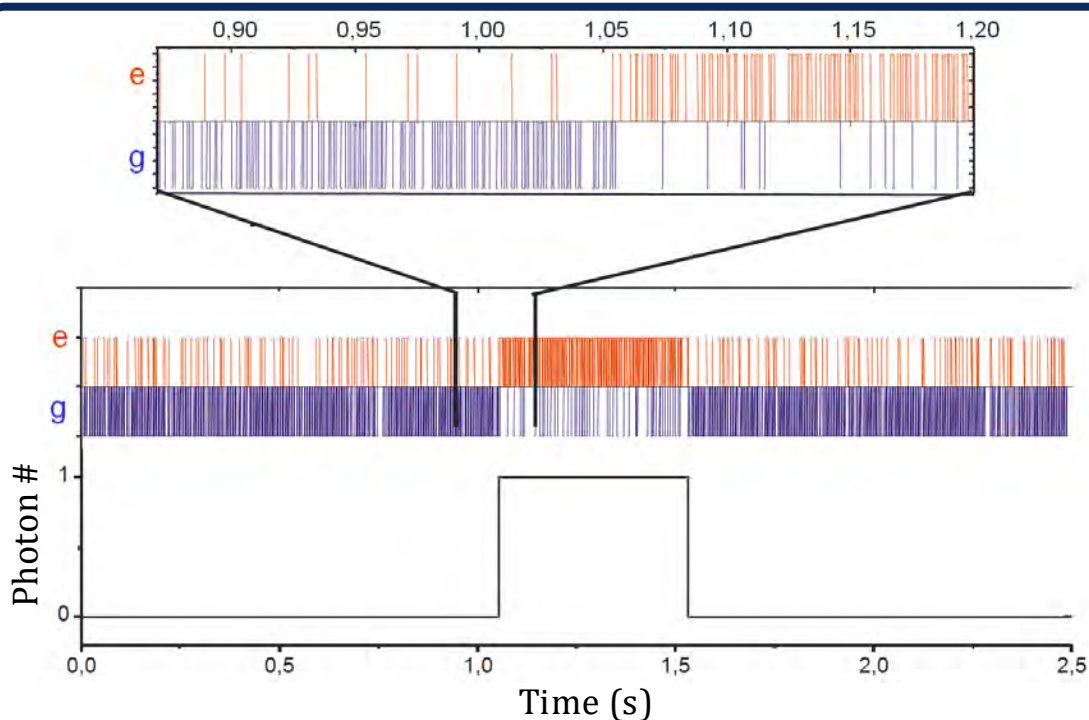
Qubit measurements only yield partial information upon the mechanical state



- Weak coupling regime

## from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state



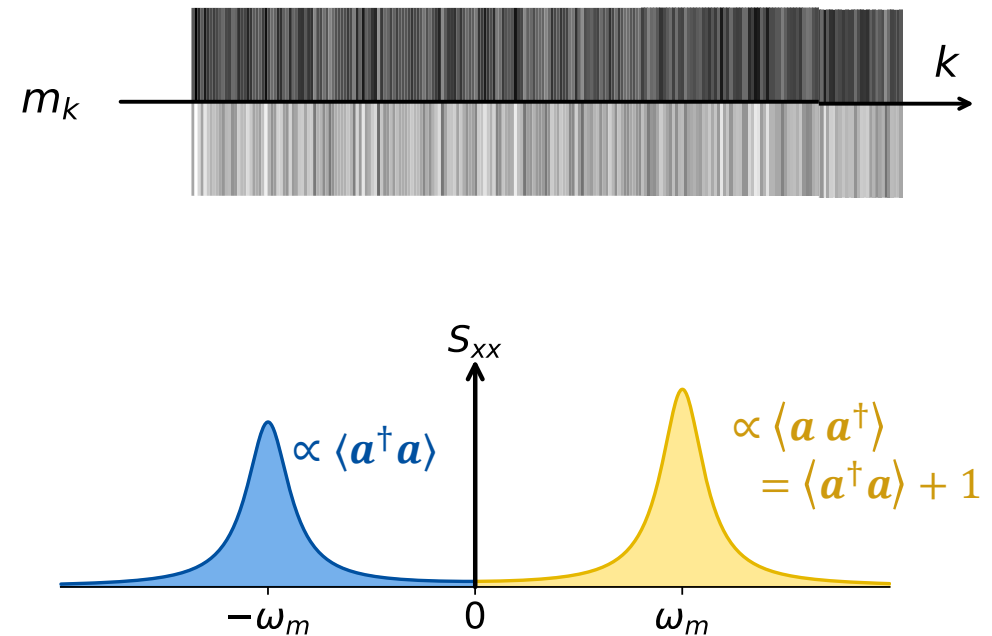
Birth and death of a photon in a cavity

Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

- Strong coupling regime

## ...to circuit Quantum Acousto Dynamics

Qubit measurements only yield partial information upon the mechanical state



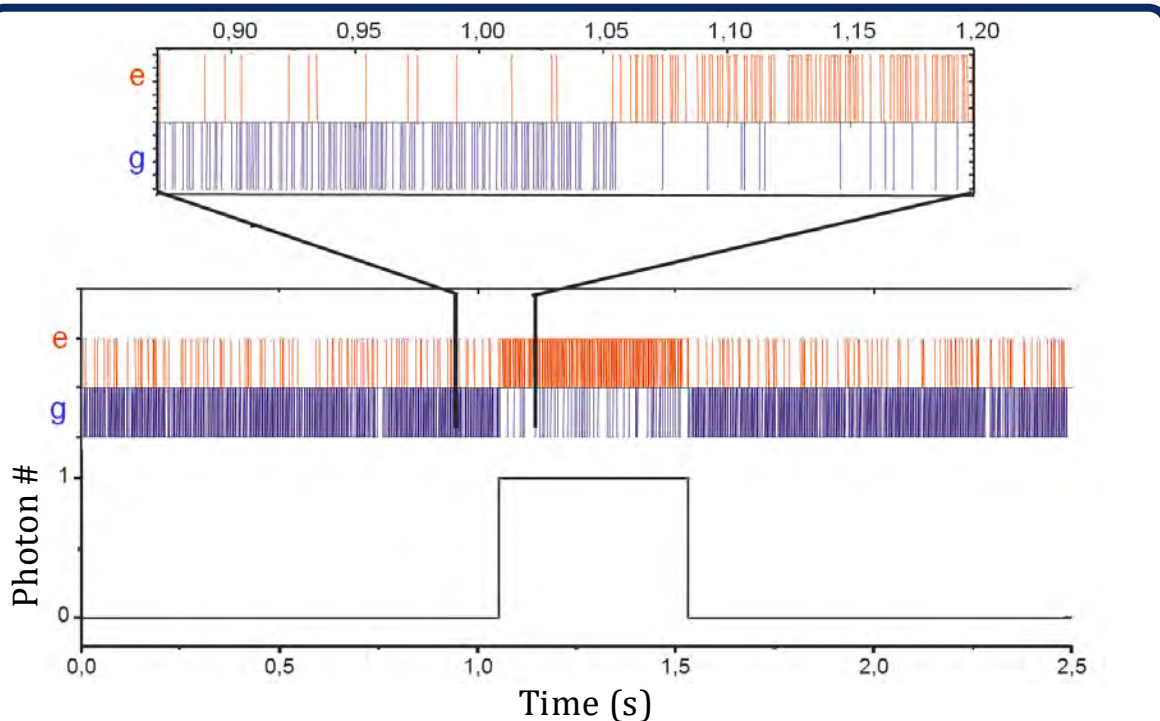
Quantum spectrum asymmetry

Gerashchenko *et al.*, arXiv 2505.21481 (2025).

- Weak coupling regime

## from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state



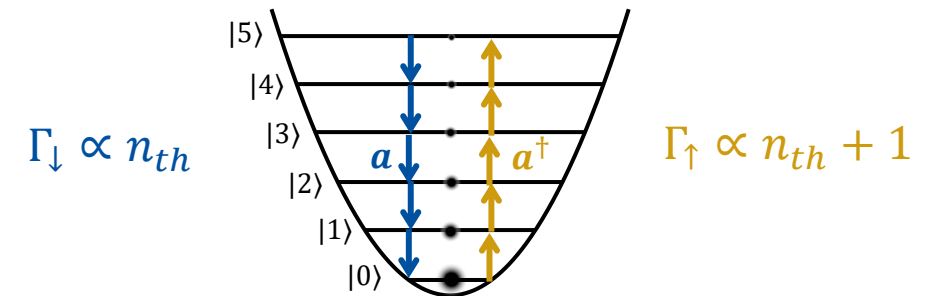
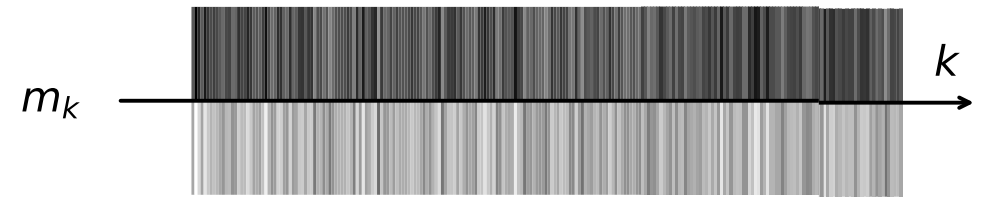
Birth and death of a photon in a cavity

Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

- Strong coupling regime

## ...to circuit Quantum Acousto Dynamics

Qubit measurements only yield partial information upon the mechanical state



Quantum spectrum asymmetry

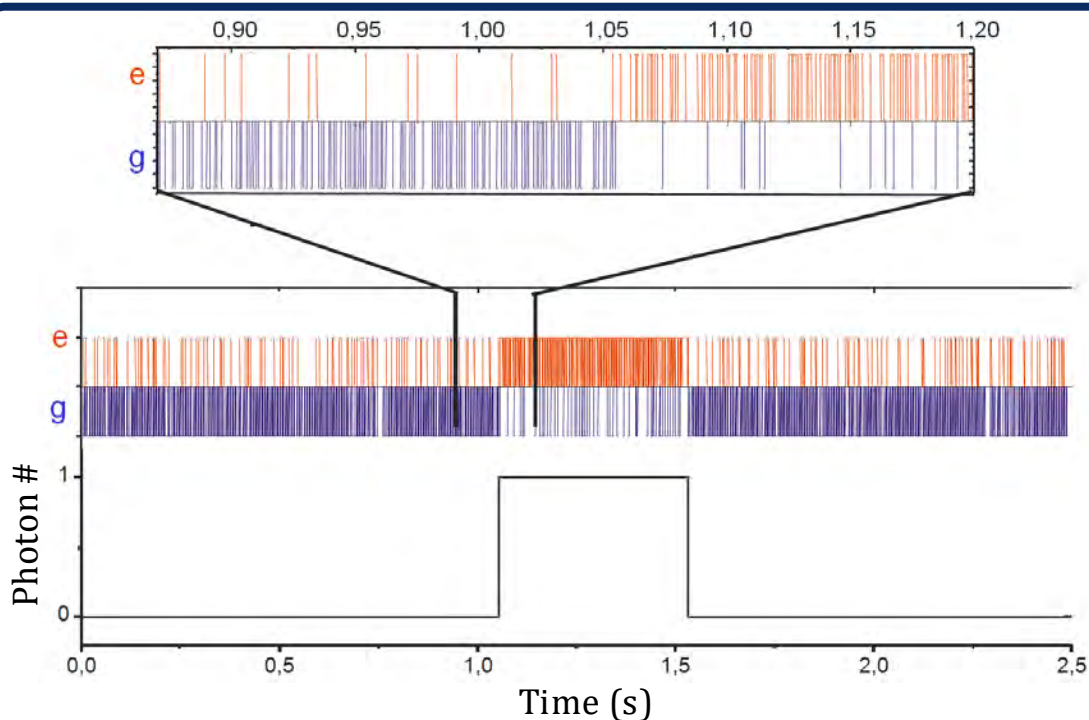
Gerashchenko *et al.*, arXiv 2505.21481 (2025).

- Weak coupling regime



## from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state



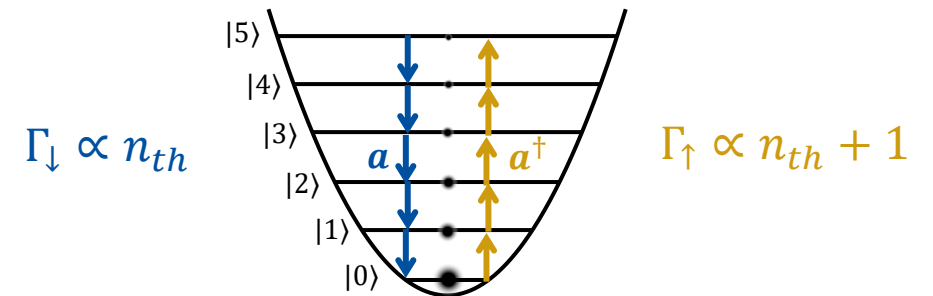
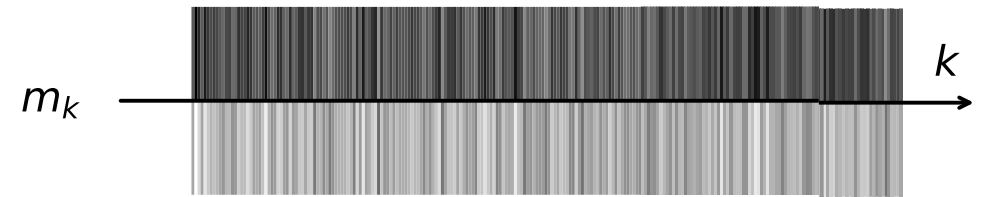
Birth and death of a photon in a cavity

Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

- Strong coupling regime

## ...to circuit Quantum Acousto Dynamics

Qubit measurements only yield partial information upon the mechanical state



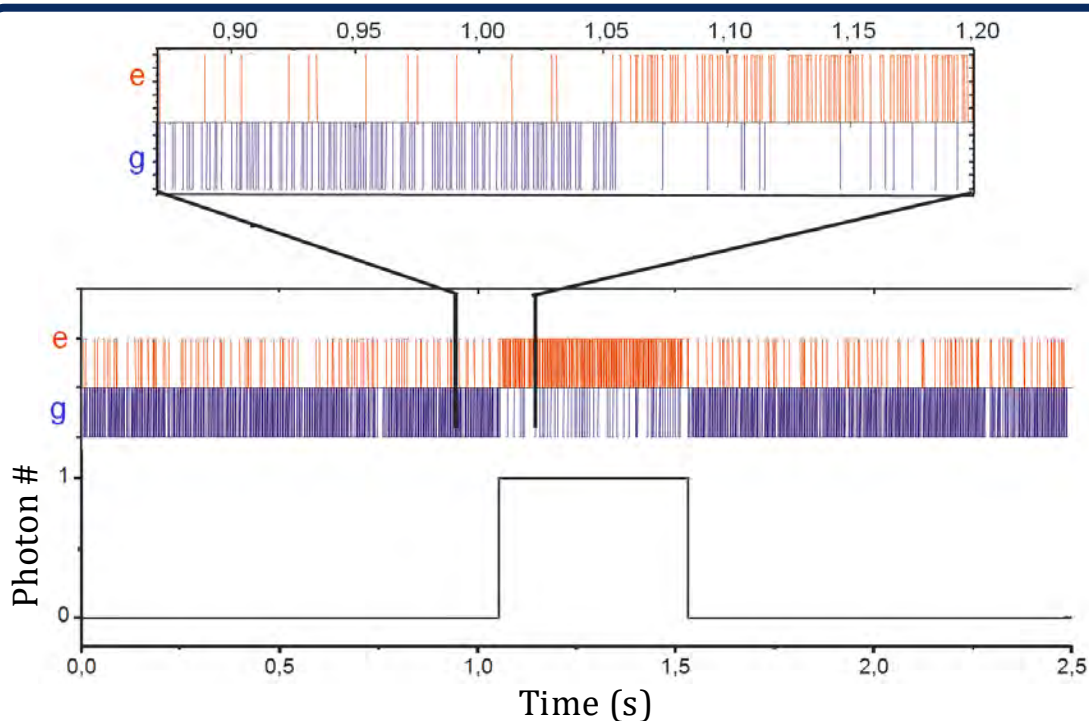
Quantum spectrum asymmetry

Gerashchenko *et al.*, arXiv 2505.21481 (2025).

Goal: experimental test of spontaneous collapse models

from Cavity Quantum Electro Dynamics ...

Each atom performs a strong measurement of the resonator state



Birth and death of a photon in a cavity

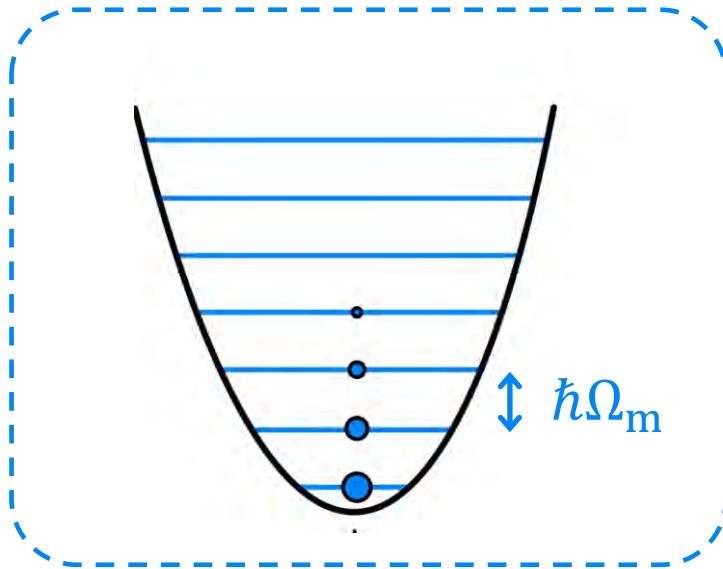
Gleyzes *et al.* *Nature*, 446(7133), 297-300 (2007).

...to circuit Quantum Acousto Dynamics

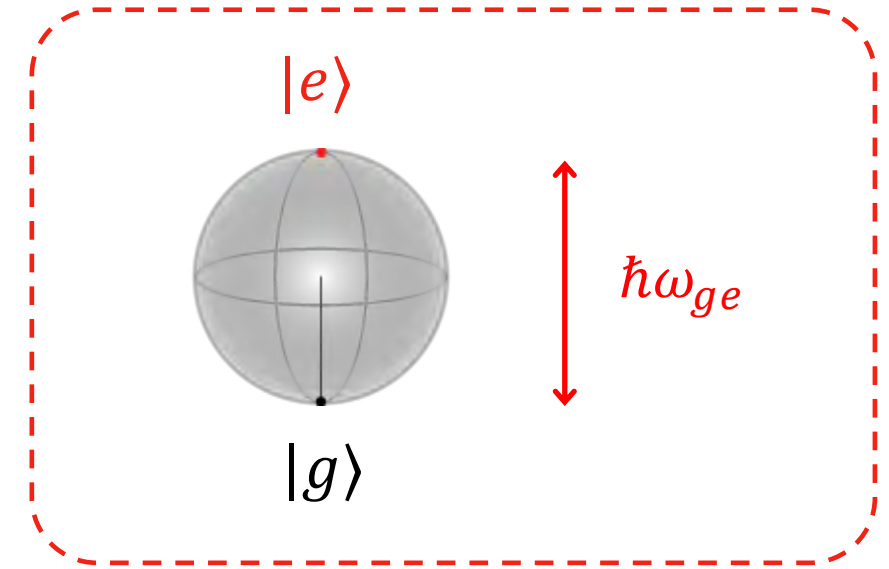
$$|\Psi\rangle = |\alpha\rangle + |-\alpha\rangle$$



Goal: experimental test of spontaneous collapse models



**Bulk acoustic mode @ 6 GHz**

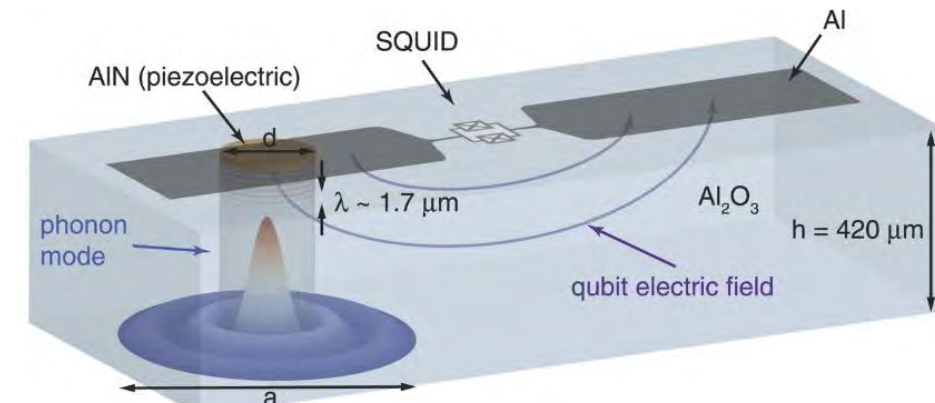


**Superconducting qubit @ 6 GHz**

## Schrödinger cat states of a 16-microgram mechanical oscillator

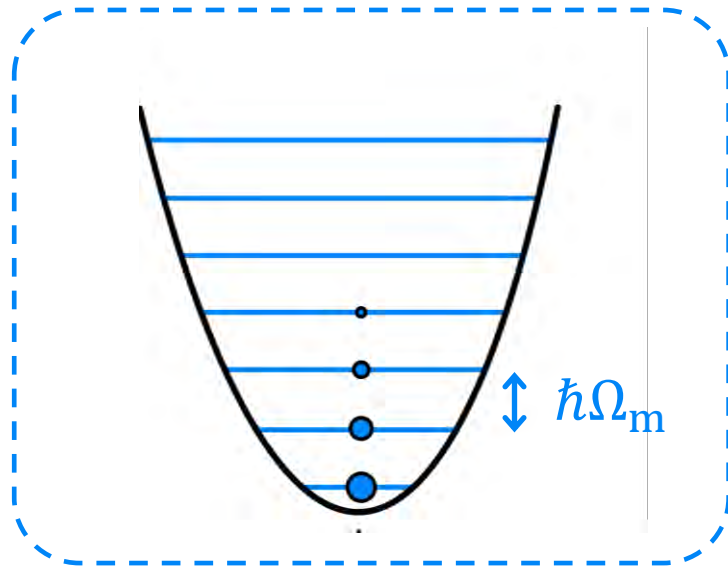
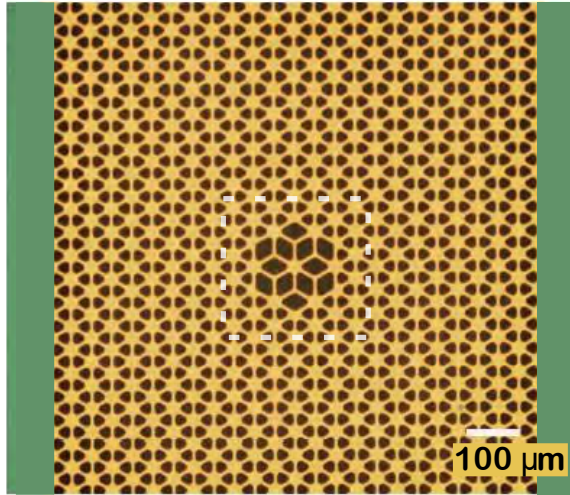
Bild et al, Vol 380, Issue 6642 pp. 274-278, Science 2023

Yiwen Chu group @ ETH

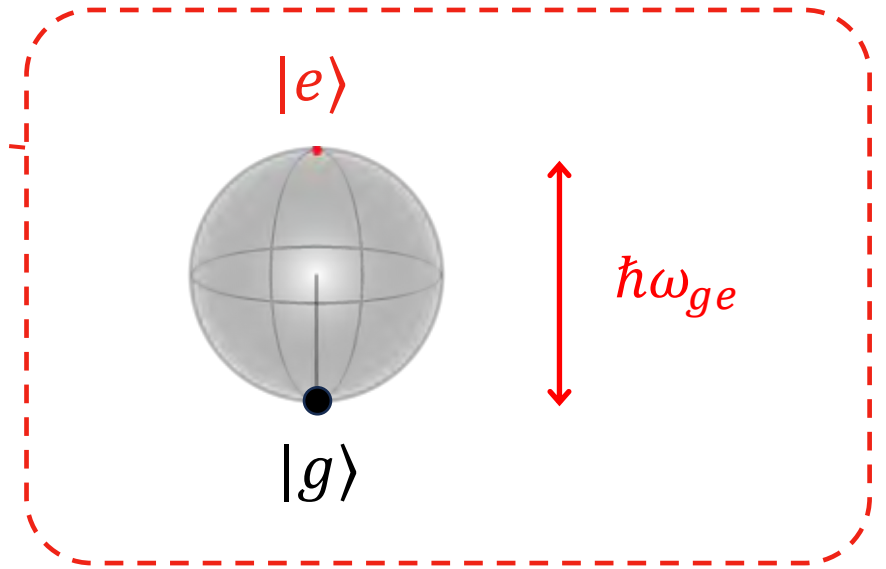
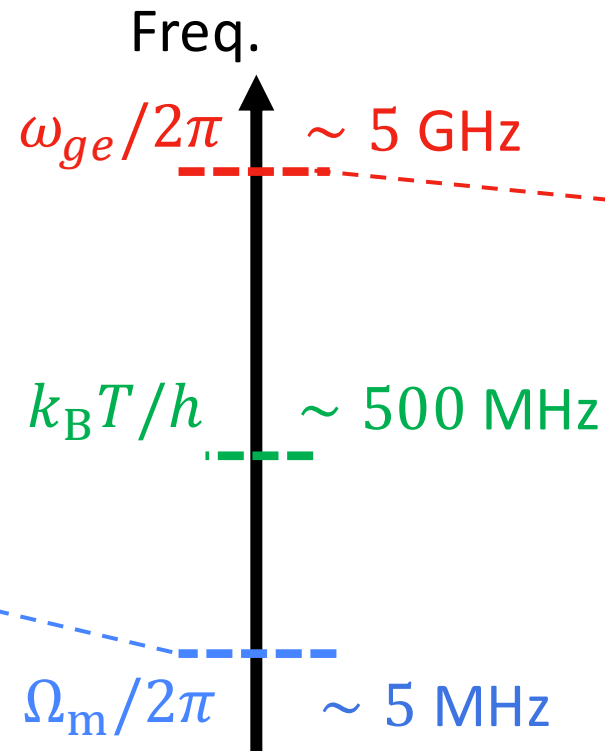




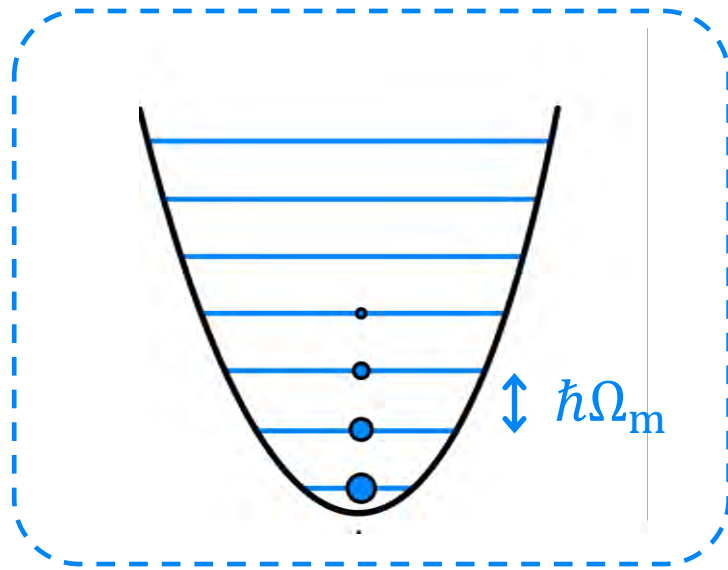
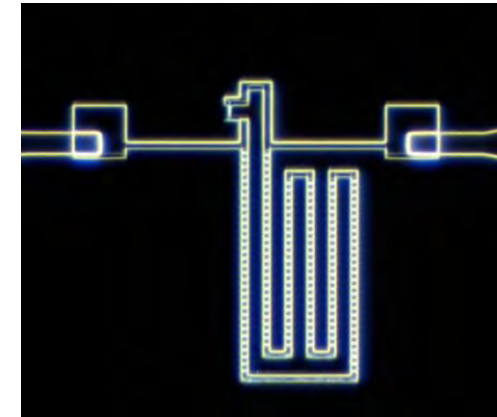
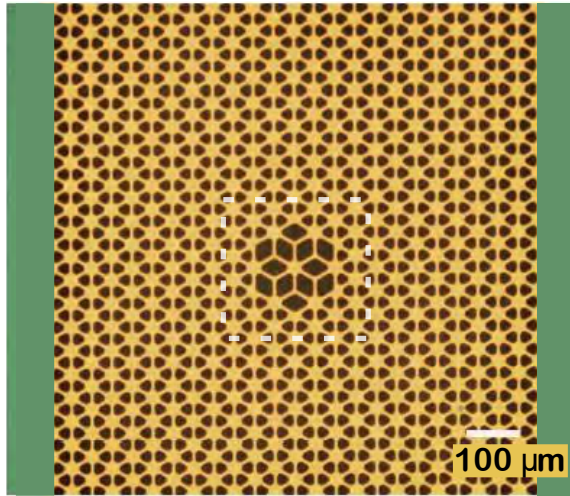
# Coupling a MHz mechanics to a superconducting qubit



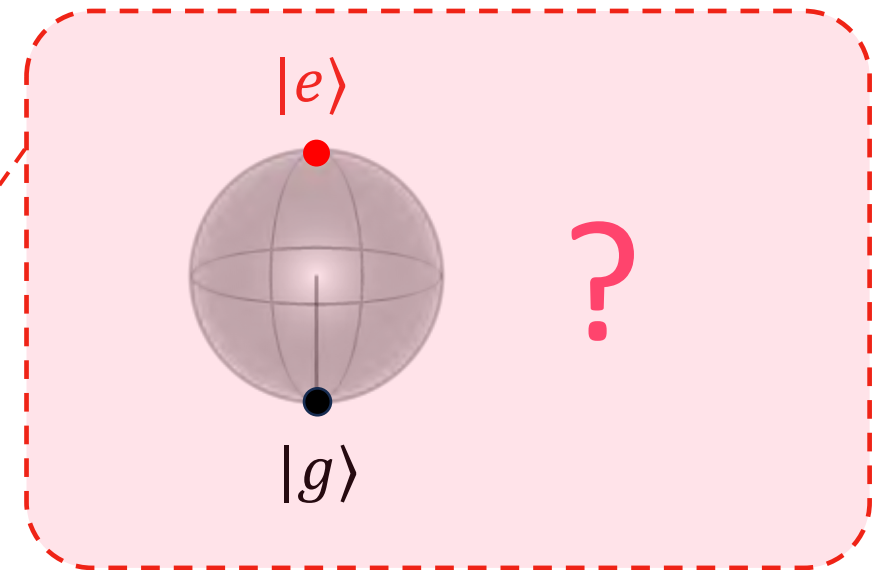
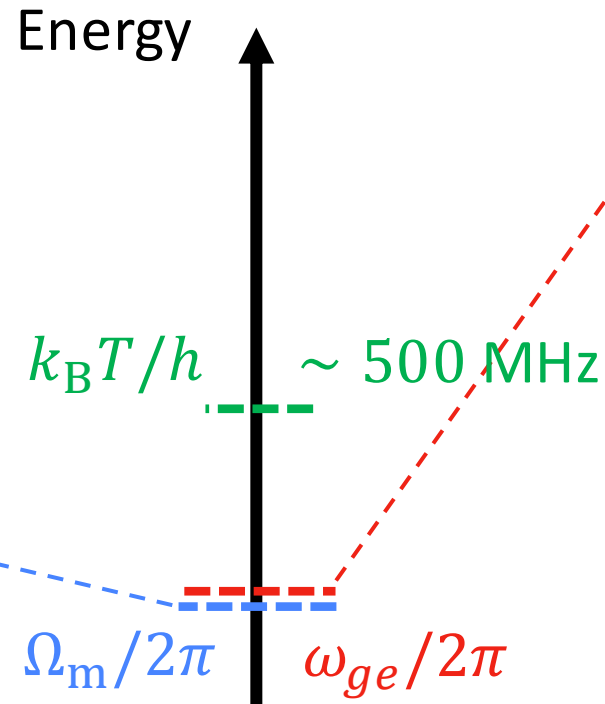
Ultra high-Q resonator @ 5 MHz



Superconducting qubit @ GHz

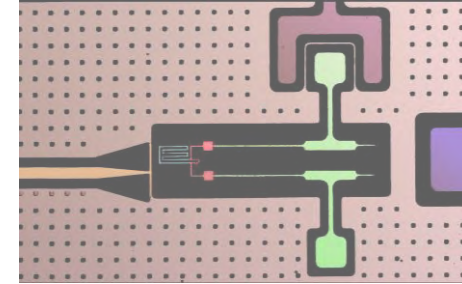


Ultra high-Q resonator @ 5 MHz



Superconducting qubit @ 5 MHz

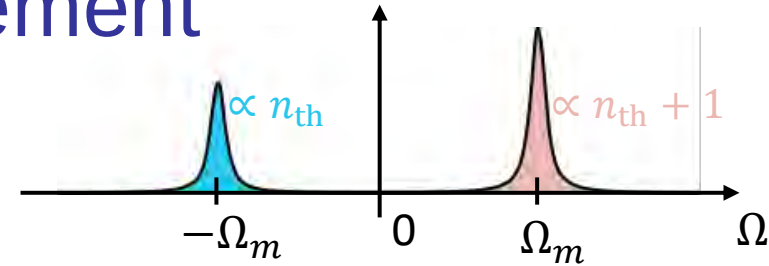
- Radio-frequency fluxonium qubits



- Hybrid qubit/mechanical system: MecaFlux



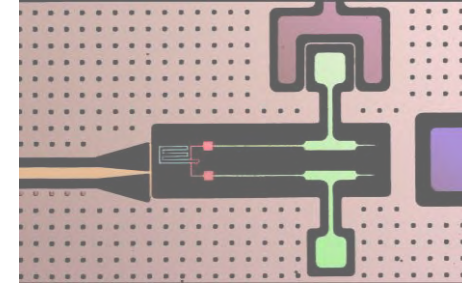
- Quantum asymmetry measurement



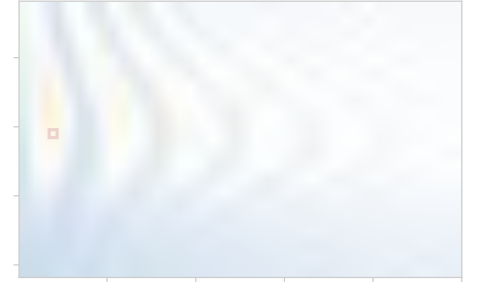
- Perspectives



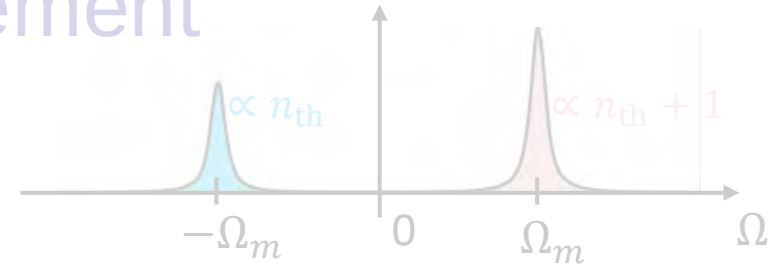
- Radio-frequency fluxonium qubits



- Hybrid qubit/mechanical system: MecaFlux



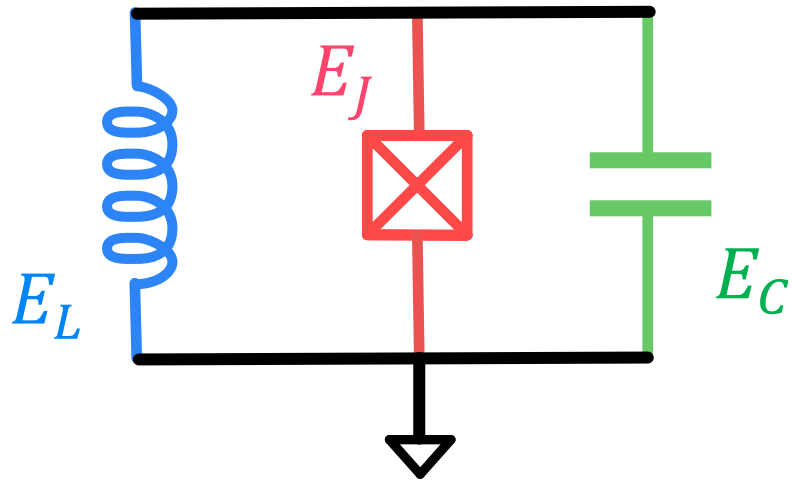
- Quantum asymmetry measurement

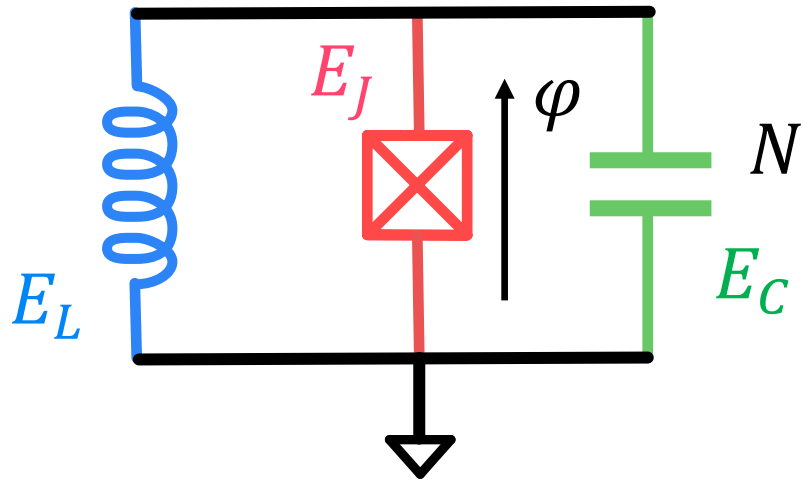


- Perspectives









Superconducting  
phase difference  
across the junction:

$$\varphi = \frac{\phi}{\Phi_0}$$

*(position-like)*

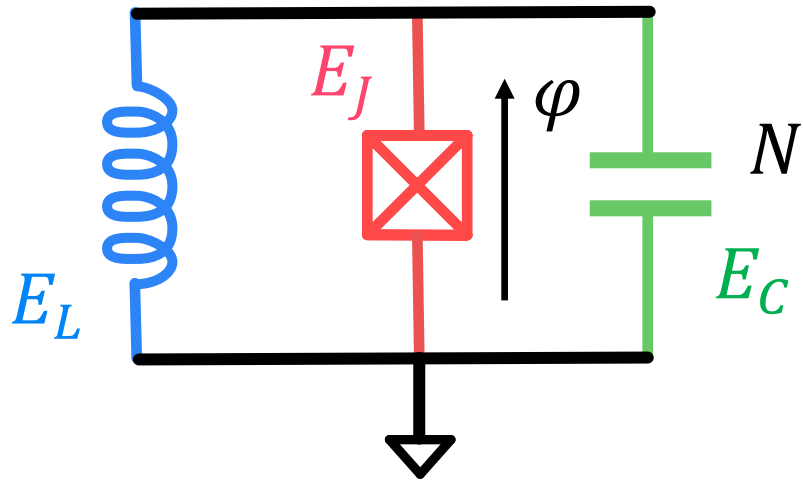
Displacement  
charge between  
junction electrodes:

$$N = \frac{Q}{2e}$$

*(momentum-like)*

# The (heavy) fluxonium

$$\hat{H} = \underbrace{\frac{1}{2} E_L \hat{\varphi}^2}_{\text{Potential energy}} - \underbrace{E_J \cos(\hat{\varphi} - \varphi_{\text{ext}})}_{\text{Potential energy}} + 4E_C \hat{N}^2$$



Superconducting  
phase difference  
across the junction:

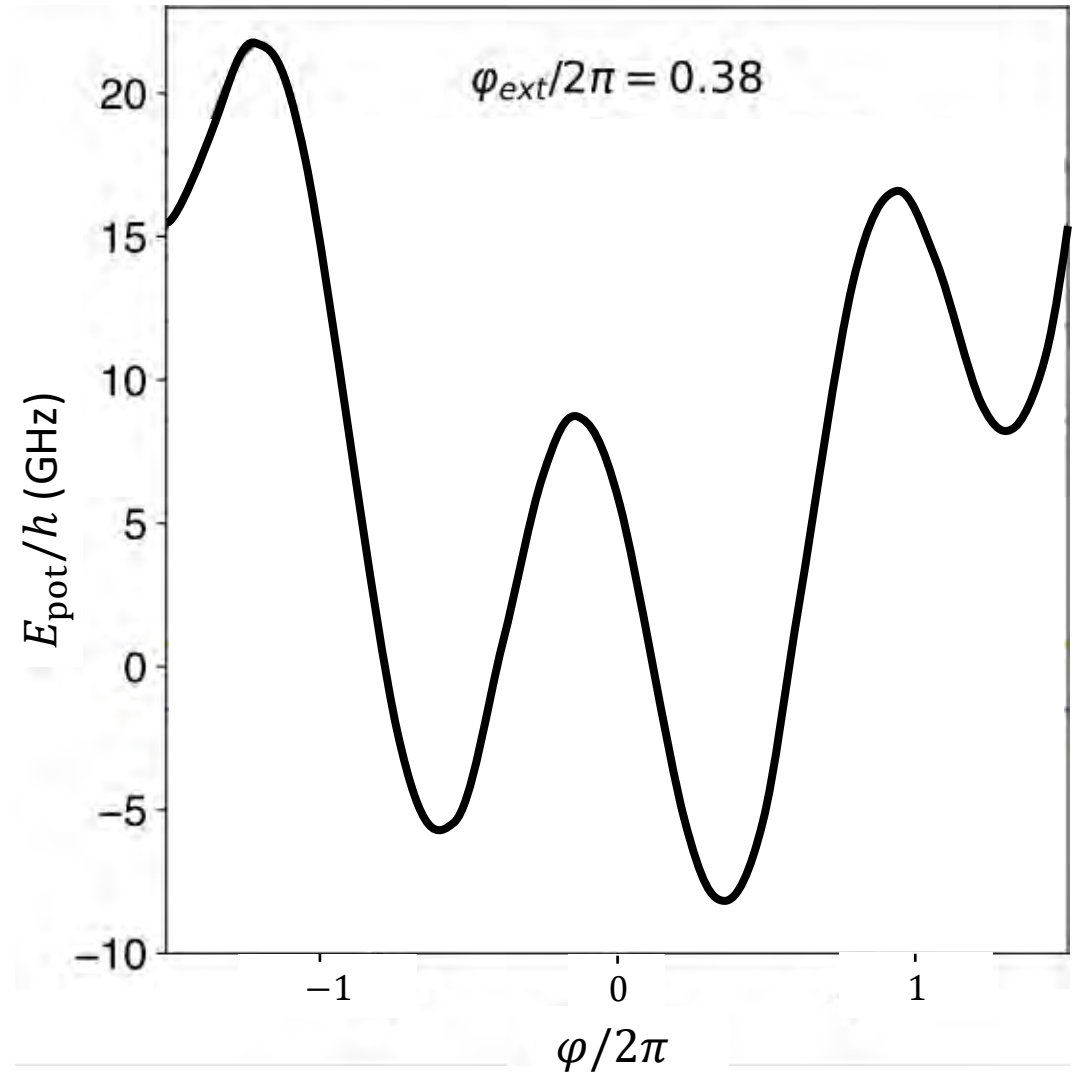
$$\varphi = \frac{\phi}{\Phi_0}$$

(position-like)

Displacement  
charge between  
junction electrodes:

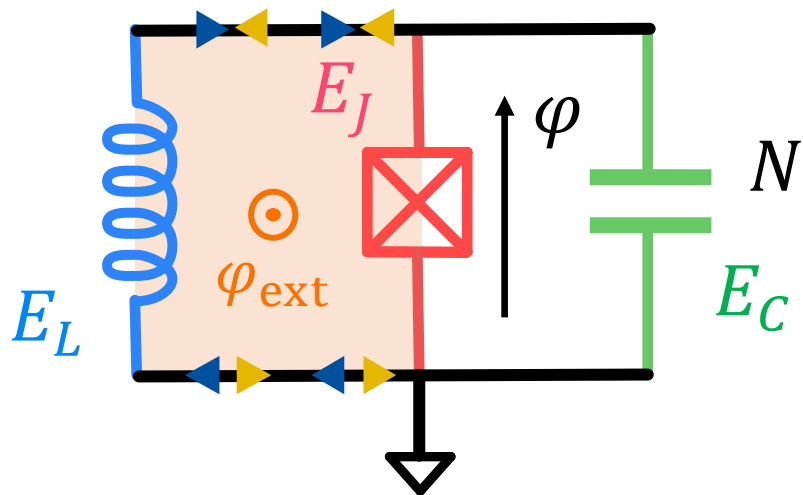
$$N = \frac{Q}{2e}$$

(momentum-like)



# The (heavy) fluxonium

$$\hat{H} = \underbrace{\frac{1}{2} E_L \hat{\varphi}^2}_{\text{Potential energy}} - \underbrace{E_J \cos(\hat{\varphi} - \varphi_{\text{ext}})}_{\text{Kinetic energy}} + \underbrace{4E_C \hat{N}^2}_{\text{Kinetic energy}}$$



Superconducting  
phase difference  
across the junction:

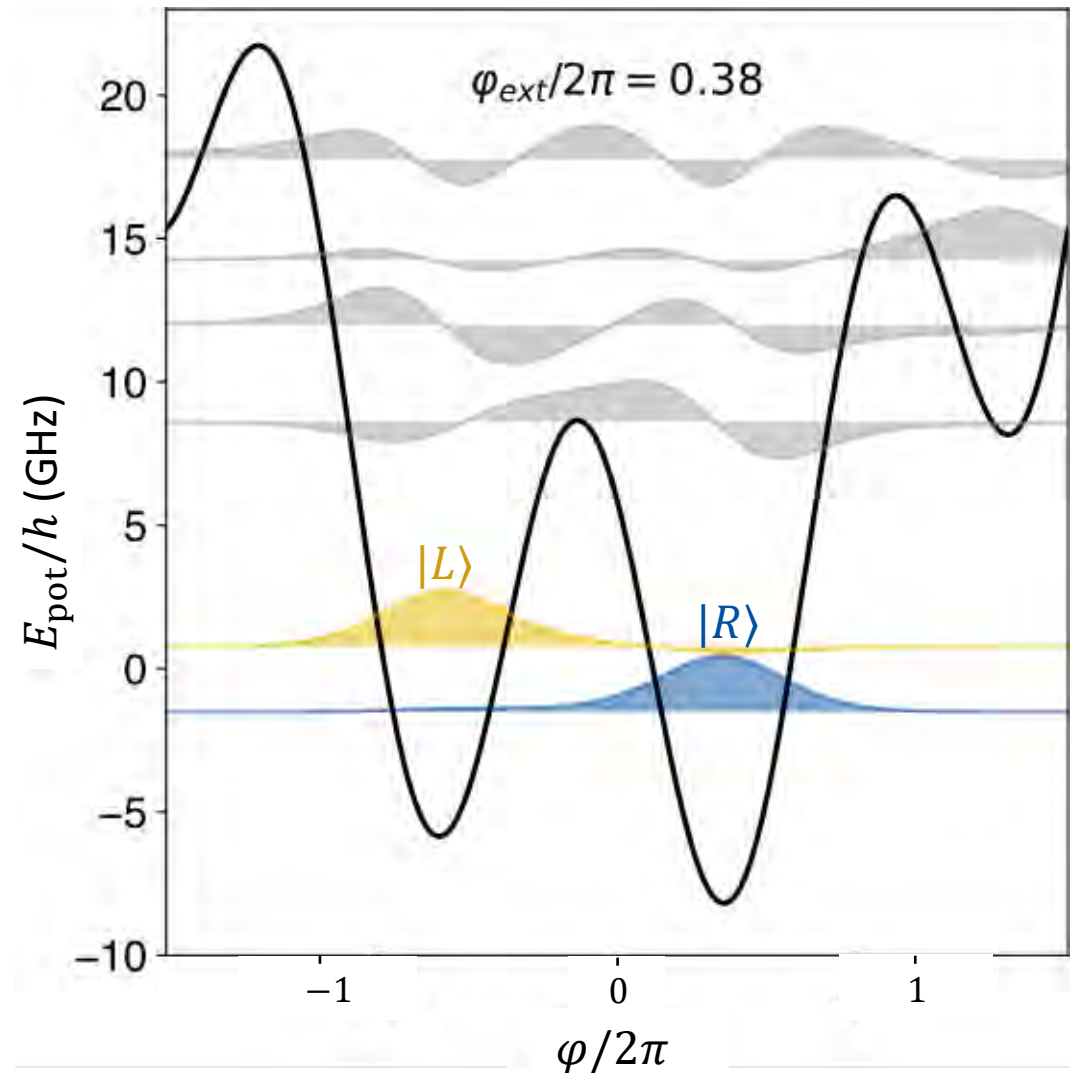
$$\varphi = \frac{\phi}{\Phi_0}$$

(position-like)

Displacement  
charge between  
junction electrodes:

$$N = \frac{Q}{2e}$$

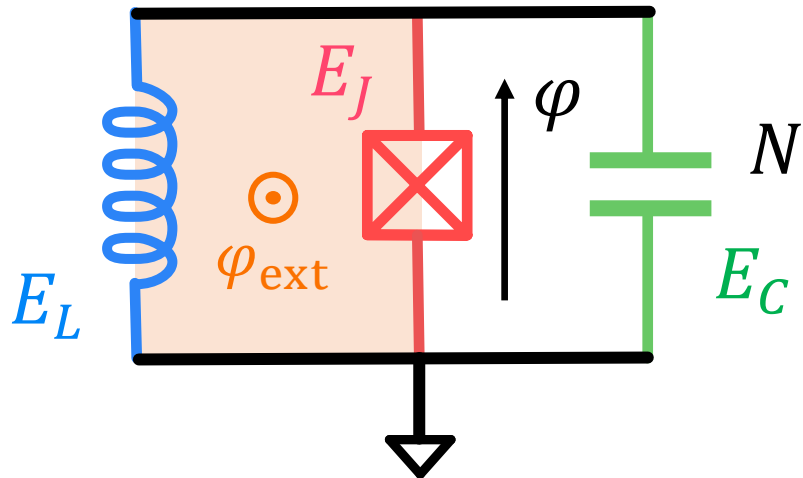
(momentum-like)





# The (heavy) fluxonium

$$\hat{H} = \underbrace{\frac{1}{2} E_L \hat{\varphi}^2 - E_J \cos(\hat{\varphi} - \varphi_{\text{ext}})}_{\text{Potential energy}} + \underbrace{4E_C \hat{N}^2}_{\text{Kinetic energy}}$$

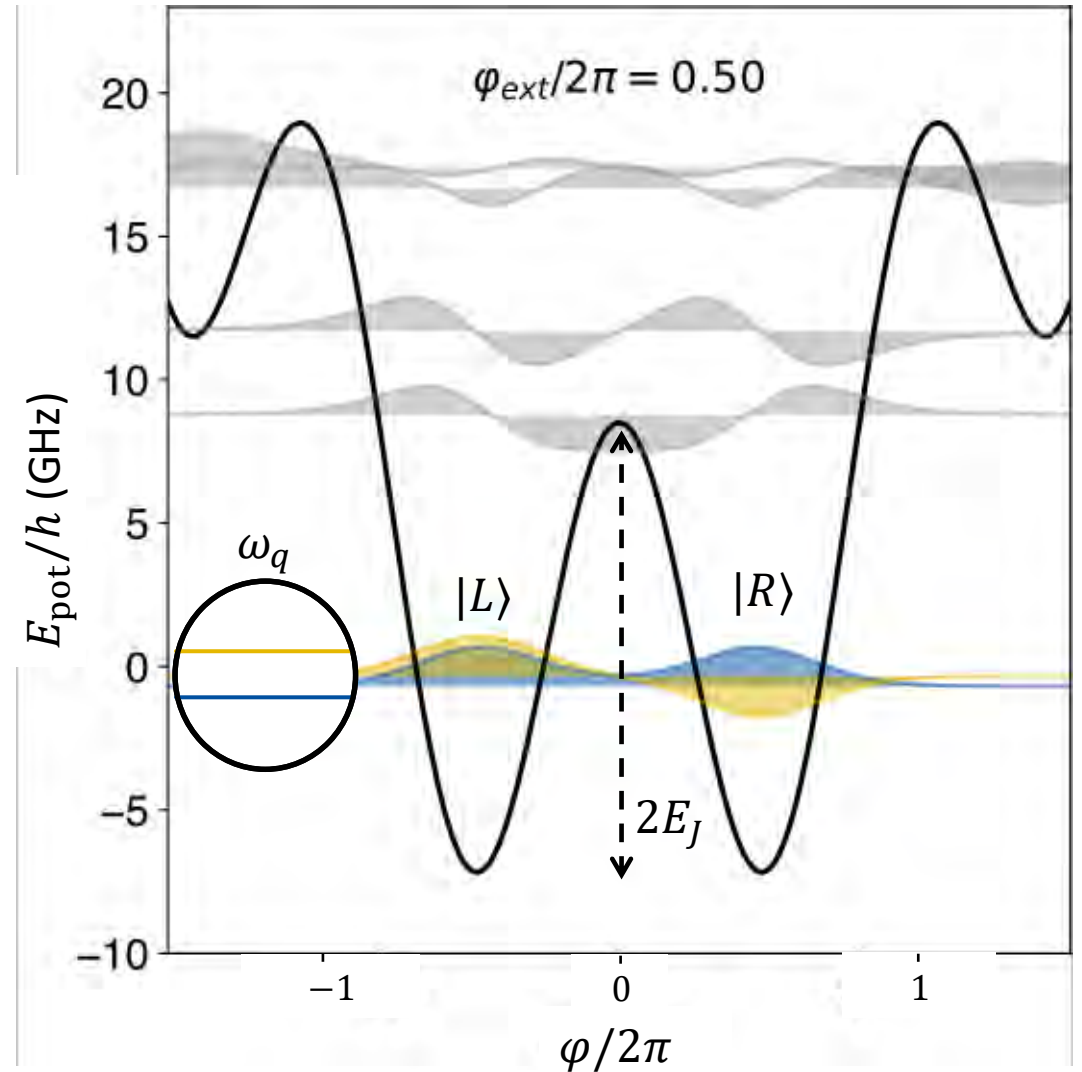


Flux frustration point ( $\varphi_{\text{ext}}/2\pi = 0.5$ ) :

The localized states hybridize:

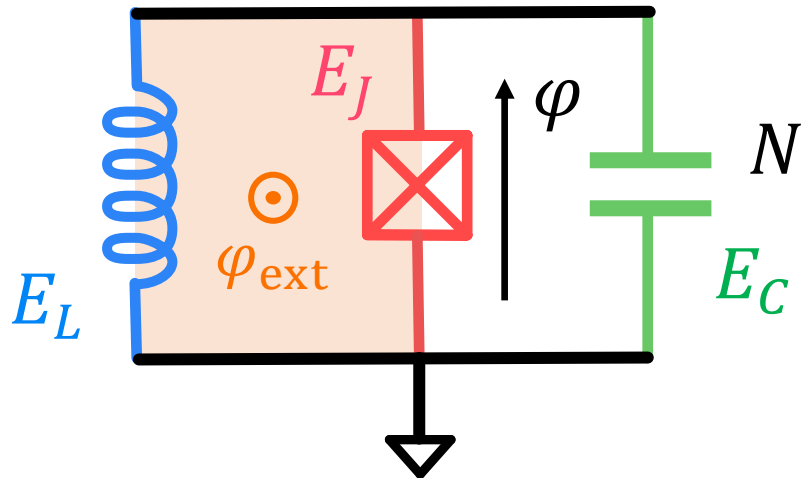
$$|e\rangle \rightarrow |L\rangle - |R\rangle$$

$$|g\rangle \rightarrow |L\rangle + |R\rangle$$



# The (heavy) fluxonium

$$\hat{H} = \underbrace{\frac{1}{2} E_L \hat{\varphi}^2 - E_J \cos(\hat{\varphi} - \varphi_{\text{ext}})}_{\text{Potential energy}} + \underbrace{4E_C \hat{N}^2}_{\text{Kinetic energy}}$$



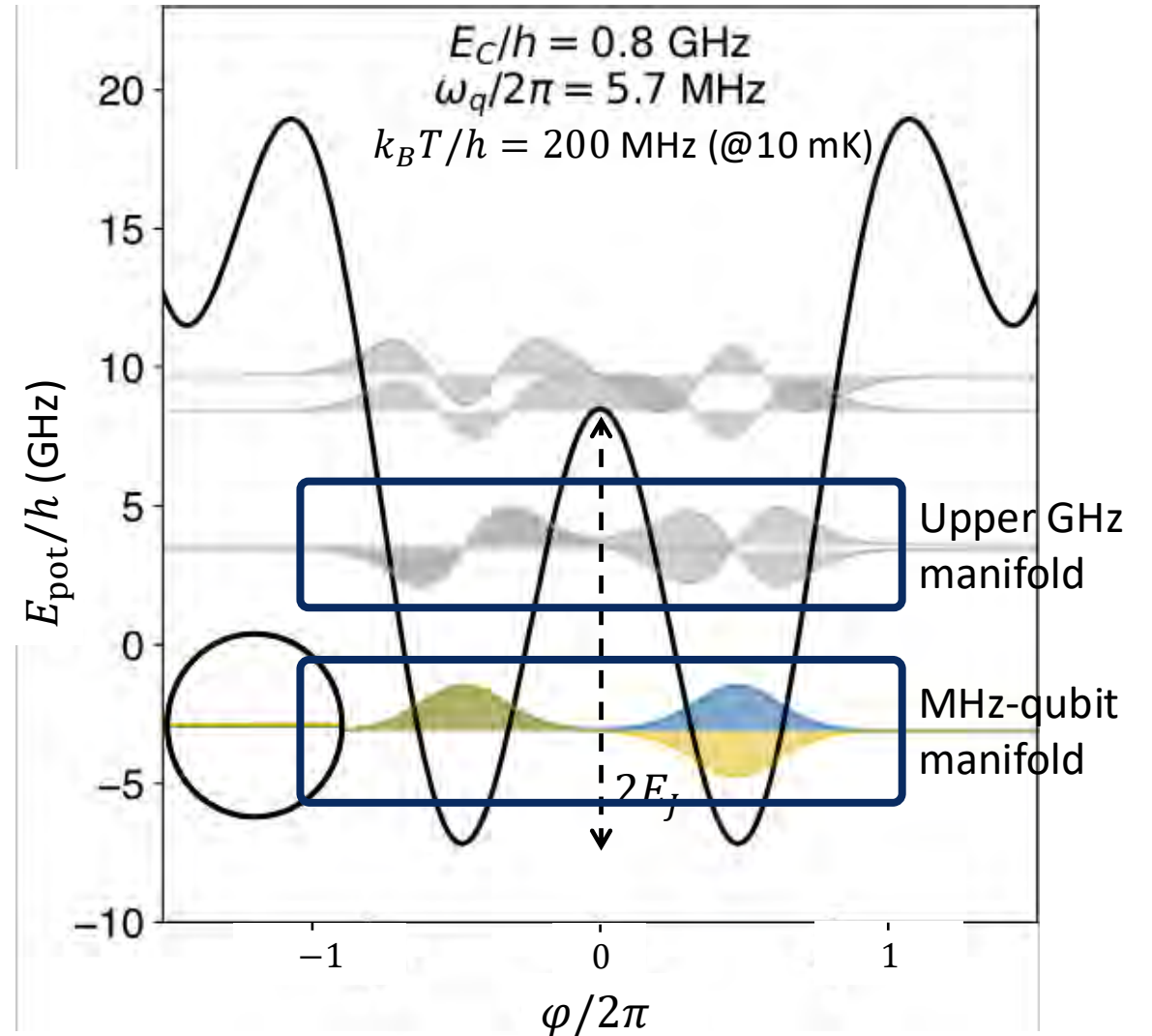
Flux frustration point ( $\varphi_{\text{ext}}/2\pi = 0.5$ ) :

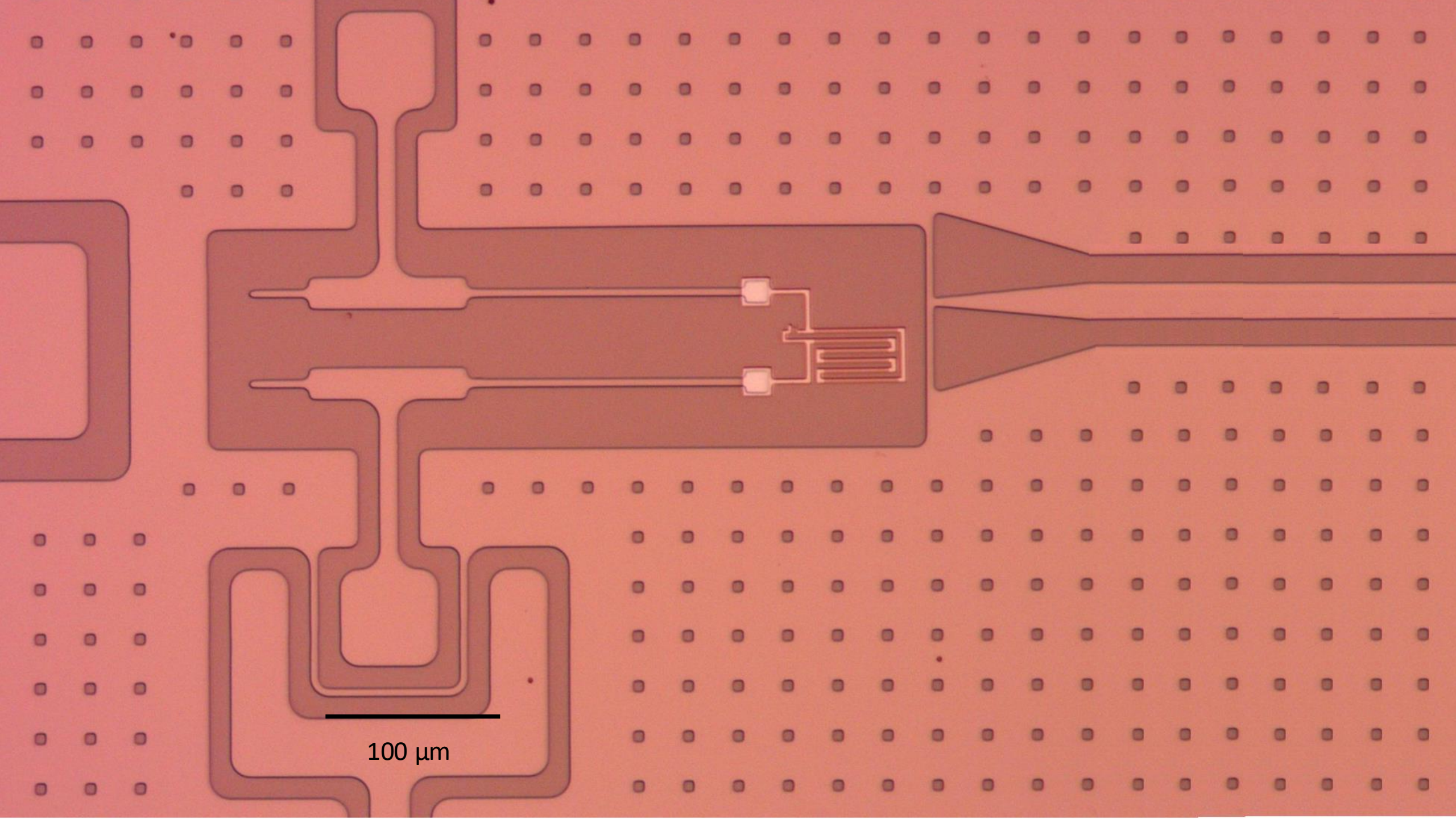
The localized states hybridize:

$$|e\rangle \rightarrow |L\rangle - |R\rangle$$

$$|g\rangle \rightarrow |L\rangle + |R\rangle$$

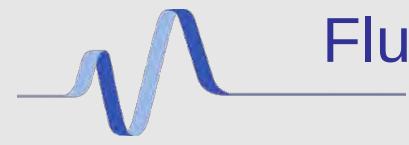
The tunnel splitting decreases exponentially with:  $E_J/E_C$



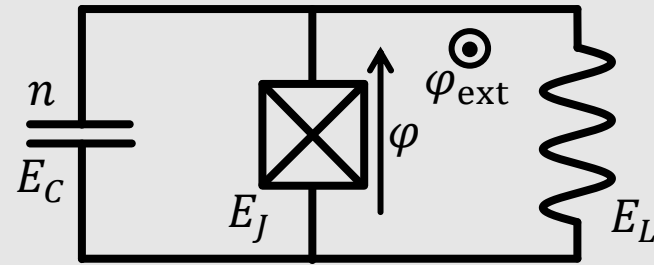


100 μm

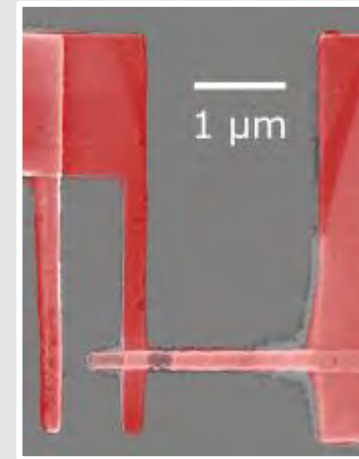




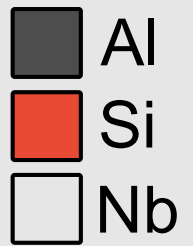
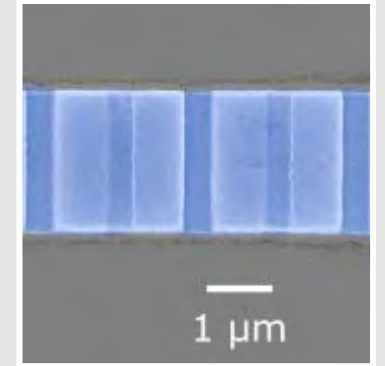
# Fluxonium design



Single small junction



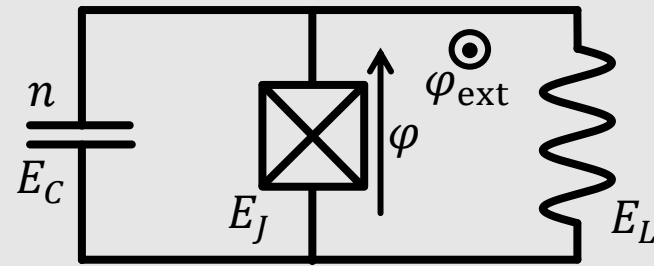
360 large junctions



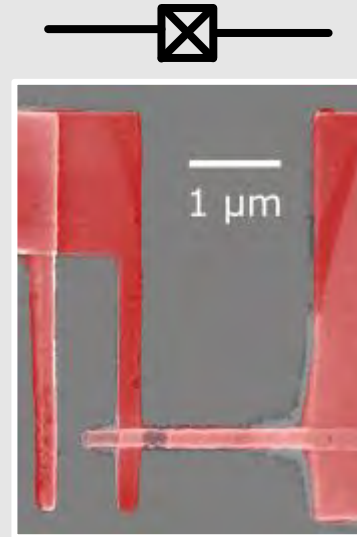
100  $\mu\text{m}$



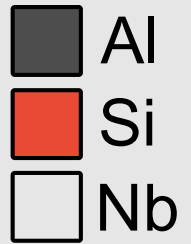
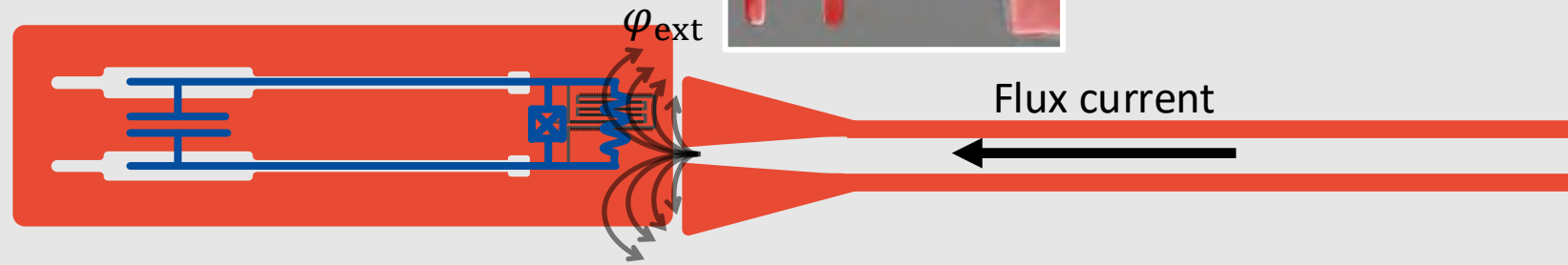
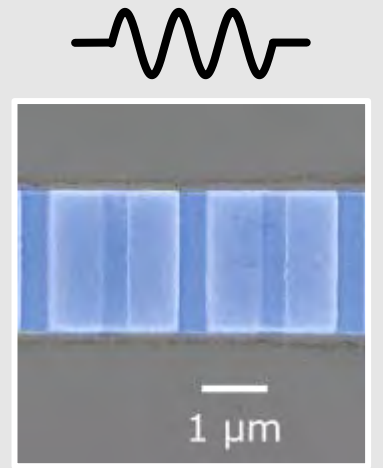
# Fluxonium design



Single small  
junction



360 large  
junctions



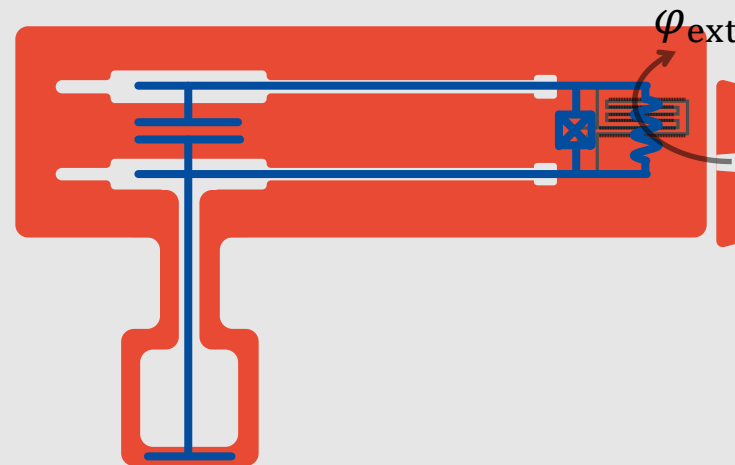
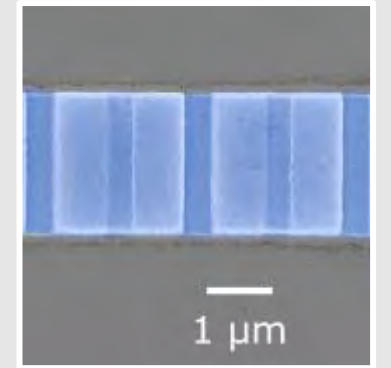
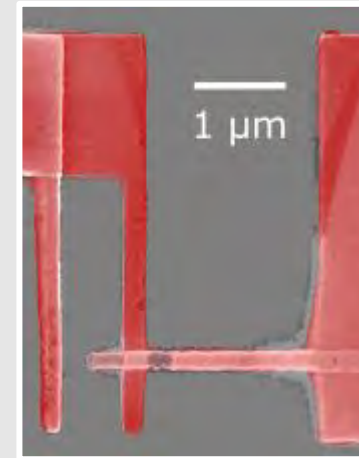
100  $\mu\text{m}$

# Fluxonium design

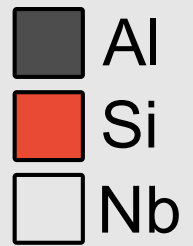
Single small  
junction



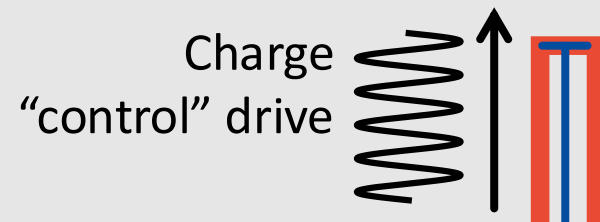
360 large  
junctions



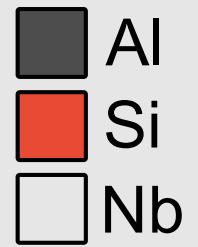
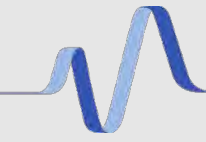
Flux current



100 μm

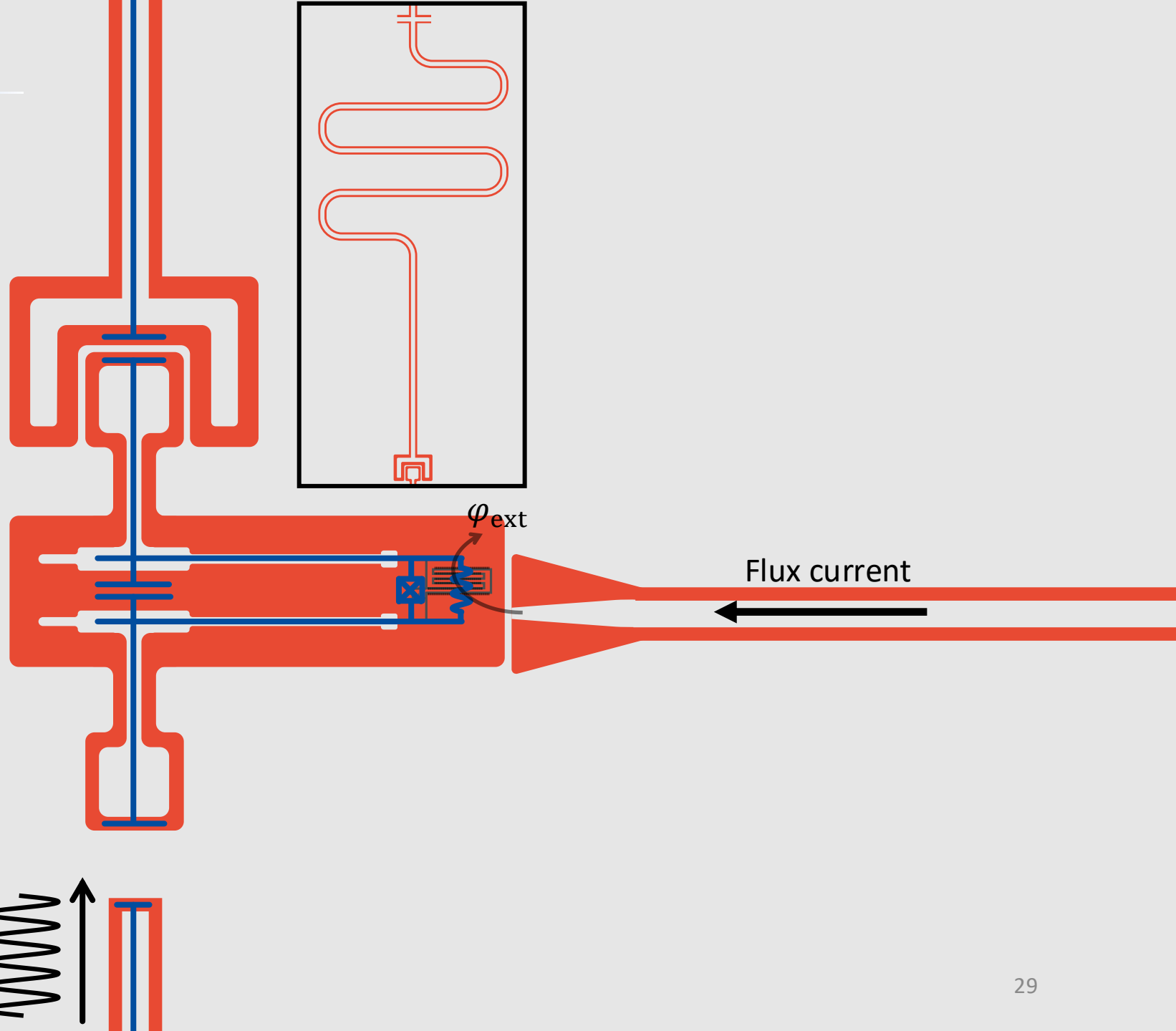
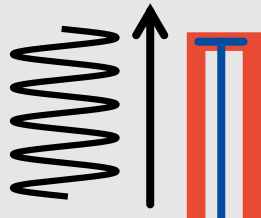


# Fluxonium design

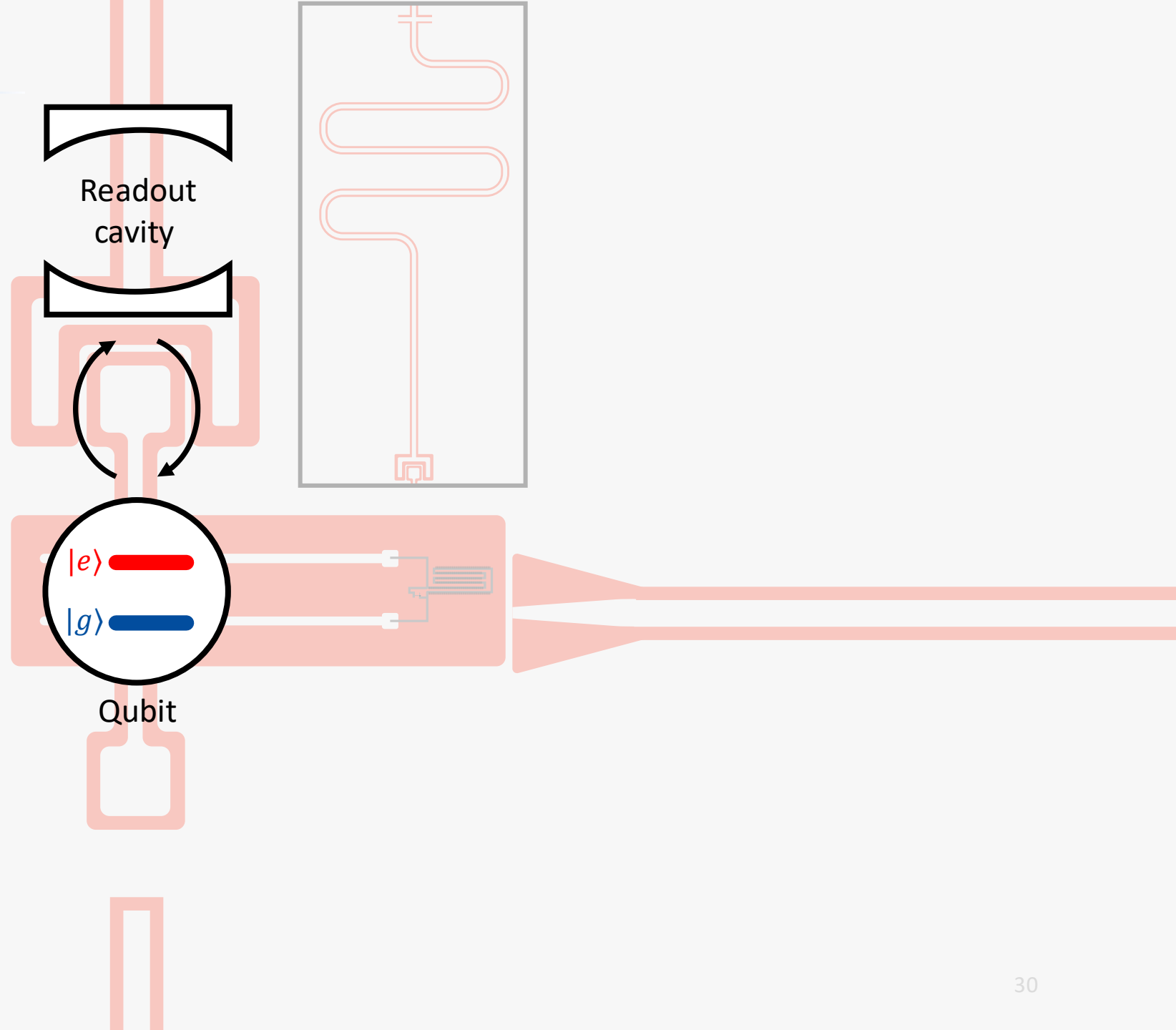


100  $\mu\text{m}$

Charge  
"control" drive

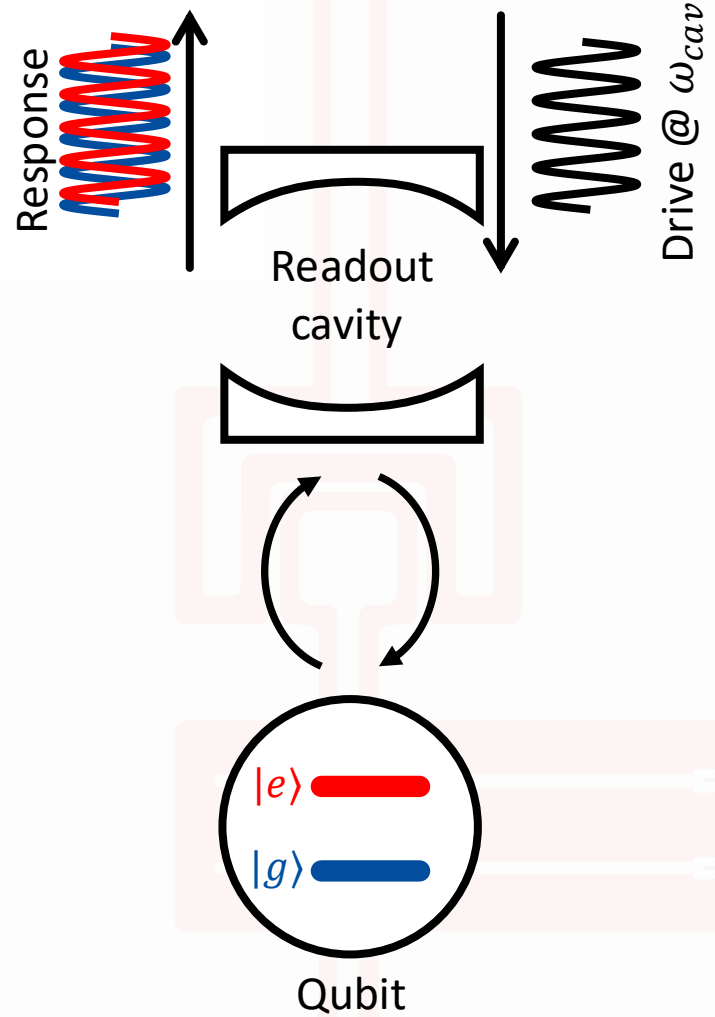


# Fluxonium readout





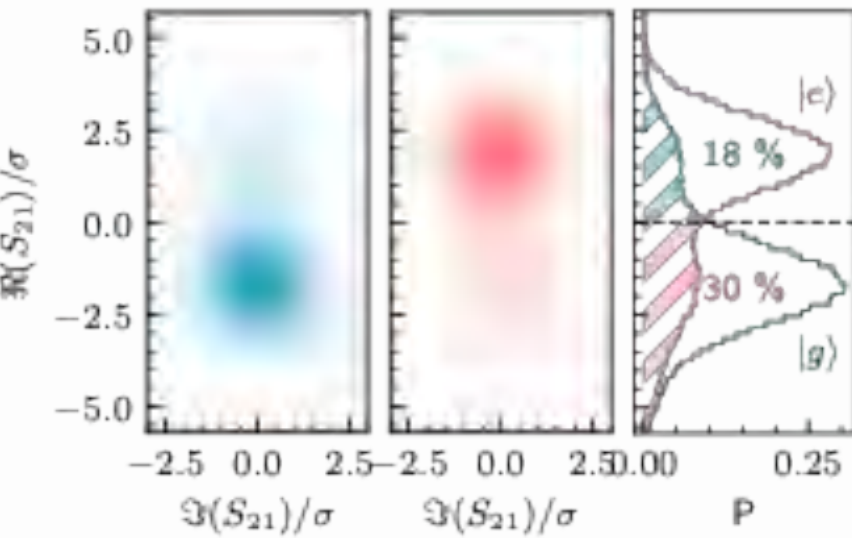
Readout response depends  
on the qubit state



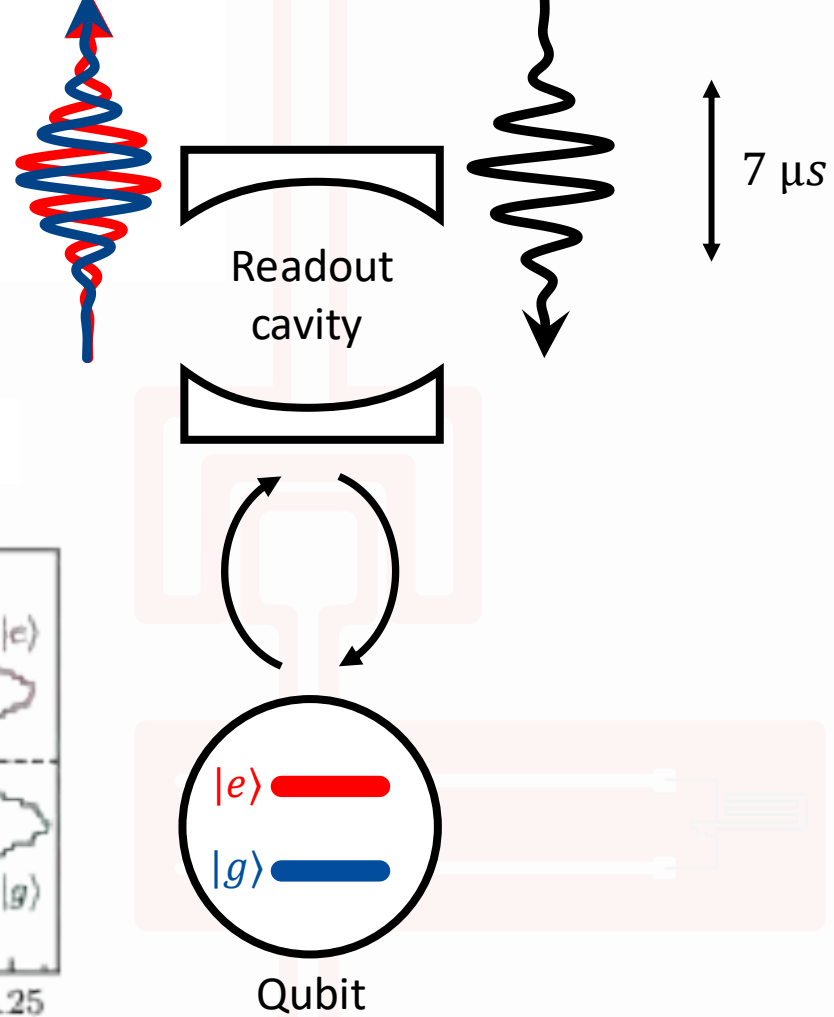
## Fluxonium readout

Readout response depends  
on the qubit state

Complex reflection coefficient



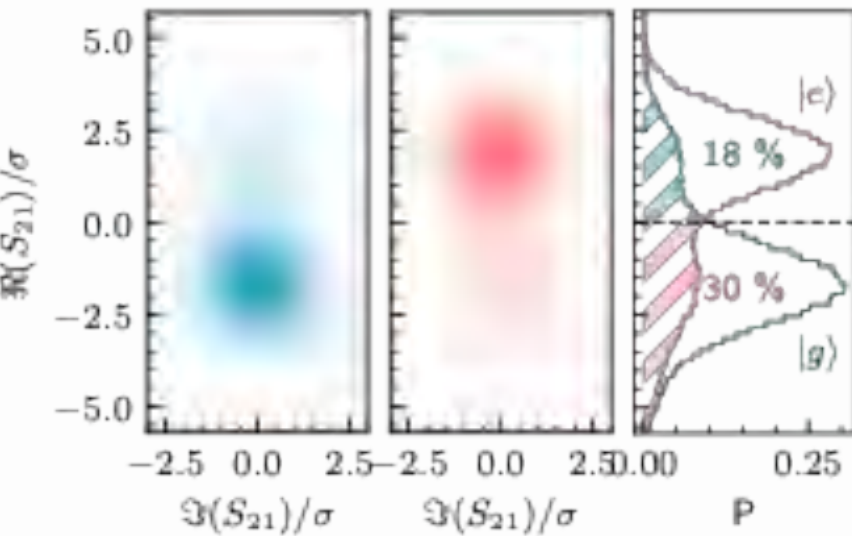
Single-shot readout  
with 77 % fidelity



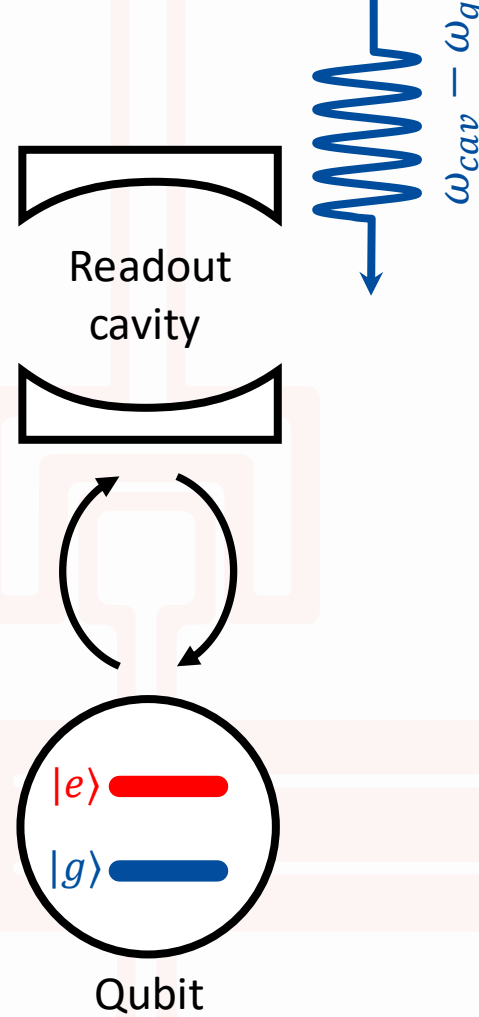
# Sideband preparation

Readout response depends  
on the qubit state

Complex reflection coefficient



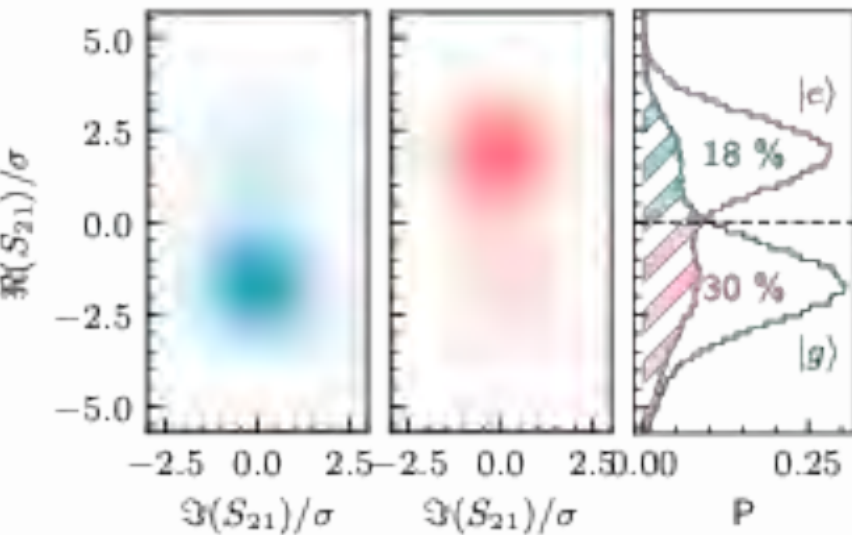
Single-shot readout  
with 77 % fidelity



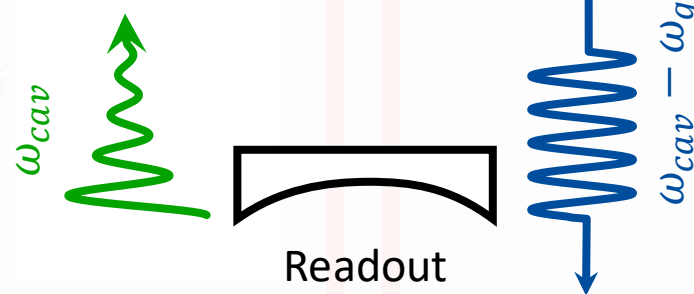
Sideband preparation

Readout response depends on the qubit state

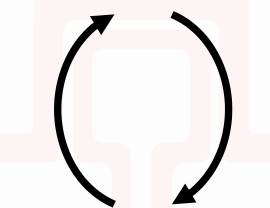
Complex reflection coefficient



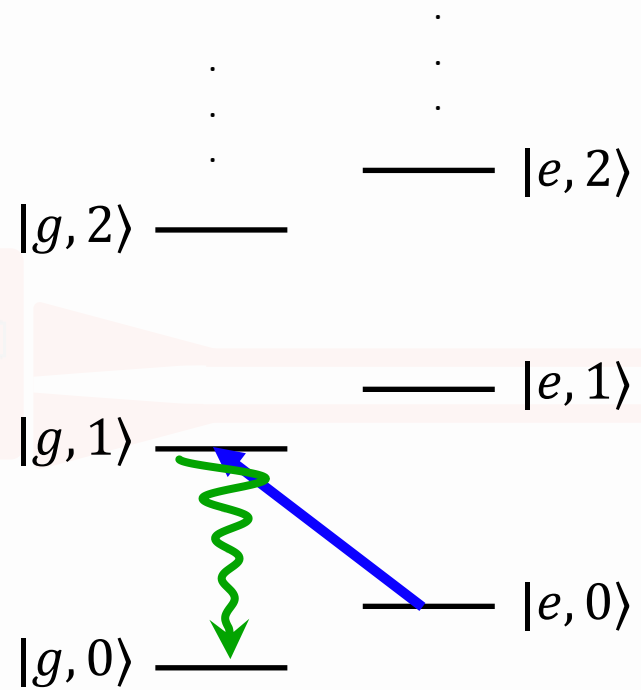
Single-shot readout with 77 % fidelity



Readout cavity



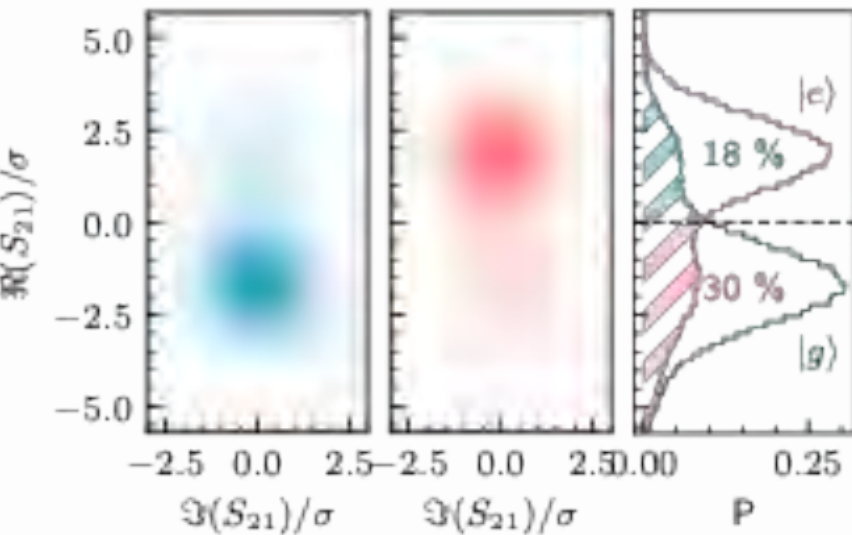
Qubit



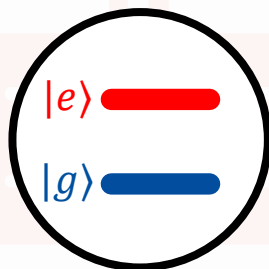
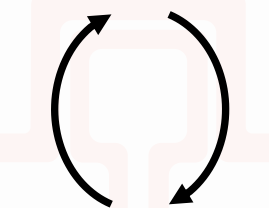
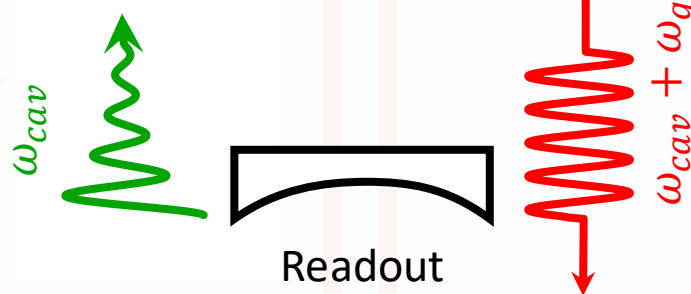
# Sideband prepa

Readout response depends  
on the qubit state

Complex reflection coefficient

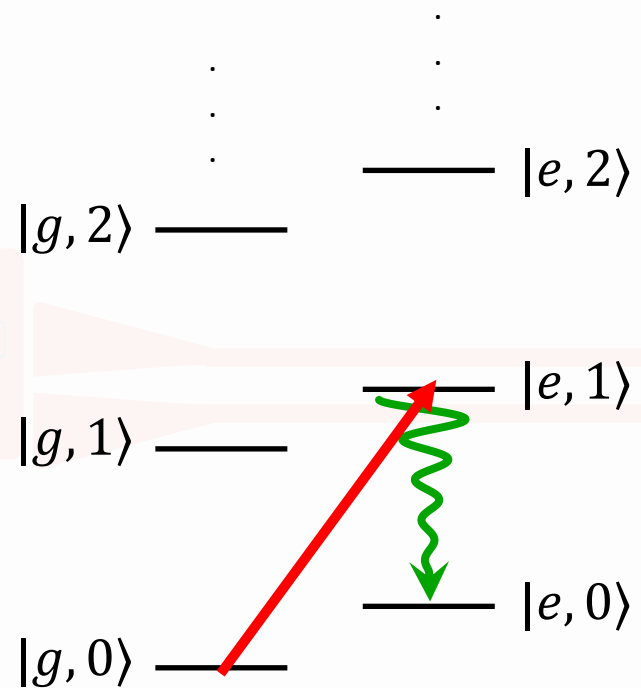


Single-shot readout  
with 77 % fidelity

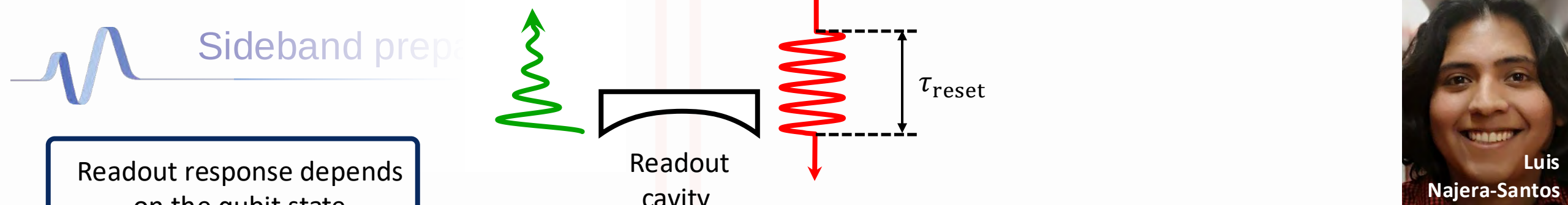


Qubit

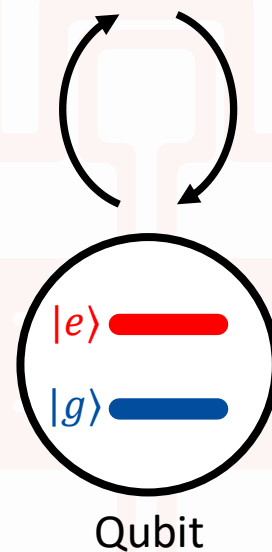
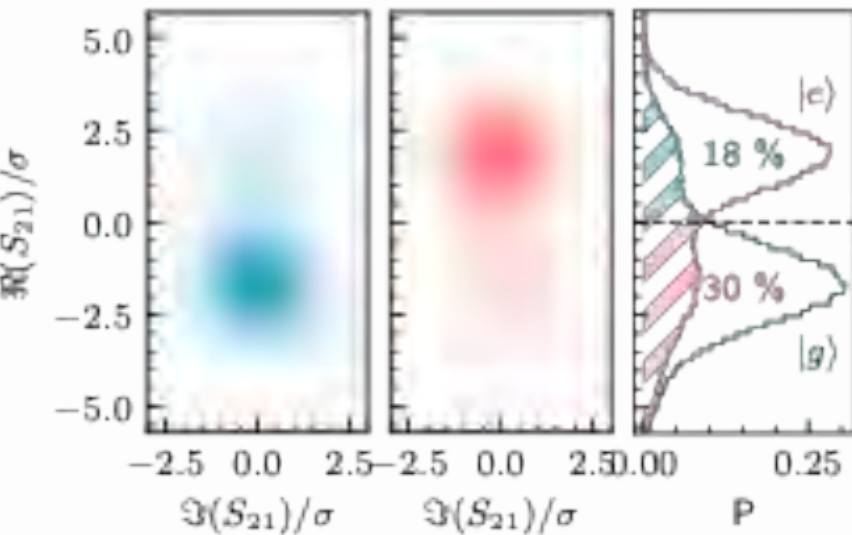
$E$



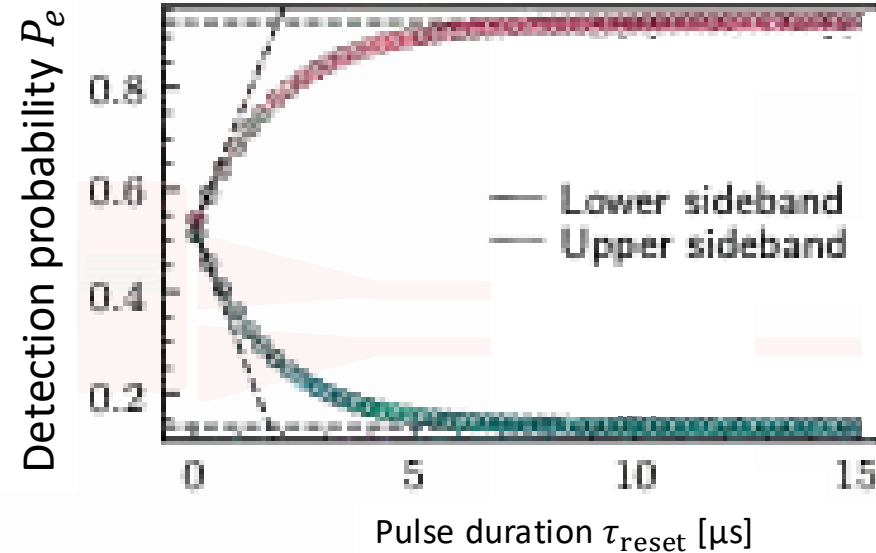




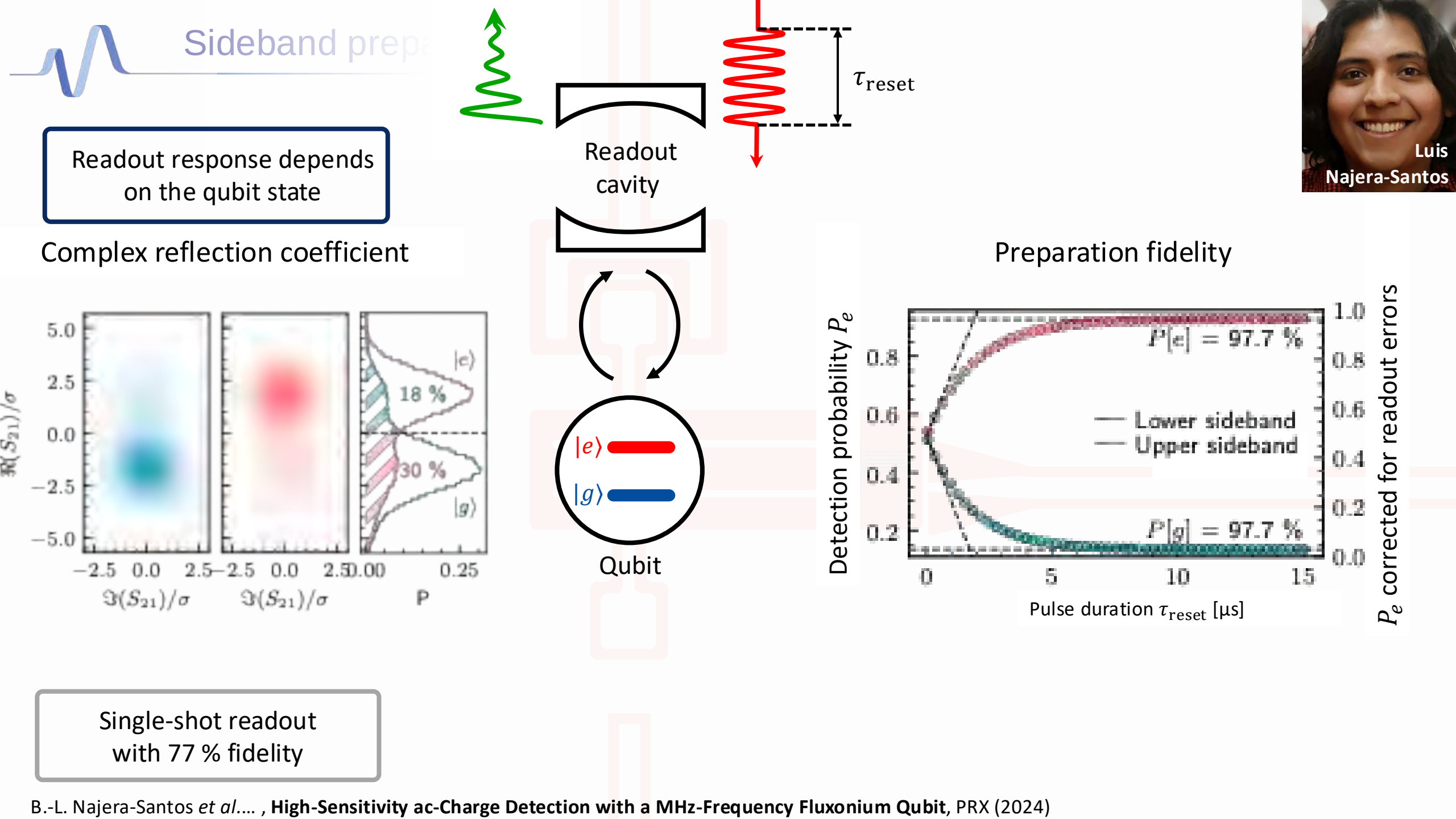
Complex reflection coefficient

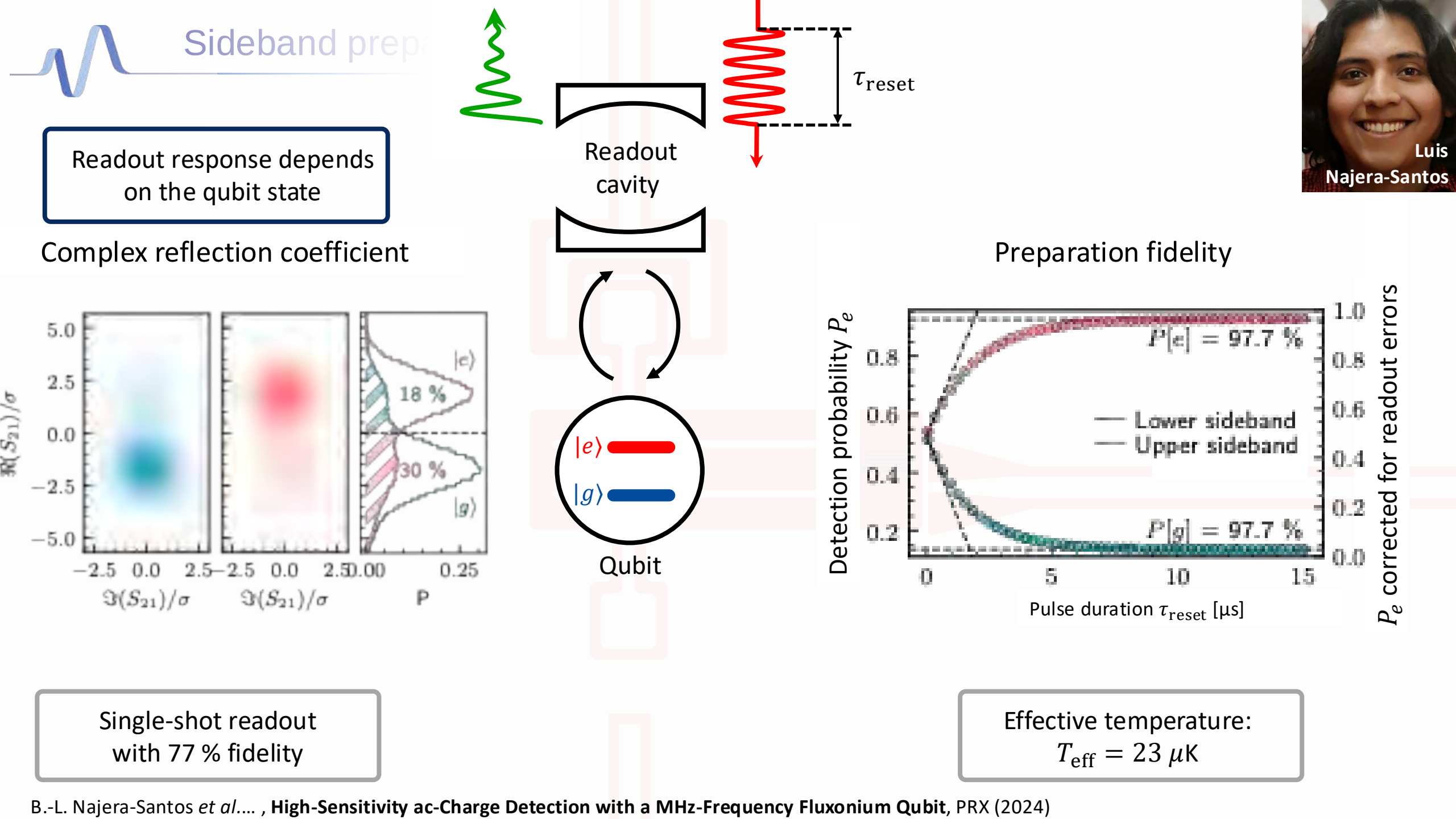


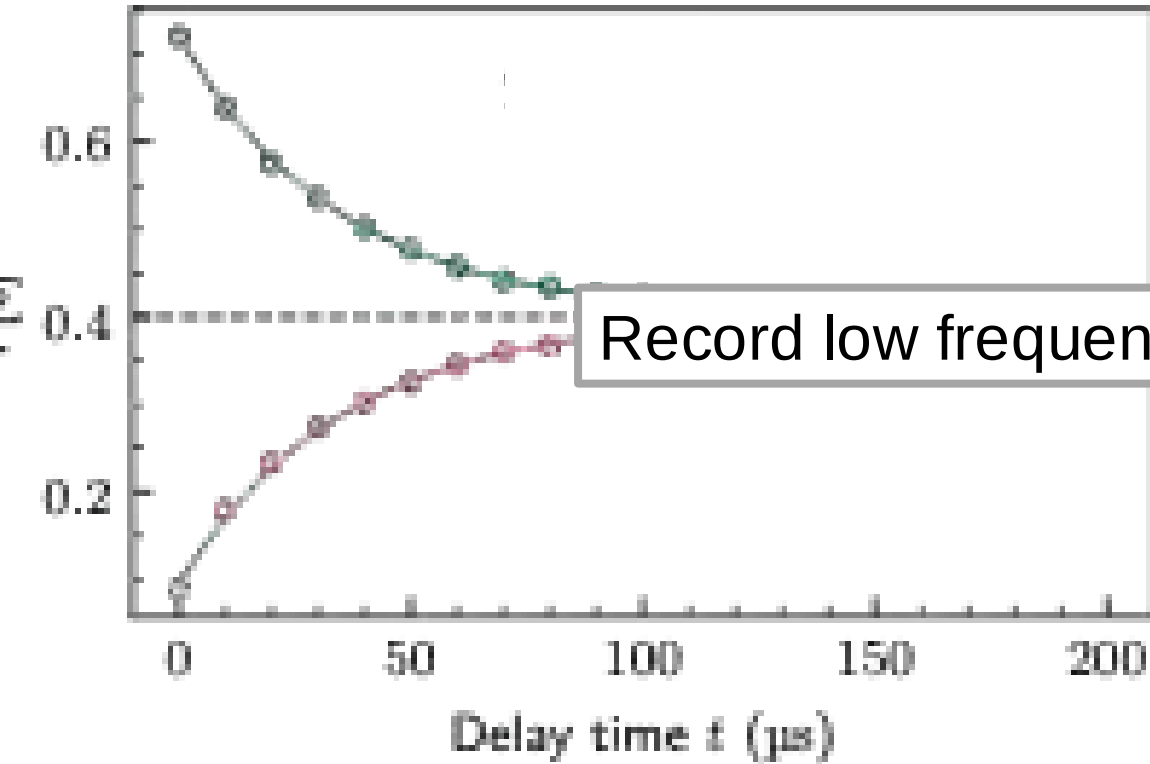
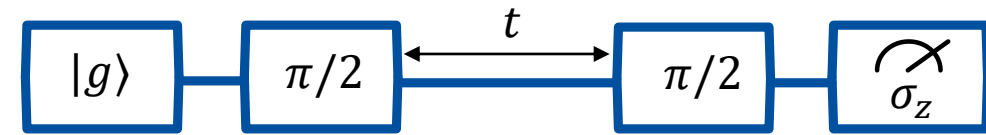
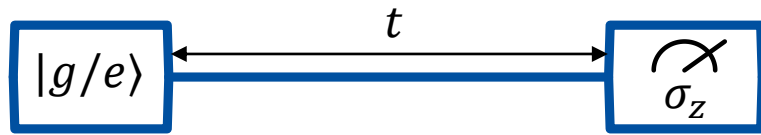
Preparation fidelity



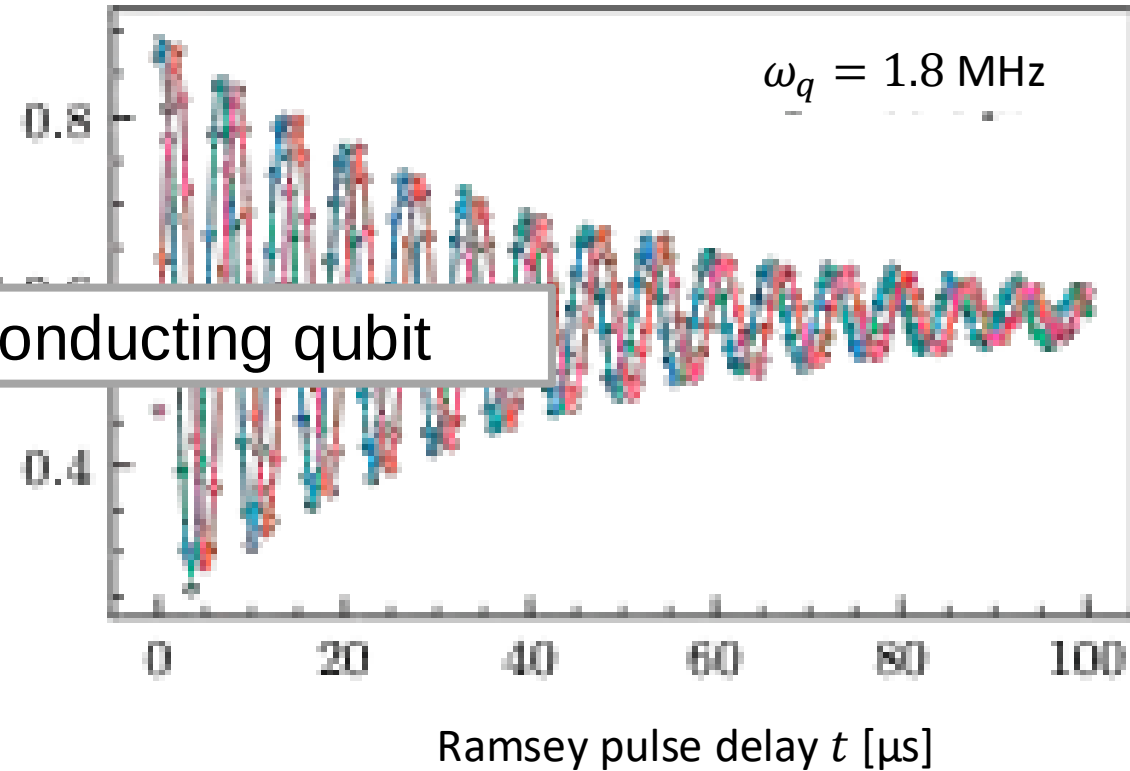
Single-shot readout with 77 % fidelity





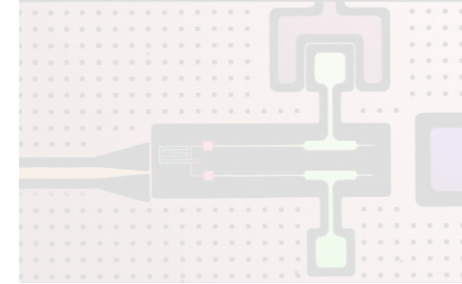


Energy relaxation time  
 $T_1 = 34 \mu\text{s}$



Ramsey decay time  
 $T_2^* = 38.6 \mu\text{s}$

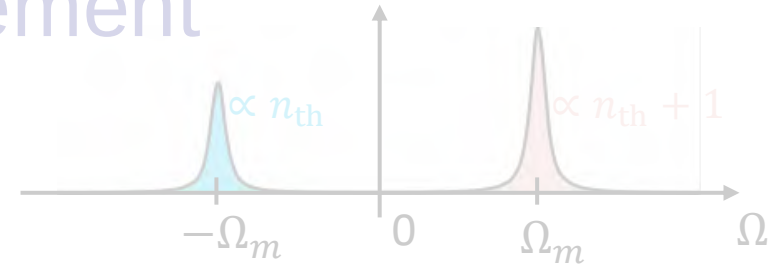
- Radio-frequency fluxonium qubits



- Hybrid qubit/mechanical system: MecaFlux



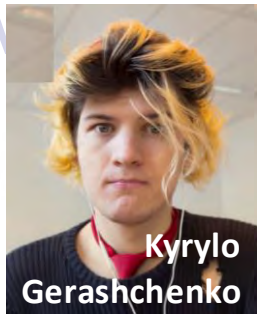
- Quantum asymmetry measurement



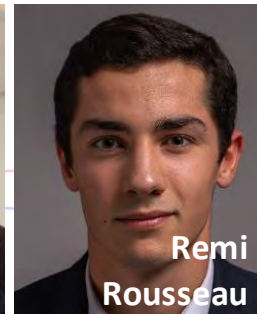
- People



Leo  
Balembois



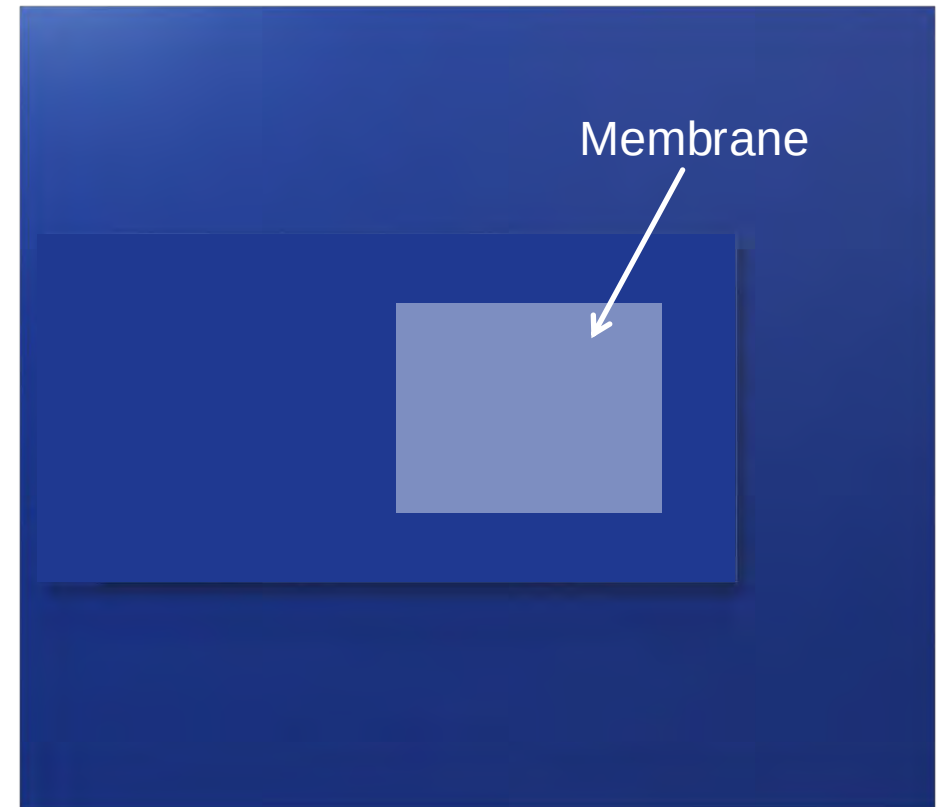
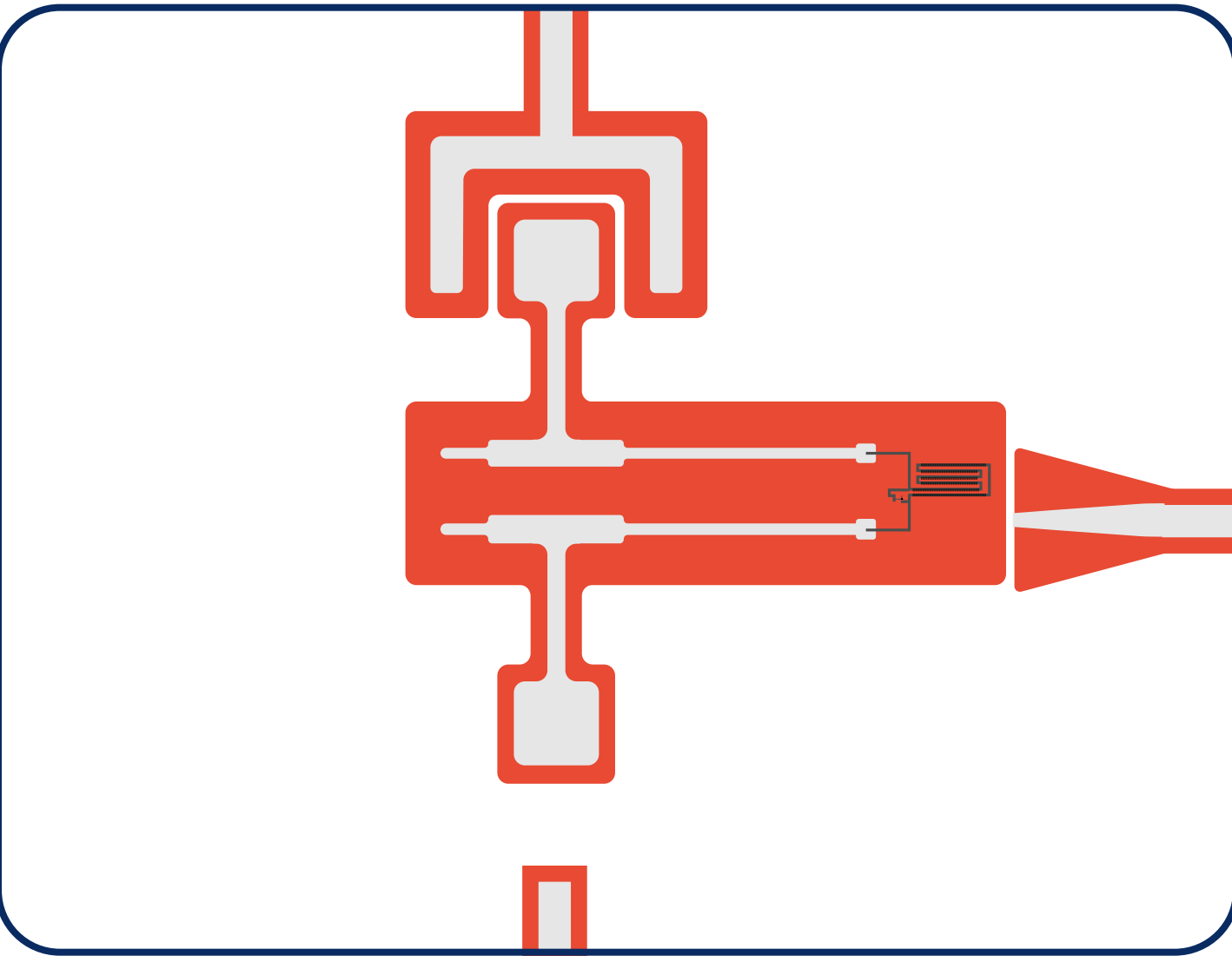
Kyrylo  
Gerashchenko



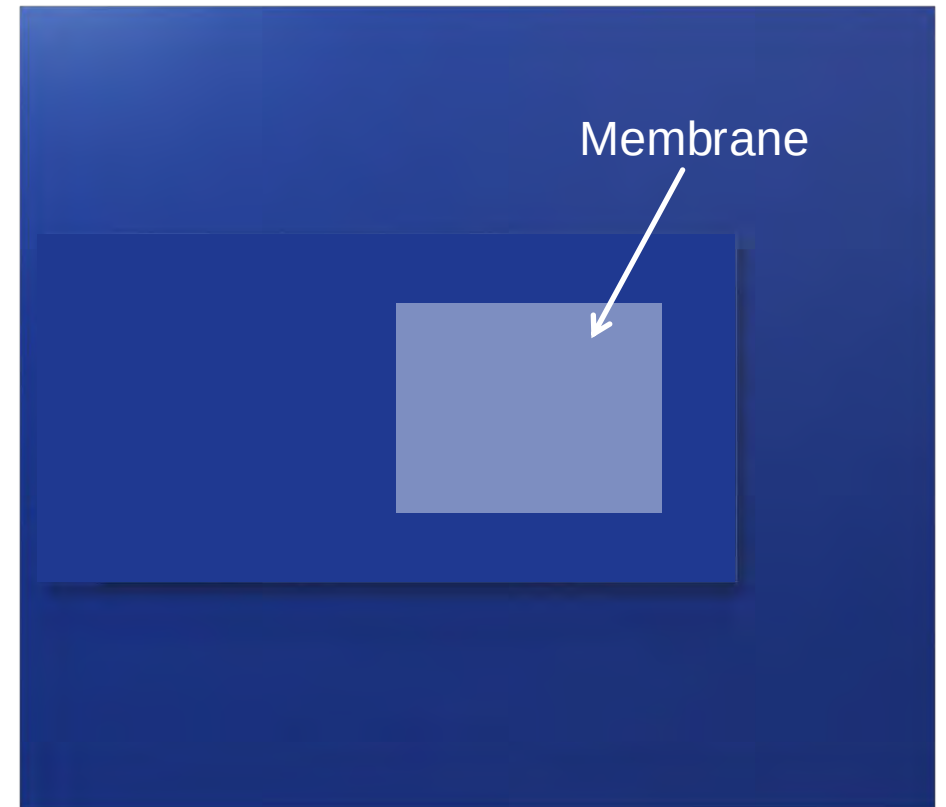
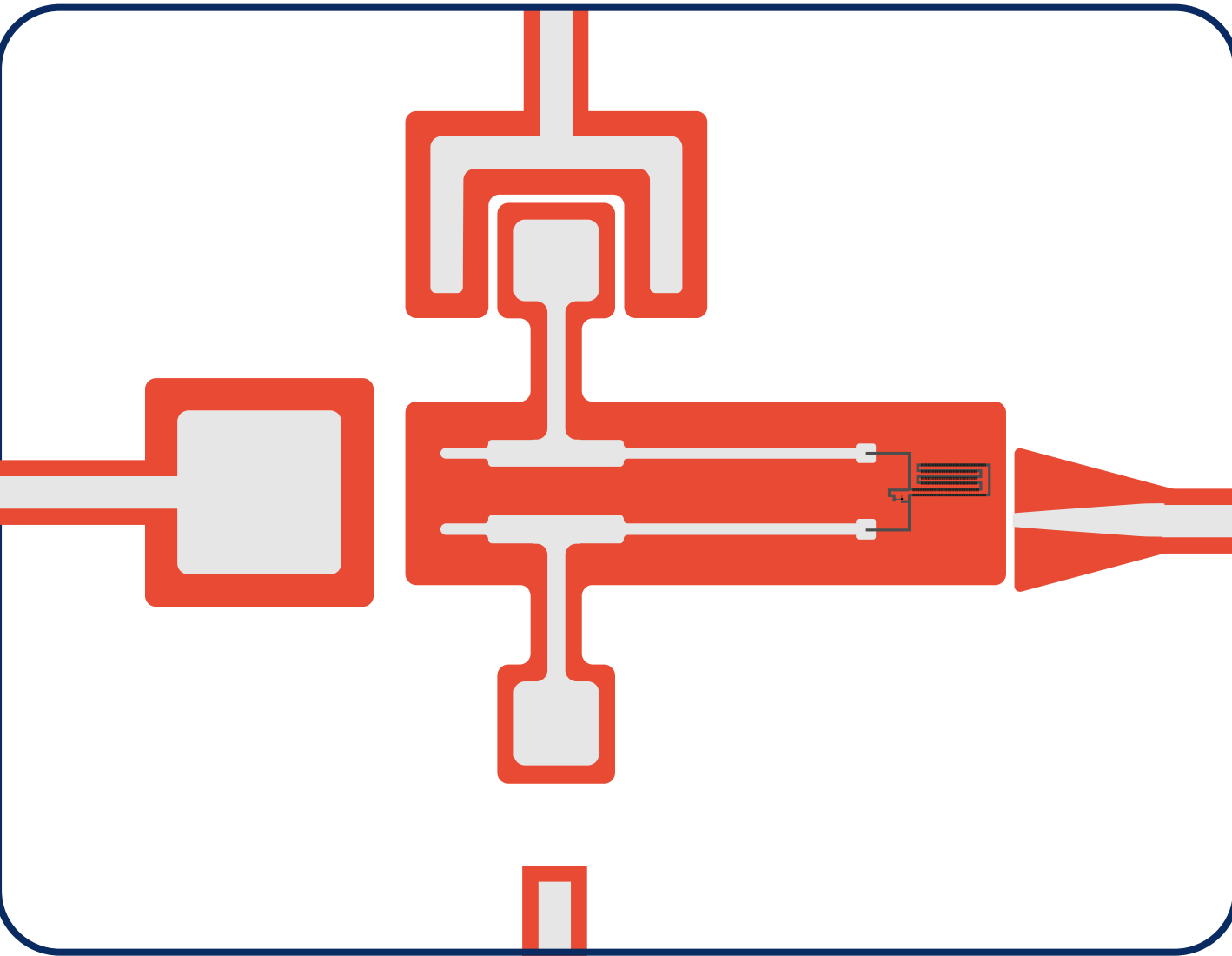
Remi  
Rousseau



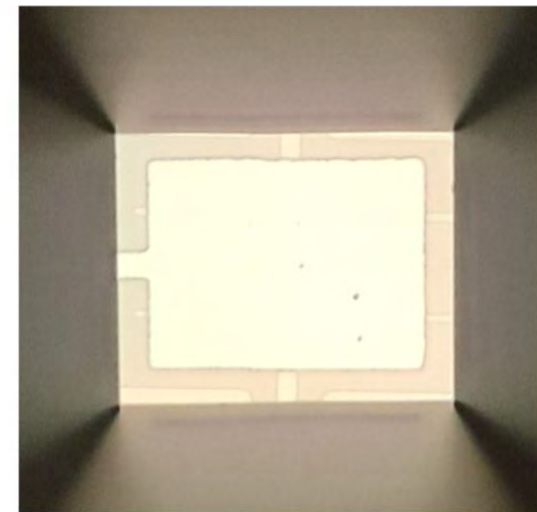
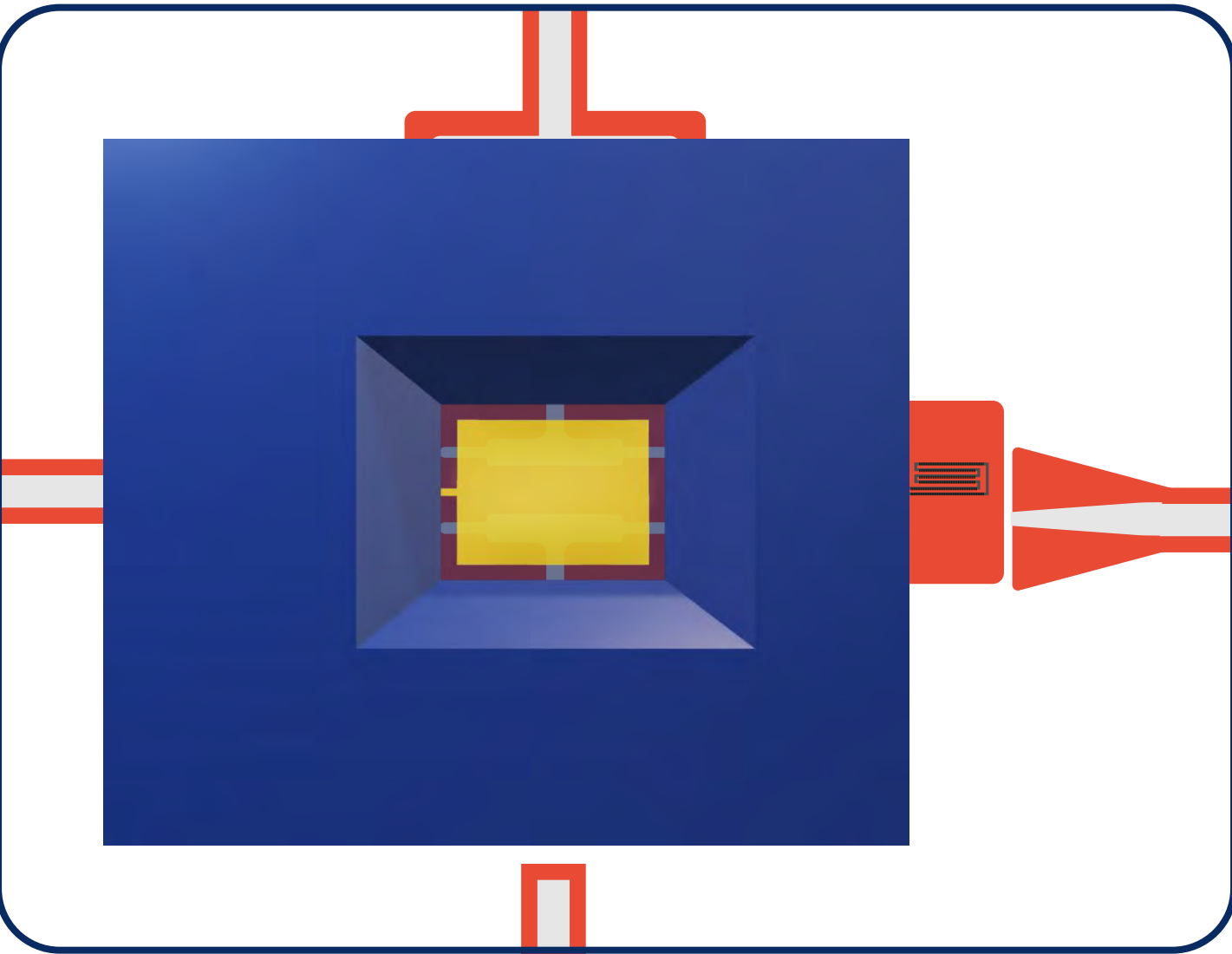
# Qubit-membrane coupling

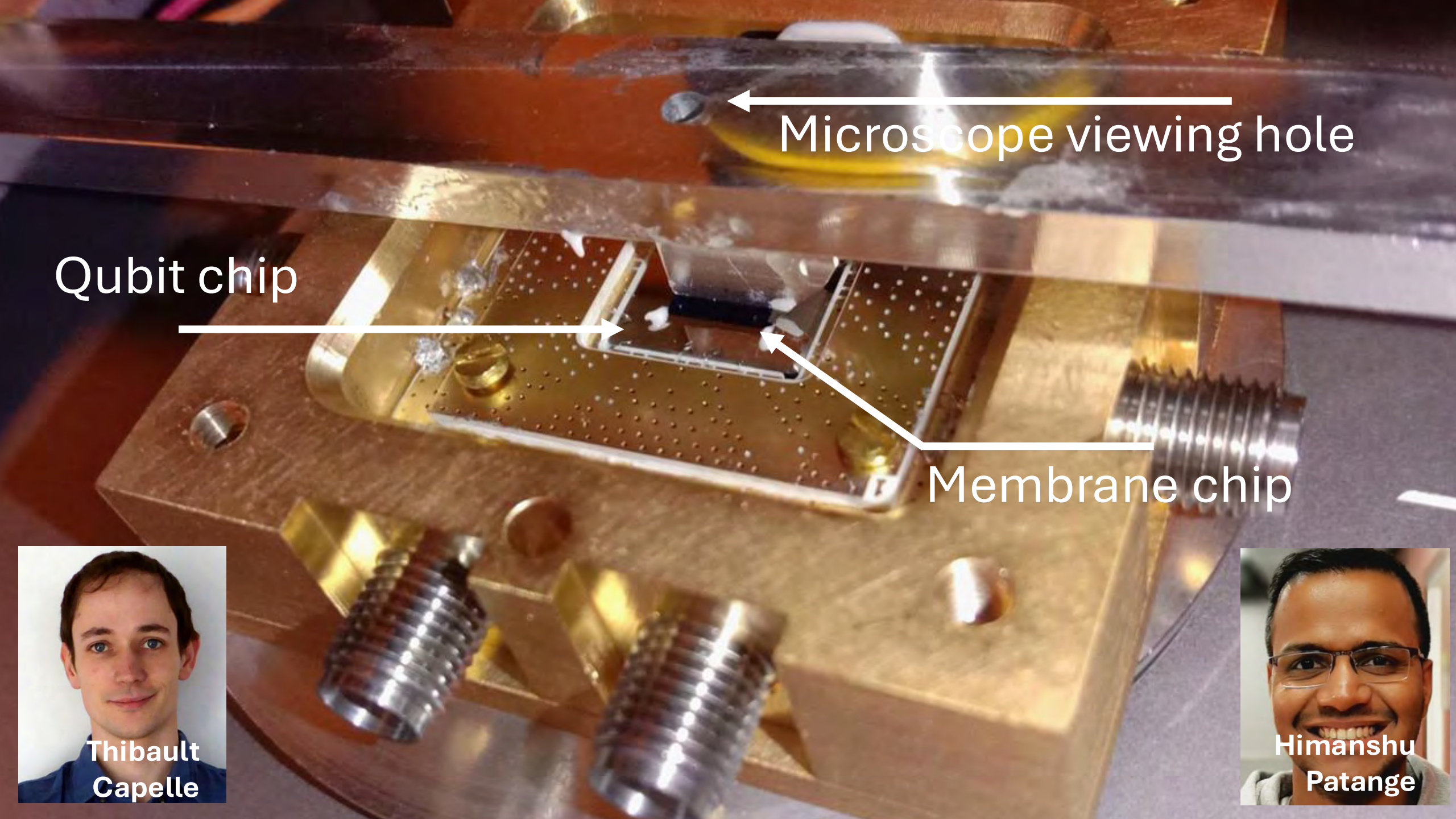


# Qubit-membrane coupling



# Qubit-membrane coupling





Microscope viewing hole

Qubit chip

Membrane chip



Thibault  
Capelle



Himanshu  
Patange

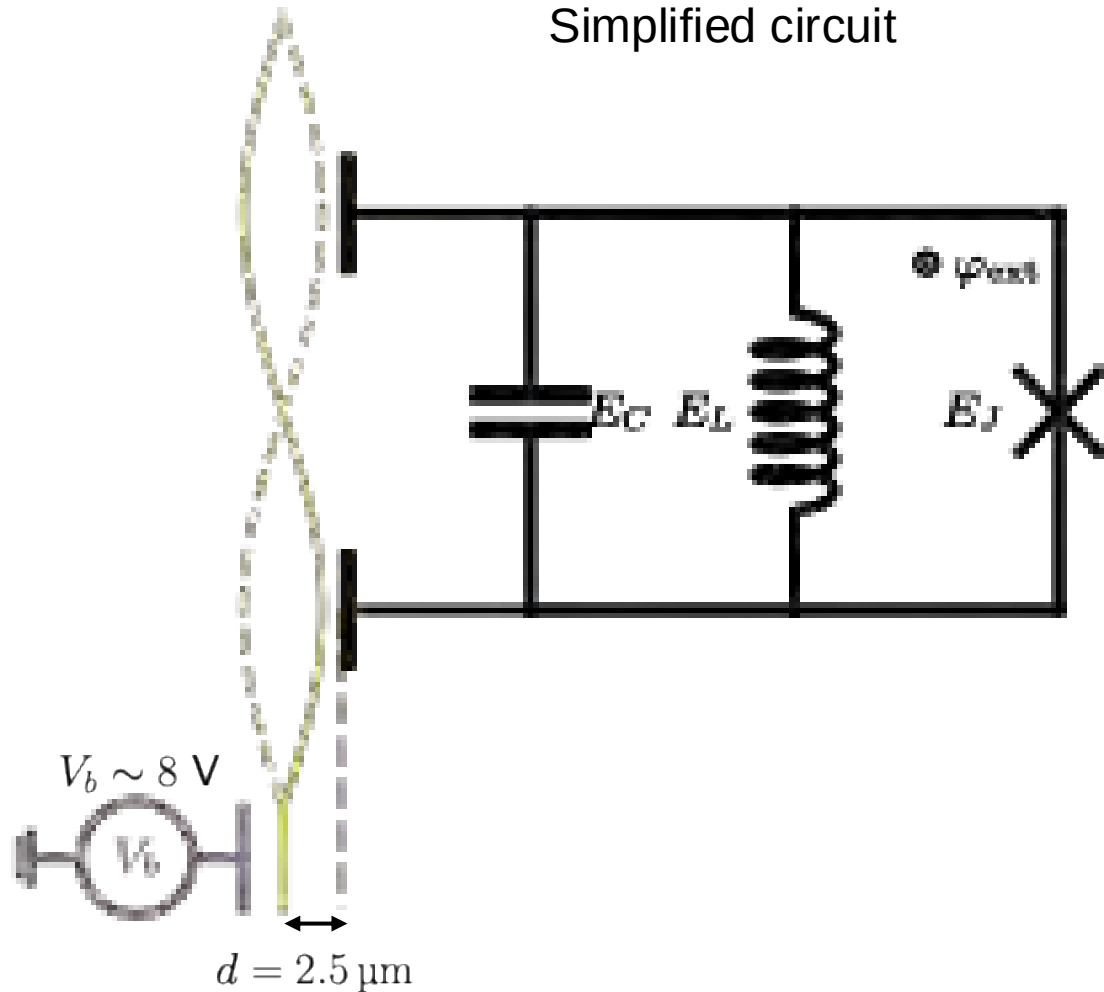






# Coupling description

Simplified circuit



System Hamiltonian

$$H = \frac{1}{2} E_L \varphi^2 - E_J \cos(\varphi - \varphi_{\text{ext}}) + 4E_C (\mathbf{n} - n_g(\mathbf{x}))^2$$

Position-dependent offset charge :

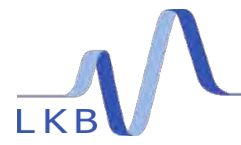
$$n_g(\mathbf{x}) = n_{\text{zpf}} \mathbf{x} \propto V_b \mathbf{x} / d^2$$

Linear coupling :

$$H_{\text{int}} \propto \mathbf{n} \cdot \mathbf{x}$$

Jaynes-Cummings interaction :

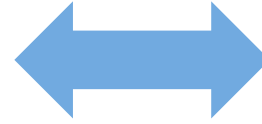
$$H_{\text{int}} / \hbar = i \frac{\Omega}{2} (a^\dagger \sigma_- - a \sigma_+)$$



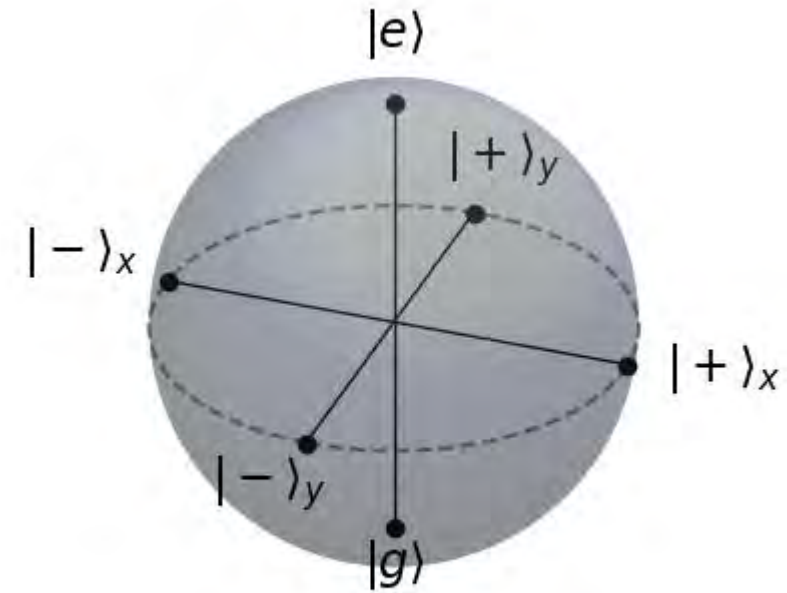
# Membrane spectroscopy (via Ramsey)



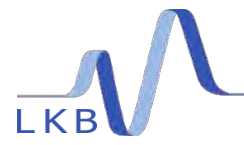
$$\Omega_m/2\pi = 4.4 \text{ MHz}$$



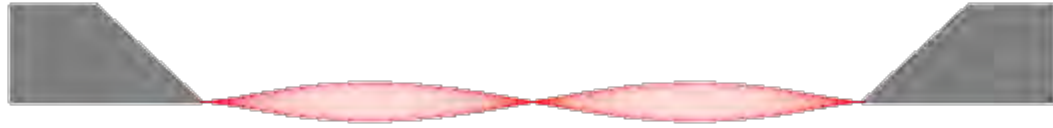
Dispersive interaction



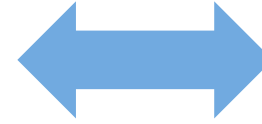
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$



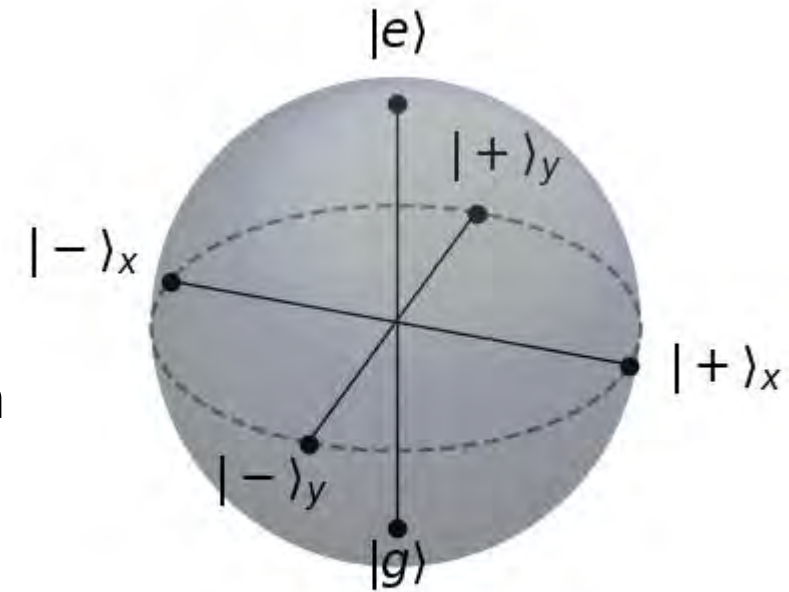
# Membrane spectroscopy (via Ramsey)



$$\Omega_m/2\pi = 4.4 \text{ MHz}$$

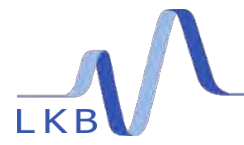


Dispersive interaction

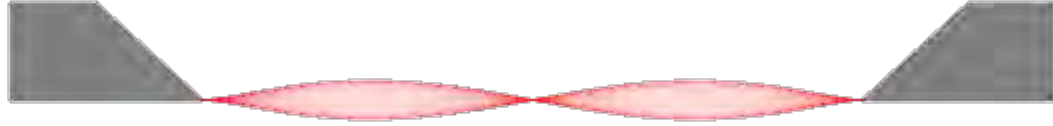


$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$

1. Continuously drive the membrane in a coherent state of mean phonon number  $\bar{n}$



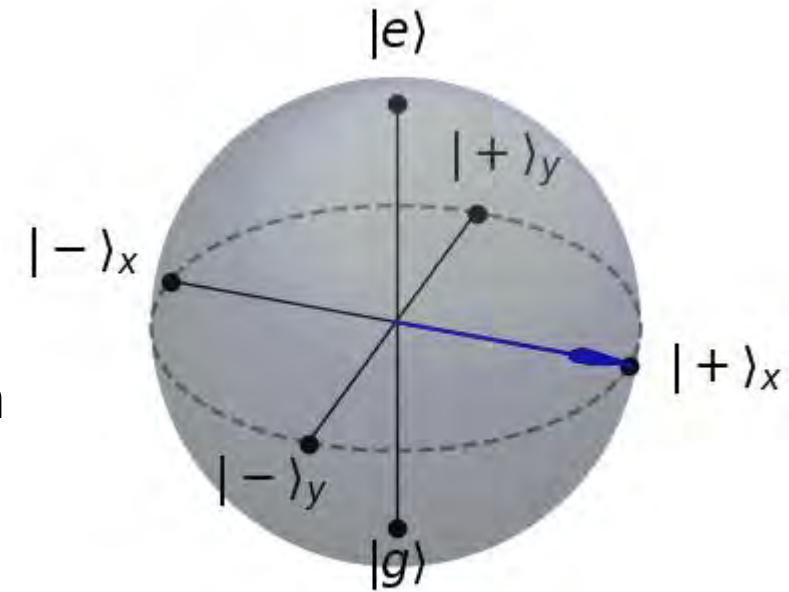
# Membrane spectroscopy (via Ramsey)



$$\Omega_m/2\pi = 4.4 \text{ MHz}$$



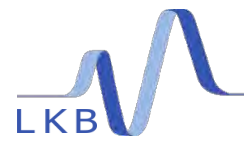
Dispersive interaction



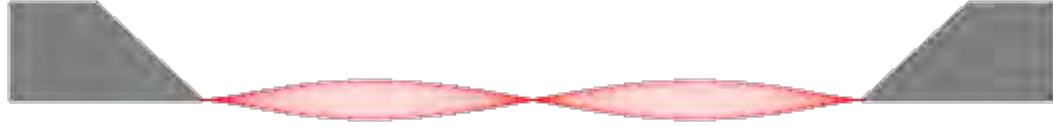
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$

1. Continuously drive the membrane in a coherent state of mean phonon number  $\bar{n}$

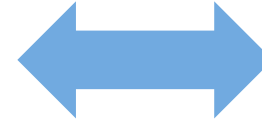
2. Detect the mechanically induced AC-Stark shift with a Ramsey sequence



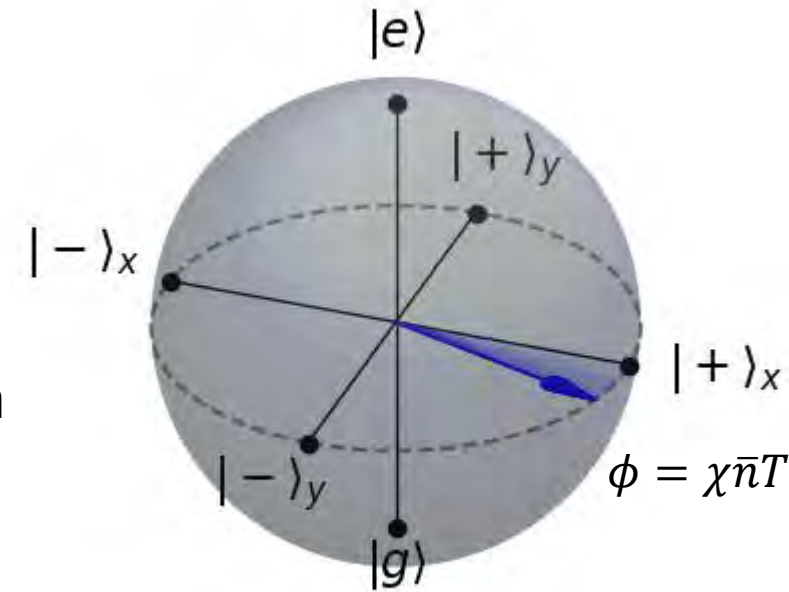
# Membrane spectroscopy (via Ramsey)



$$\Omega_m/2\pi = 4.4 \text{ MHz}$$



Dispersive interaction

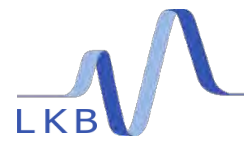


$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$

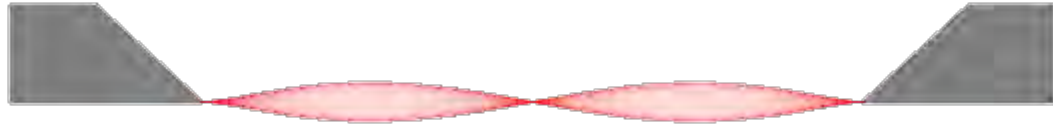
1. Continuously drive the membrane in a coherent state of mean phonon number  $\bar{n}$

2. Detect the mechanically induced AC-Stark shift with a Ramsey sequence

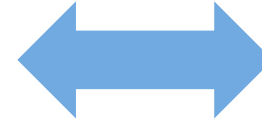




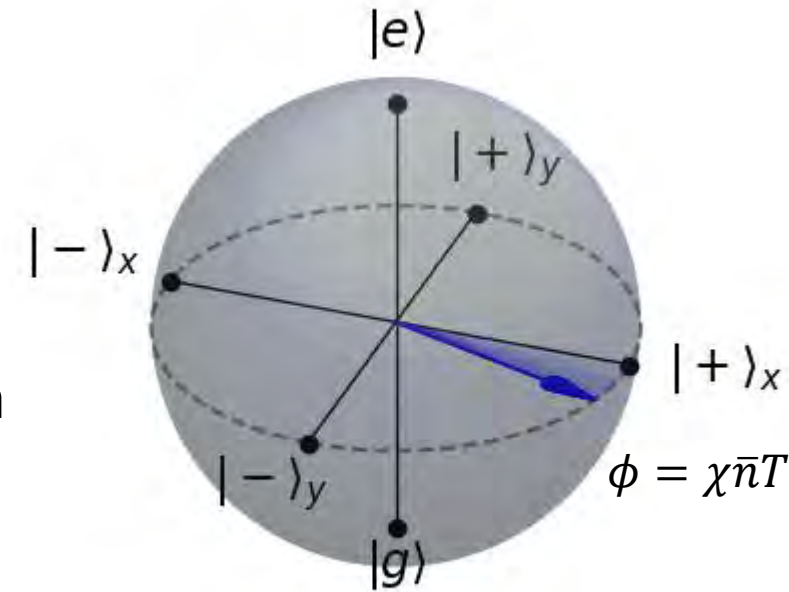
# Membrane spectroscopy (via Ramsey)



$$\Omega_m/2\pi = 4.4 \text{ MHz}$$



Dispersive interaction



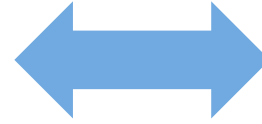
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$

1. Continuously drive the membrane in a coherent state of mean phonon number  $\bar{n}$
3. Repeat as a function of drive frequency

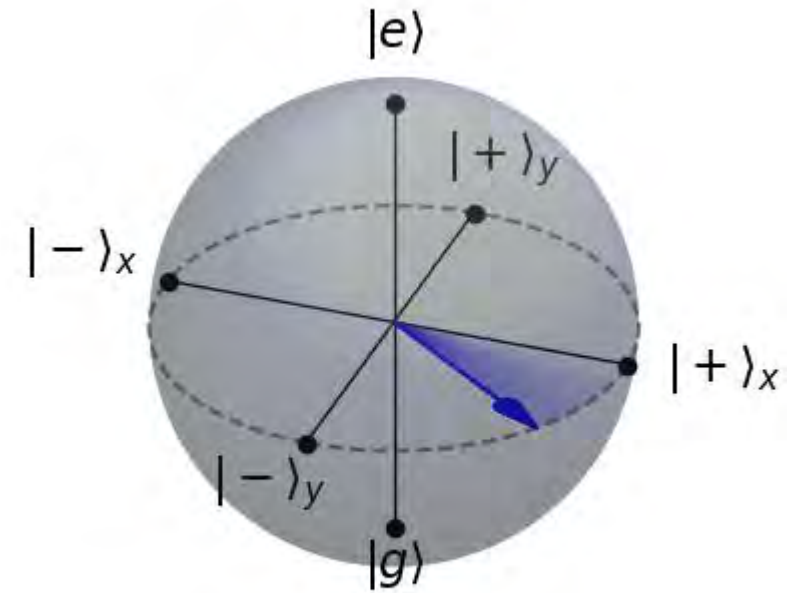
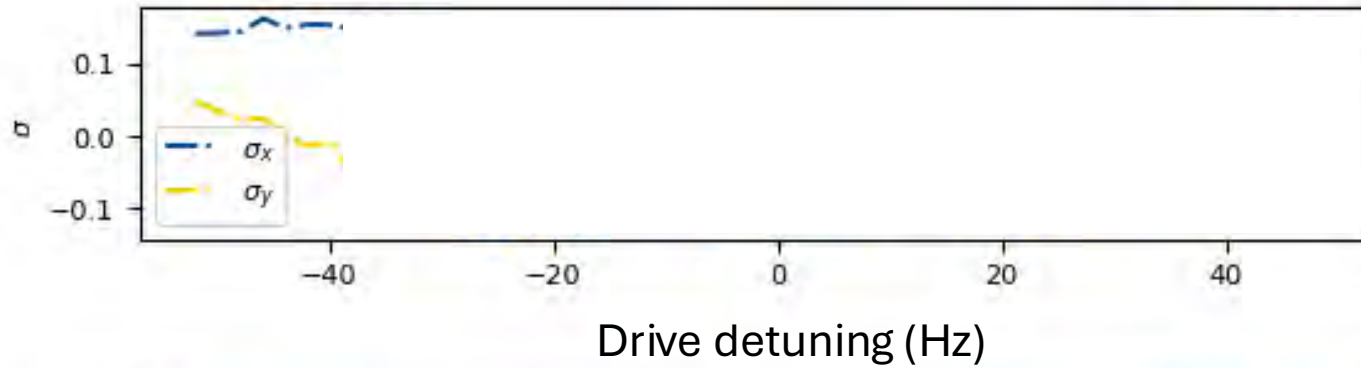
2. Detect the mechanically induced AC-Stark shift with a Ramsey sequence



# Membrane spectroscopy (via Ramsey)



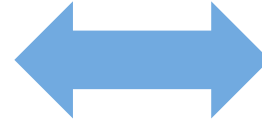
Dispersive interaction



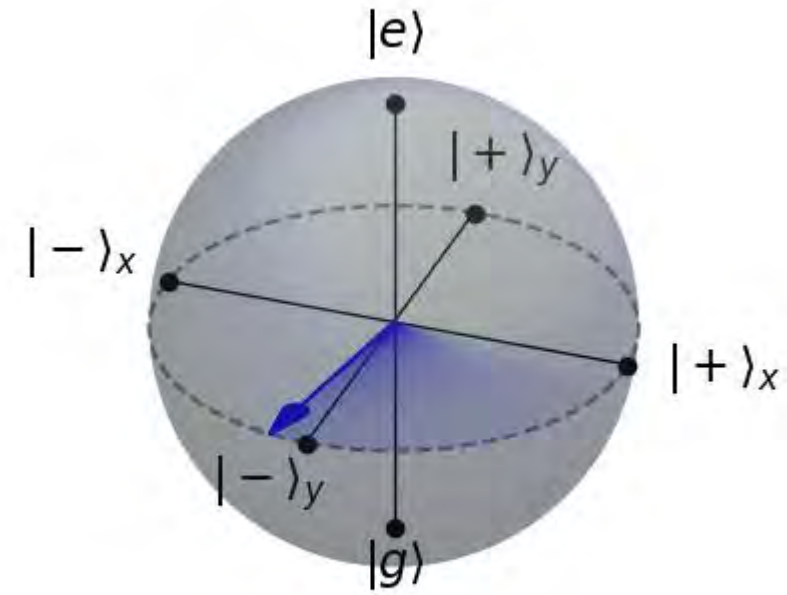
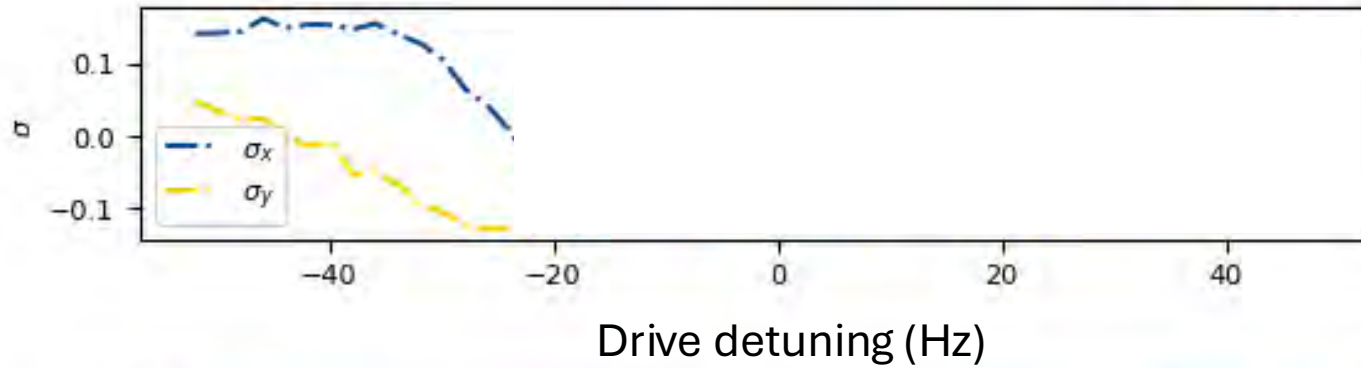
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$



# Membrane spectroscopy (via Ramsey)



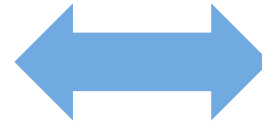
Dispersive interaction



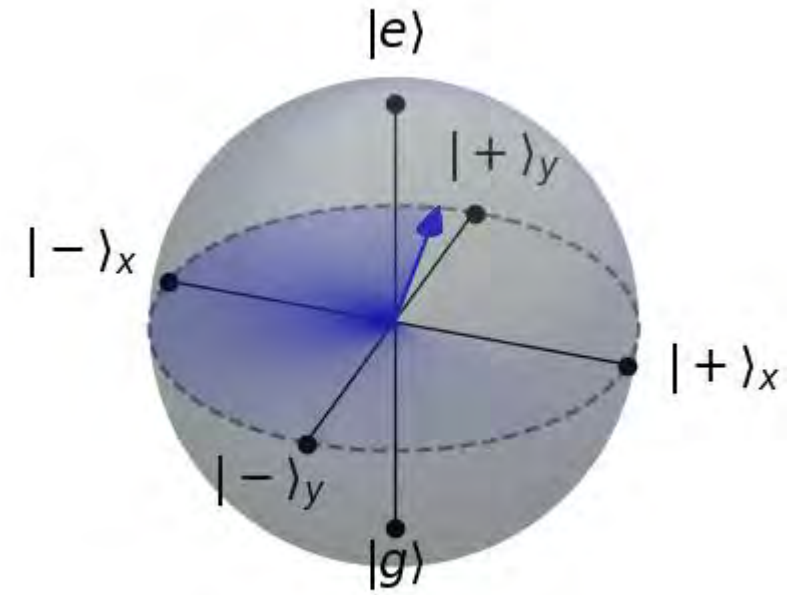
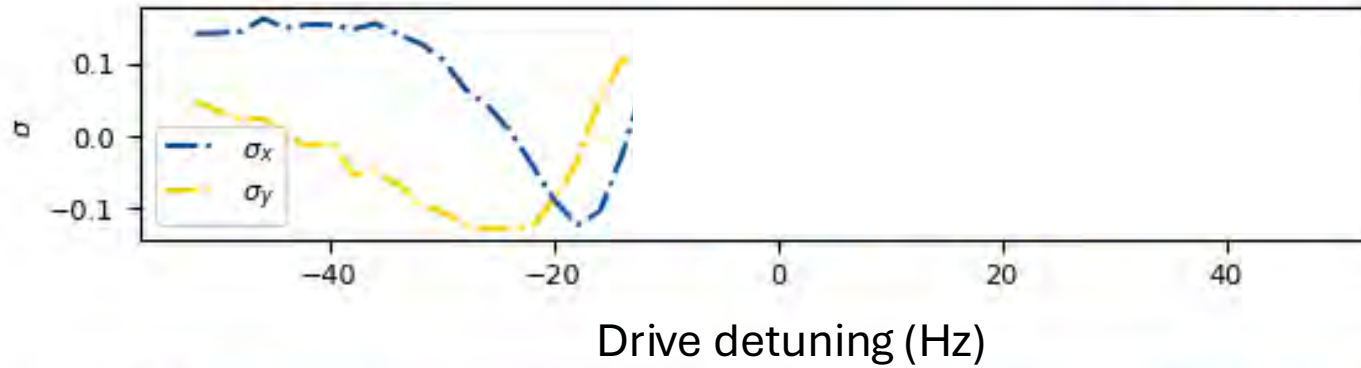
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$



# Membrane spectroscopy (via Ramsey)



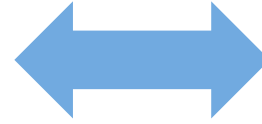
Dispersive interaction



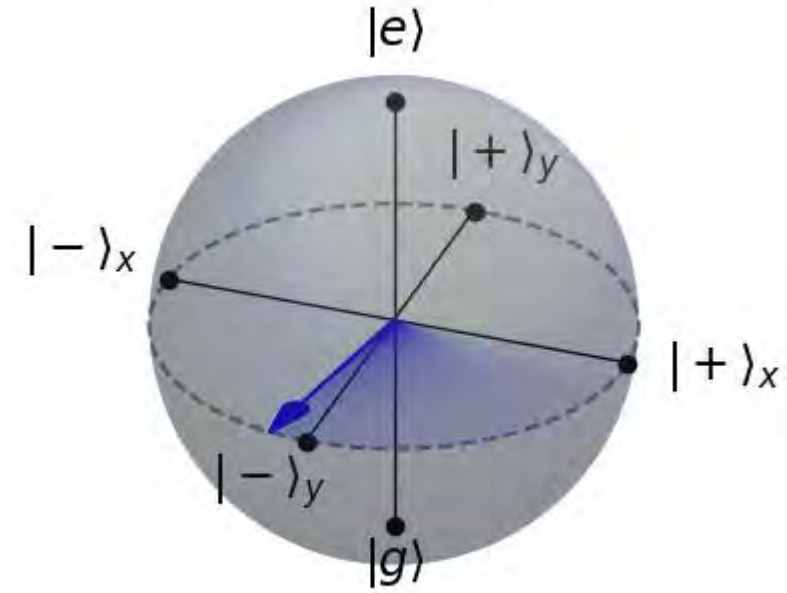
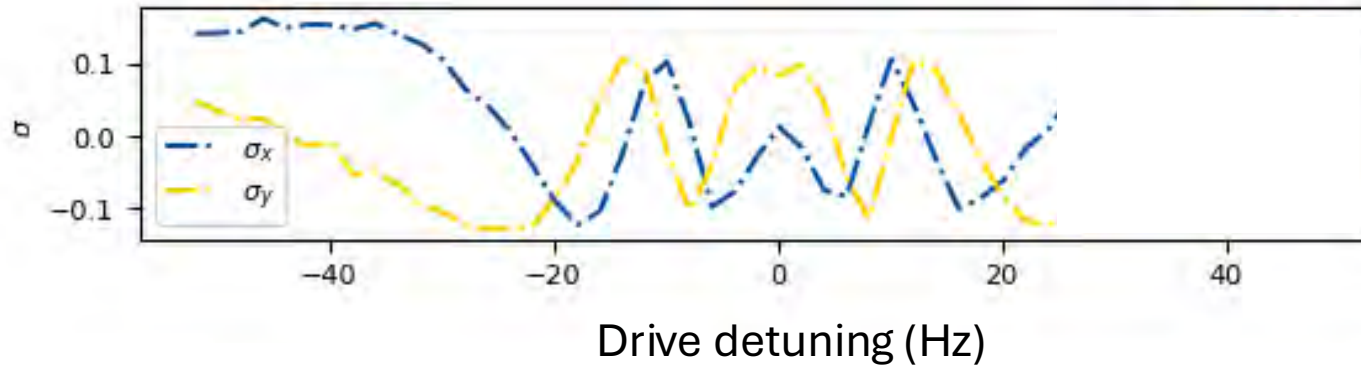
$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$



# Membrane spectroscopy (via Ramsey)

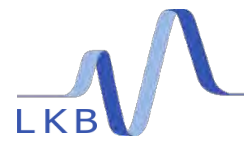


Dispersive interaction

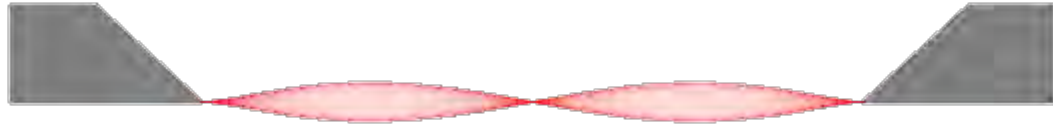


$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$





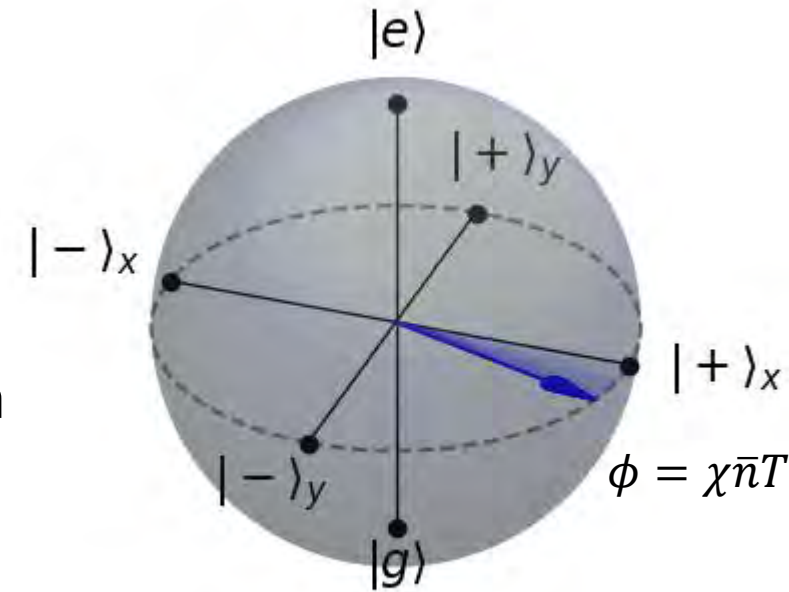
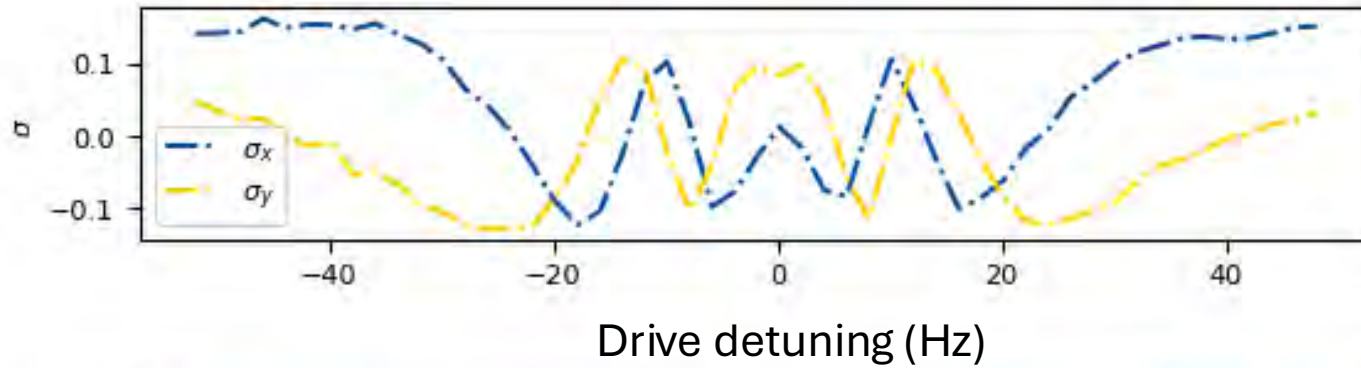
# Membrane spectroscopy (via Ramsey)



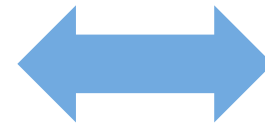
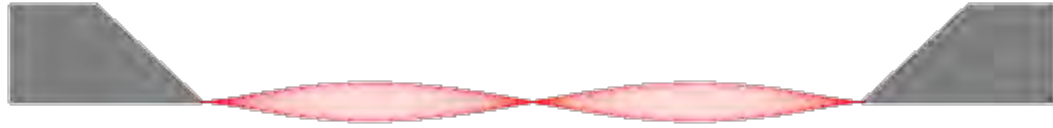
$$\Omega_m/2\pi = 4.4 \text{ MHz}$$



Dispersive interaction

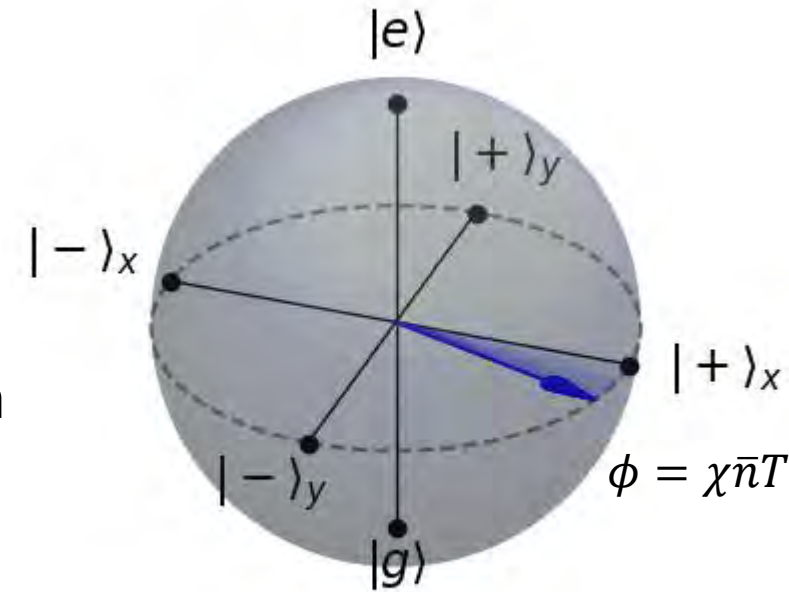
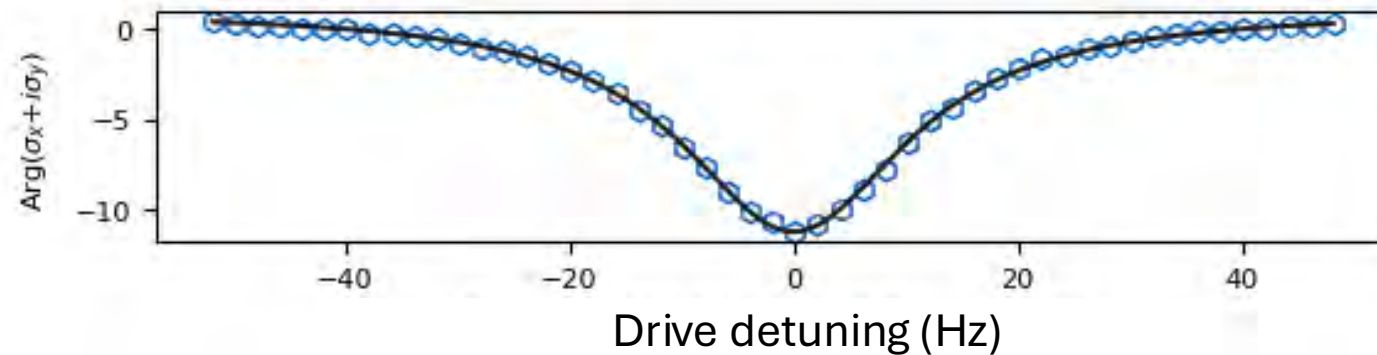
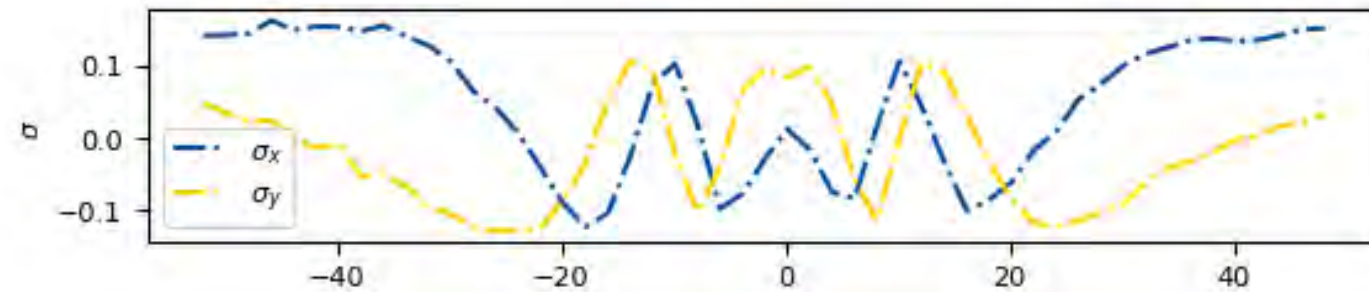


$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$



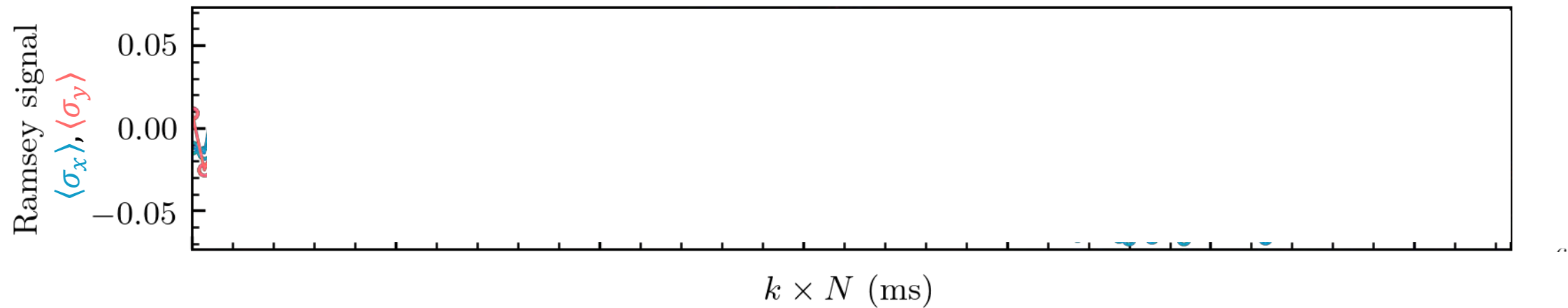
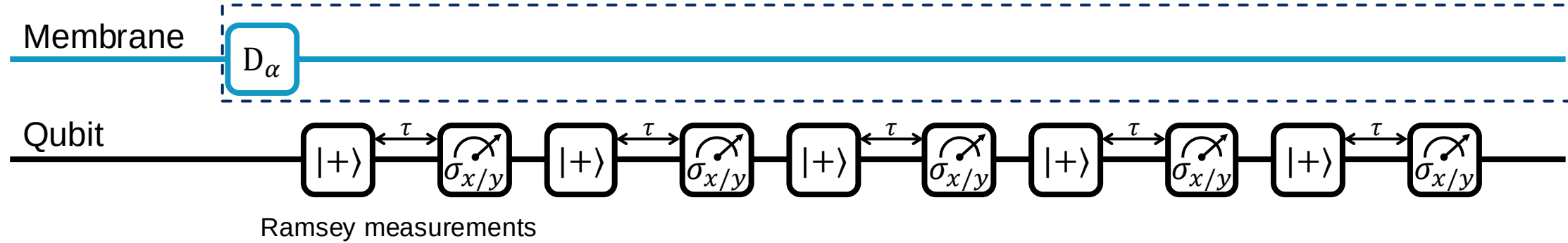
Dispersive interaction

$$\Omega_m/2\pi = 4.4 \text{ MHz}$$

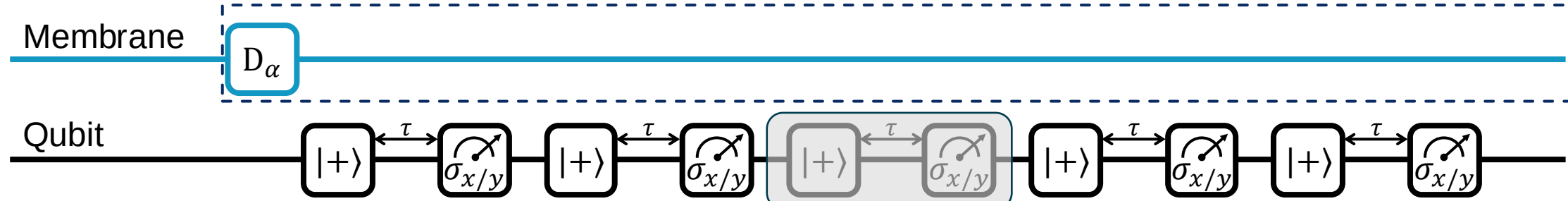


$$\omega_{eg}/2\pi = 2.2 \text{ MHz}$$

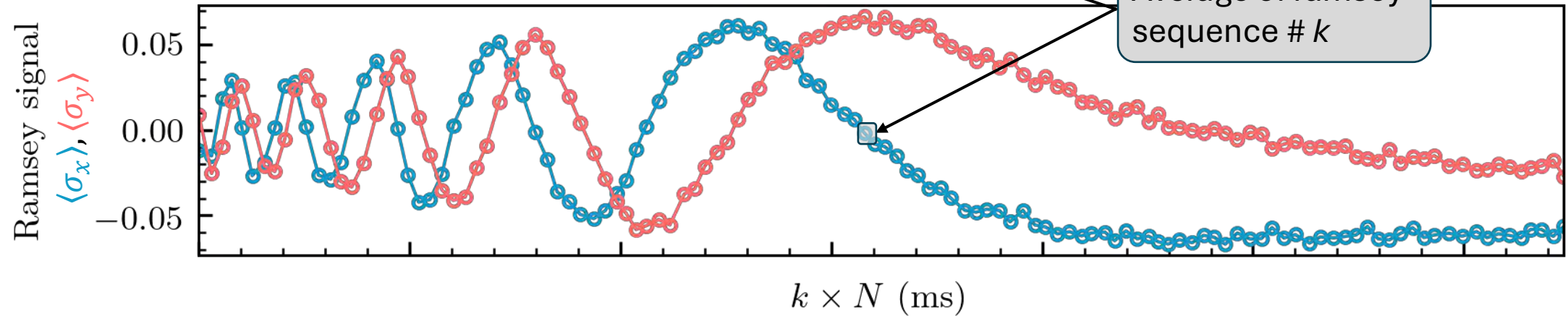
Lifetime large compared to control durations



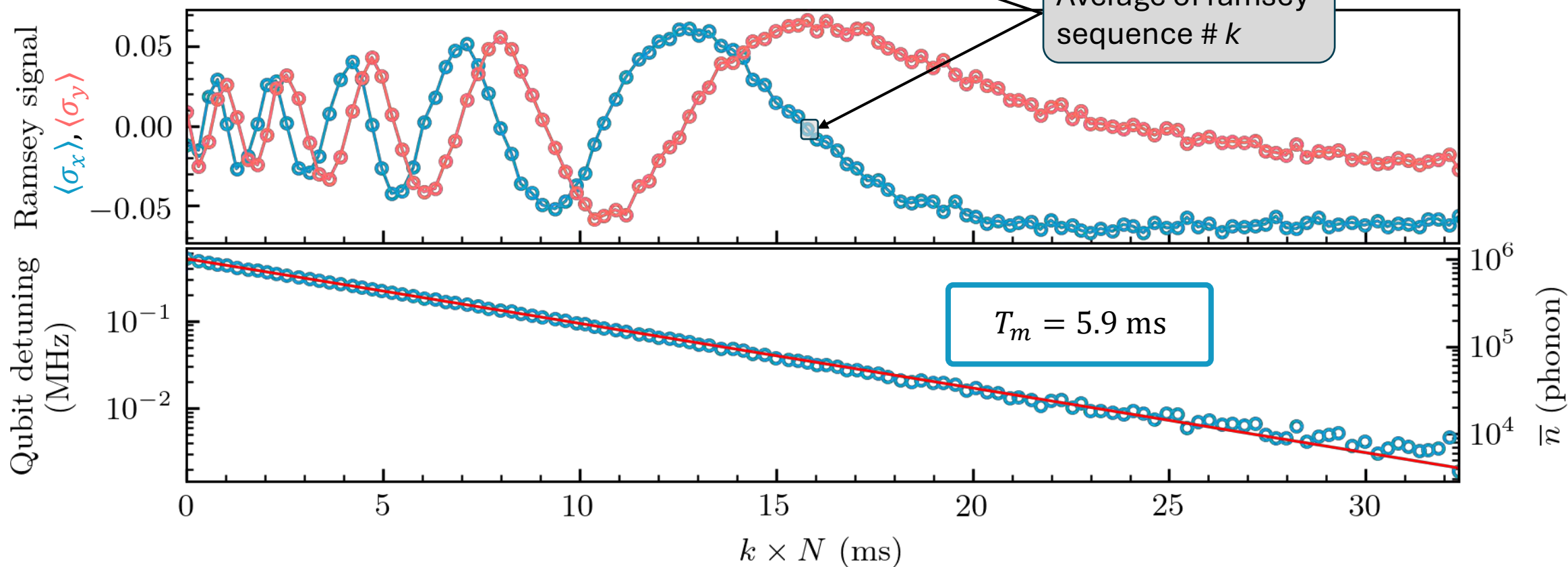
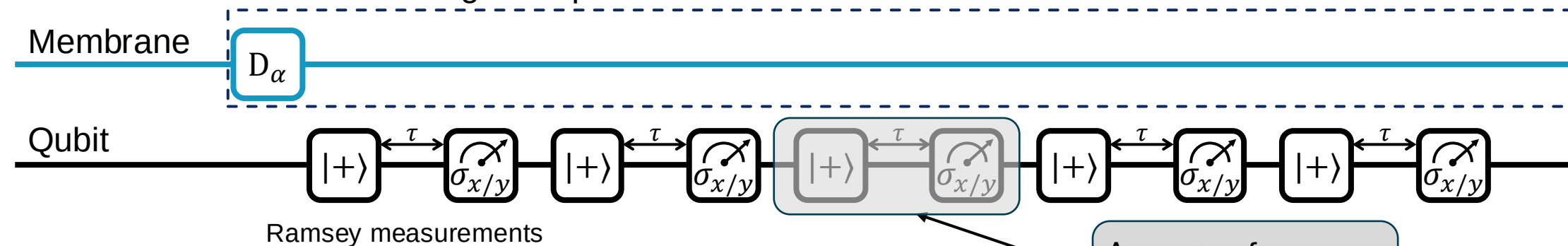
Lifetime large compared to control durations

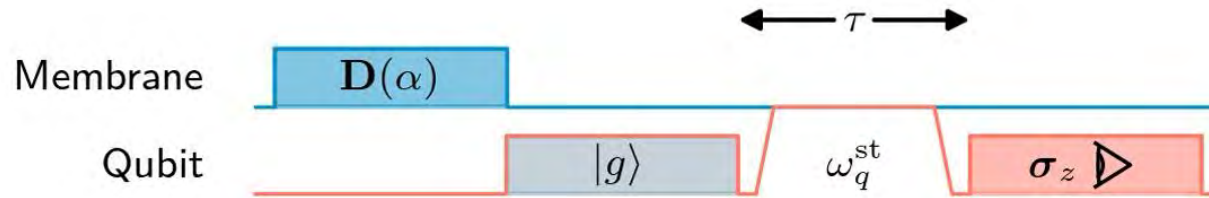


Ramsey measurements



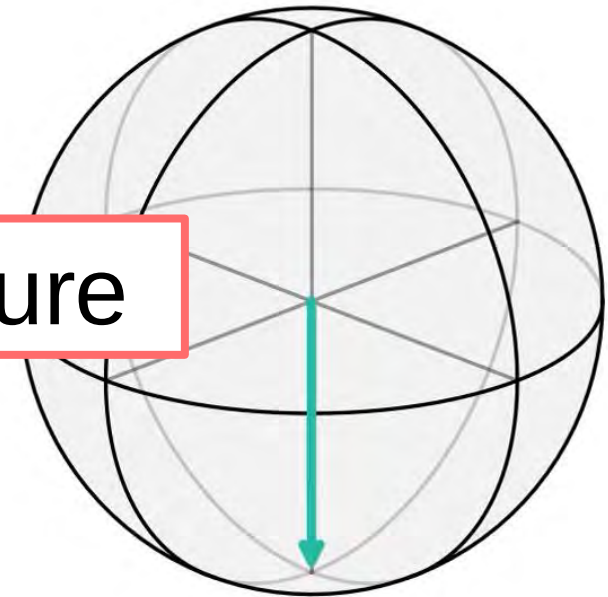
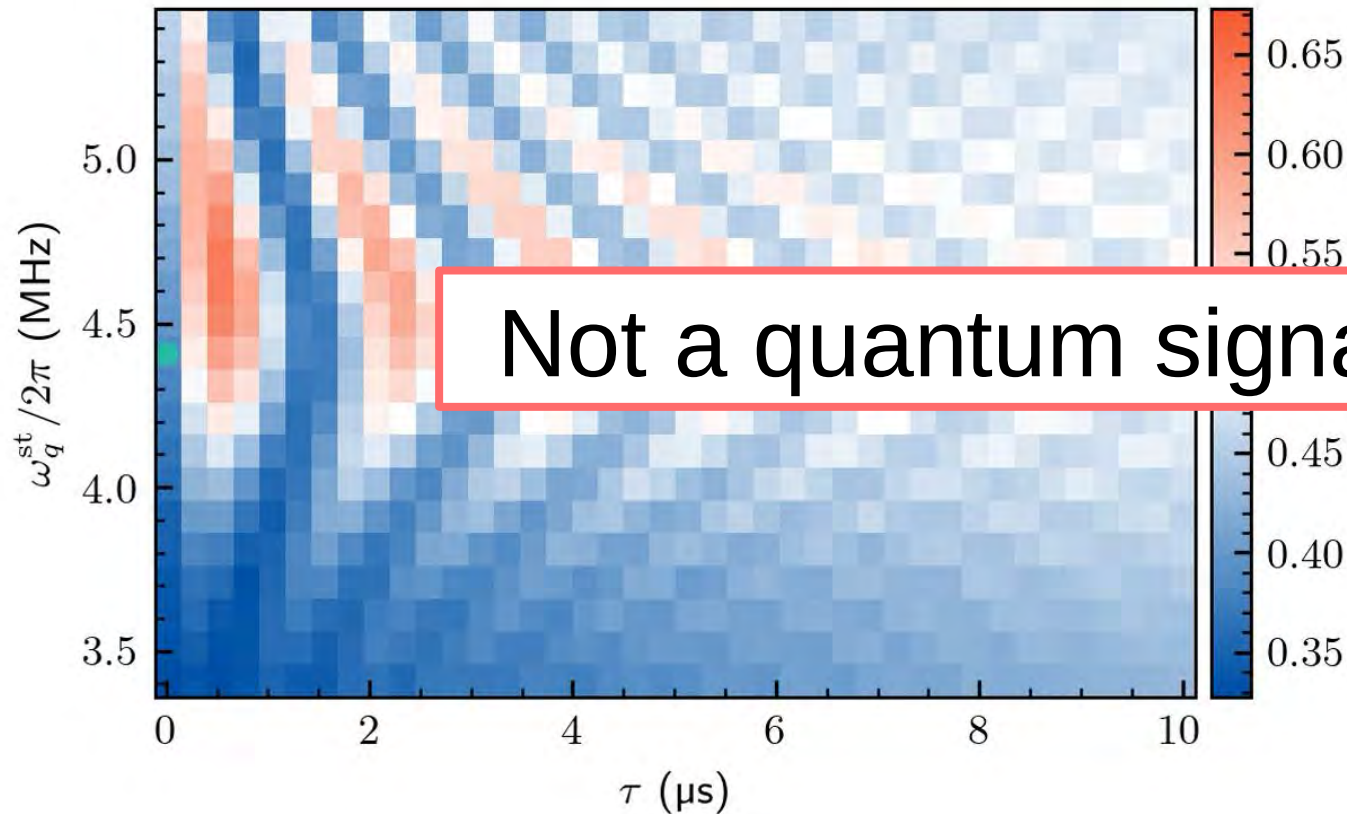
Lifetime large compared to control durations



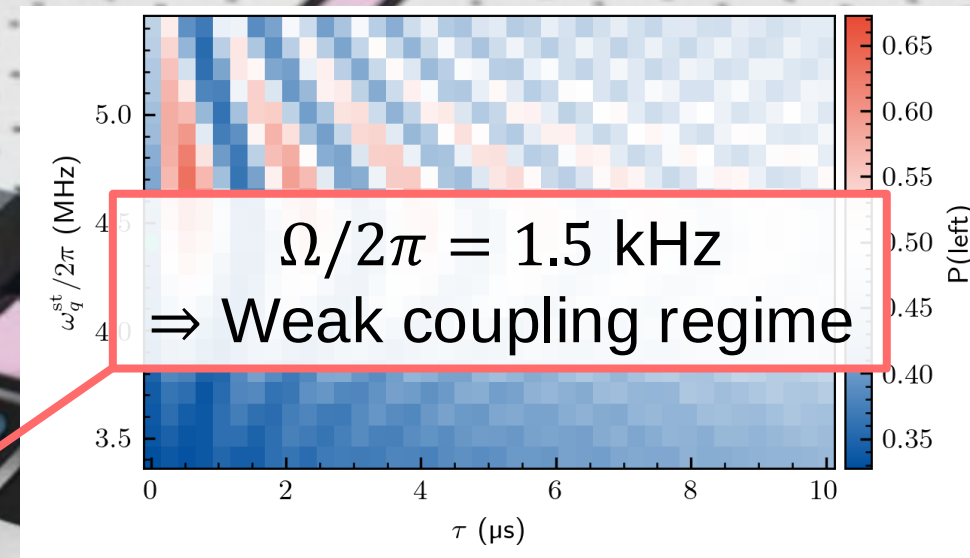
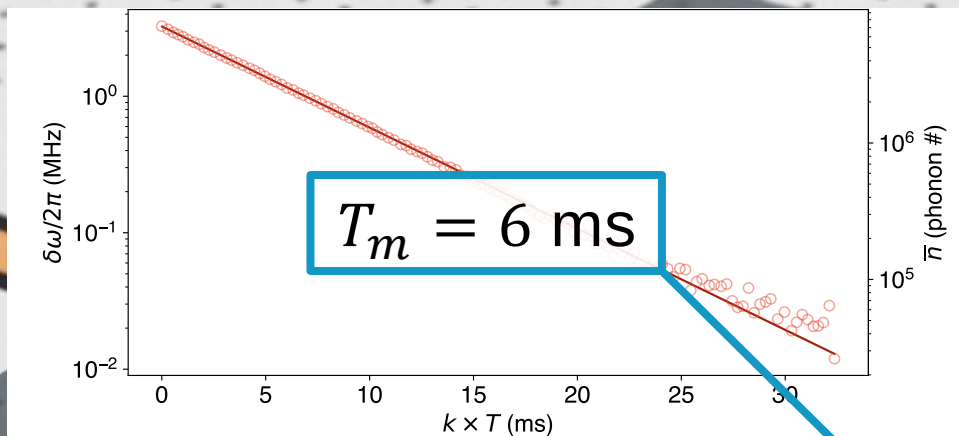


$$\theta = \Omega|\alpha|\tau$$

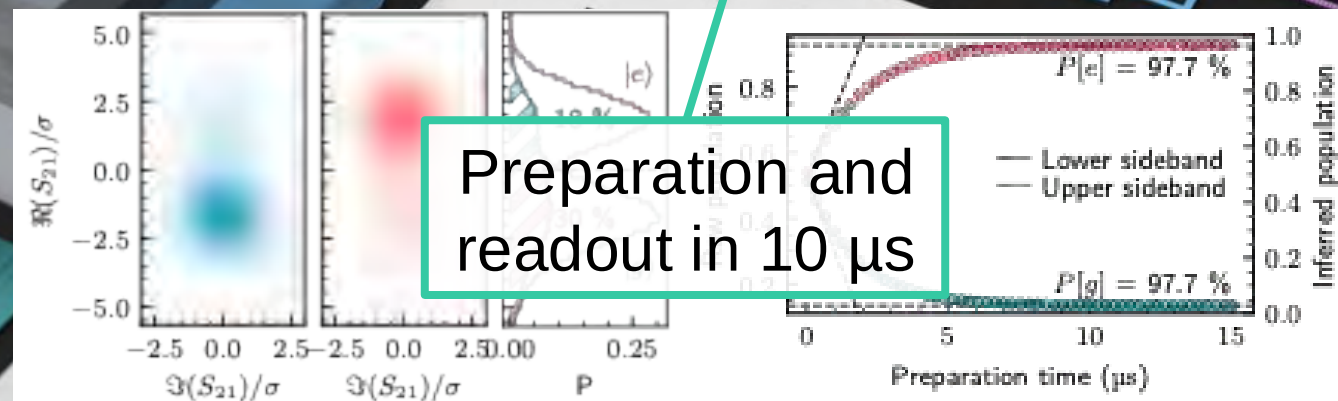
$$|\alpha|^2 \sim 10^5 \text{ phonons}$$



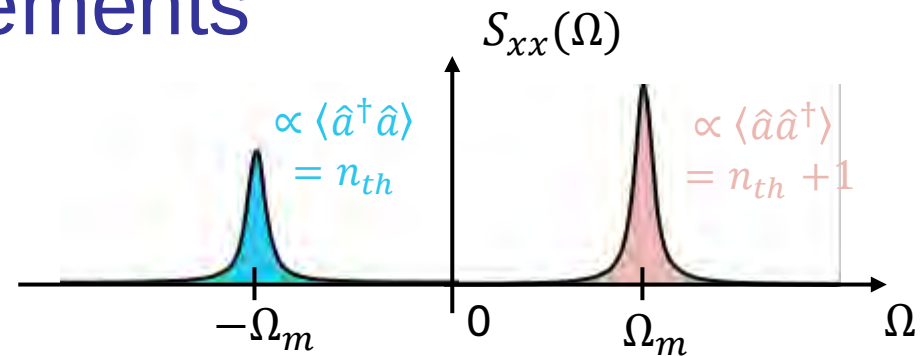
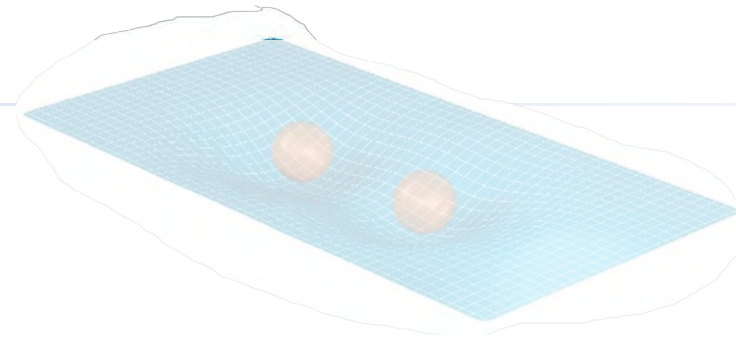




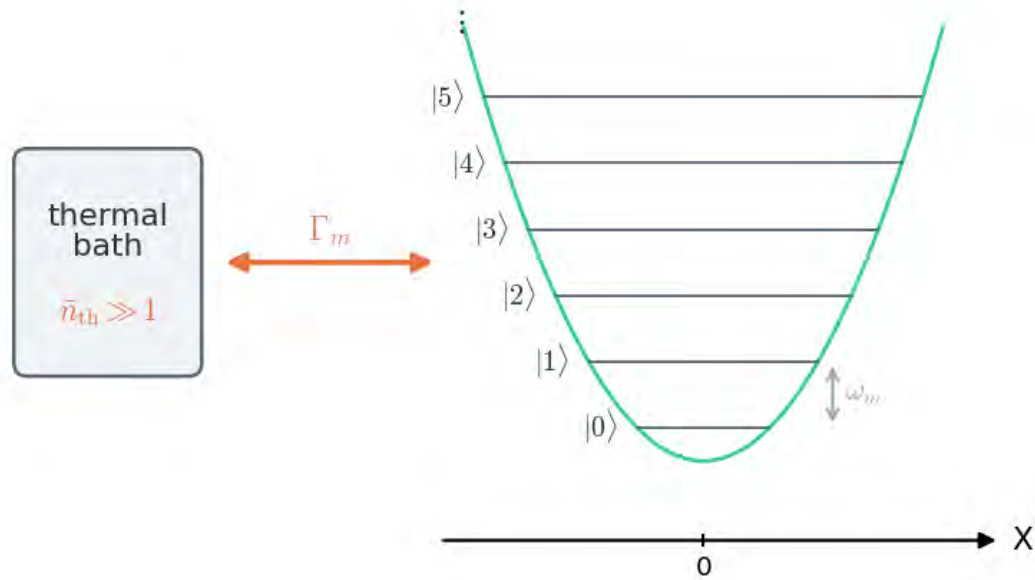
300 measurements in a lifetime



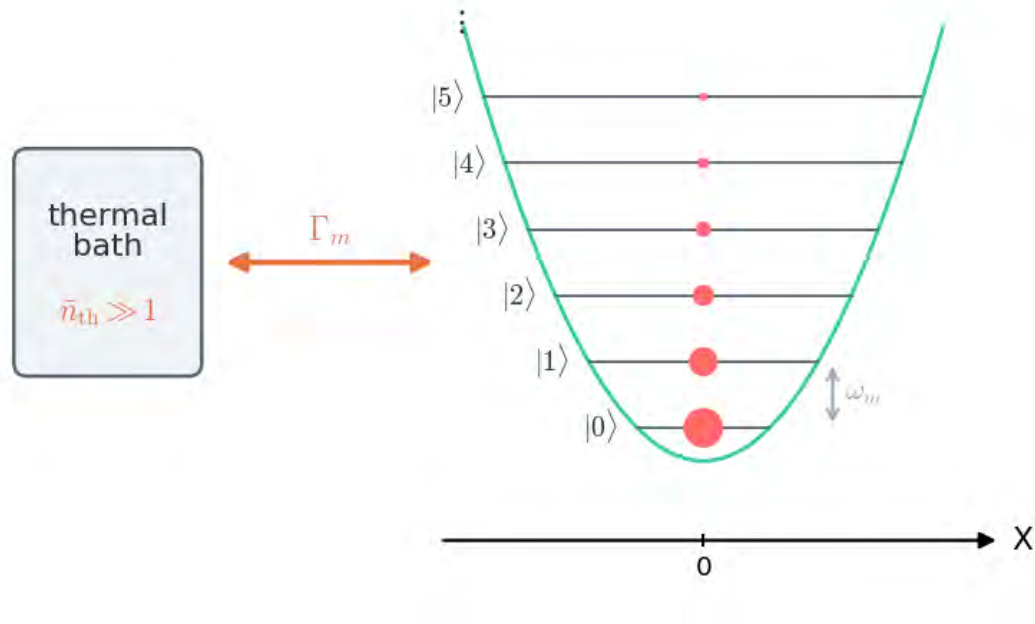
- Introduction/Motivation
- Radio-frequency fluxonium qubits
- Hybrid qubit/mechanical system: MecaFlux
- Quantum asymmetry measurements
- Perspectives



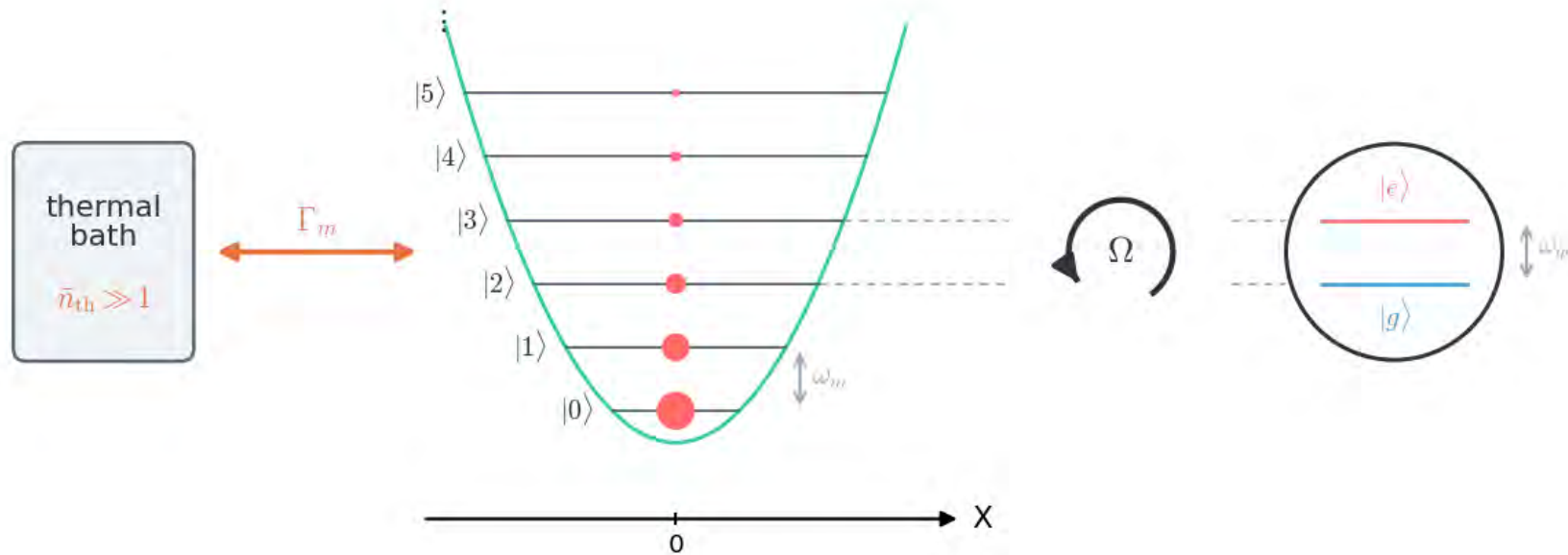
# A thought experiment ...



# A thought experiment ...



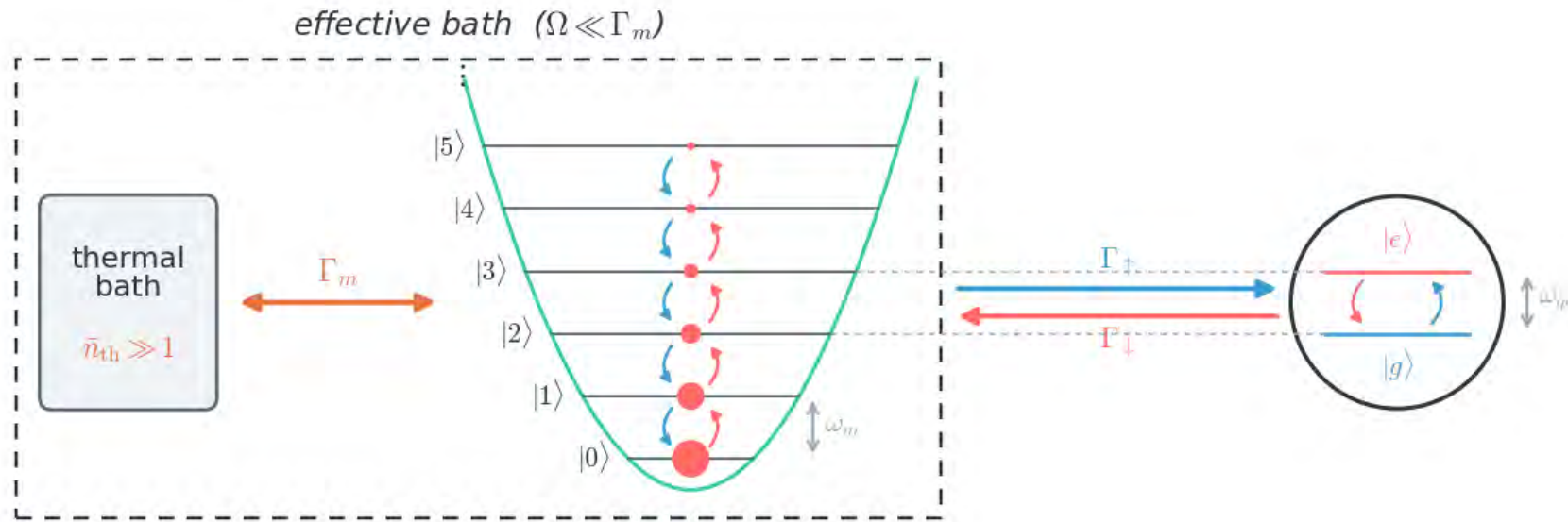
# A thought experiment ...



A qubit as a quantum spectrum analyzer :

- linearly couple the system observable to a qubit :  $H = \hbar \Omega \hat{X} \otimes \hat{\sigma}_x$

# A thought experiment ...

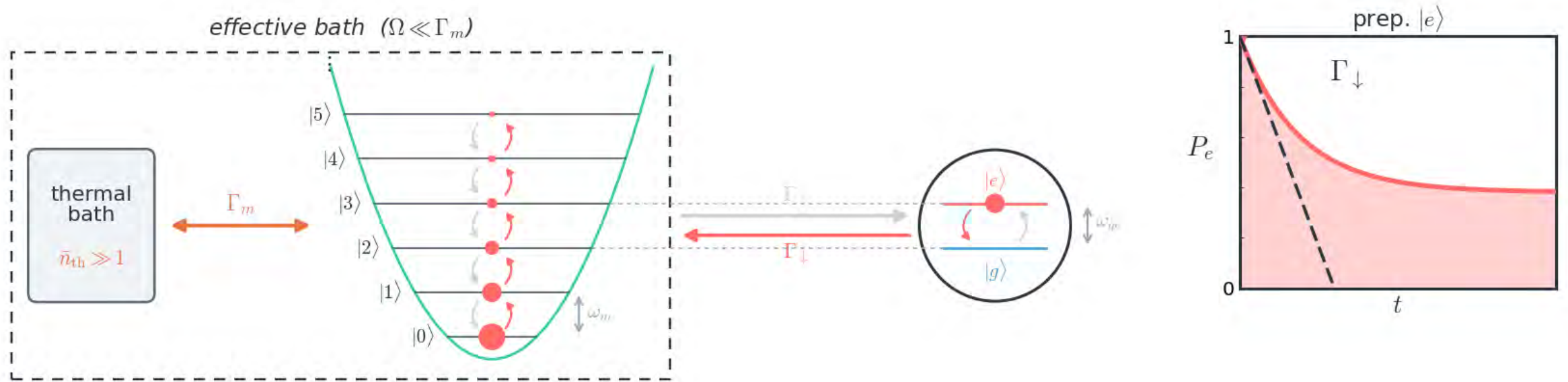


A qubit as a quantum spectrum analyzer :

- linearly couple the system observable to a qubit :  $H = \hbar \Omega \hat{X} \otimes \hat{\sigma}_x$
- in the weak-coupling regime, the system acts as an effective bath



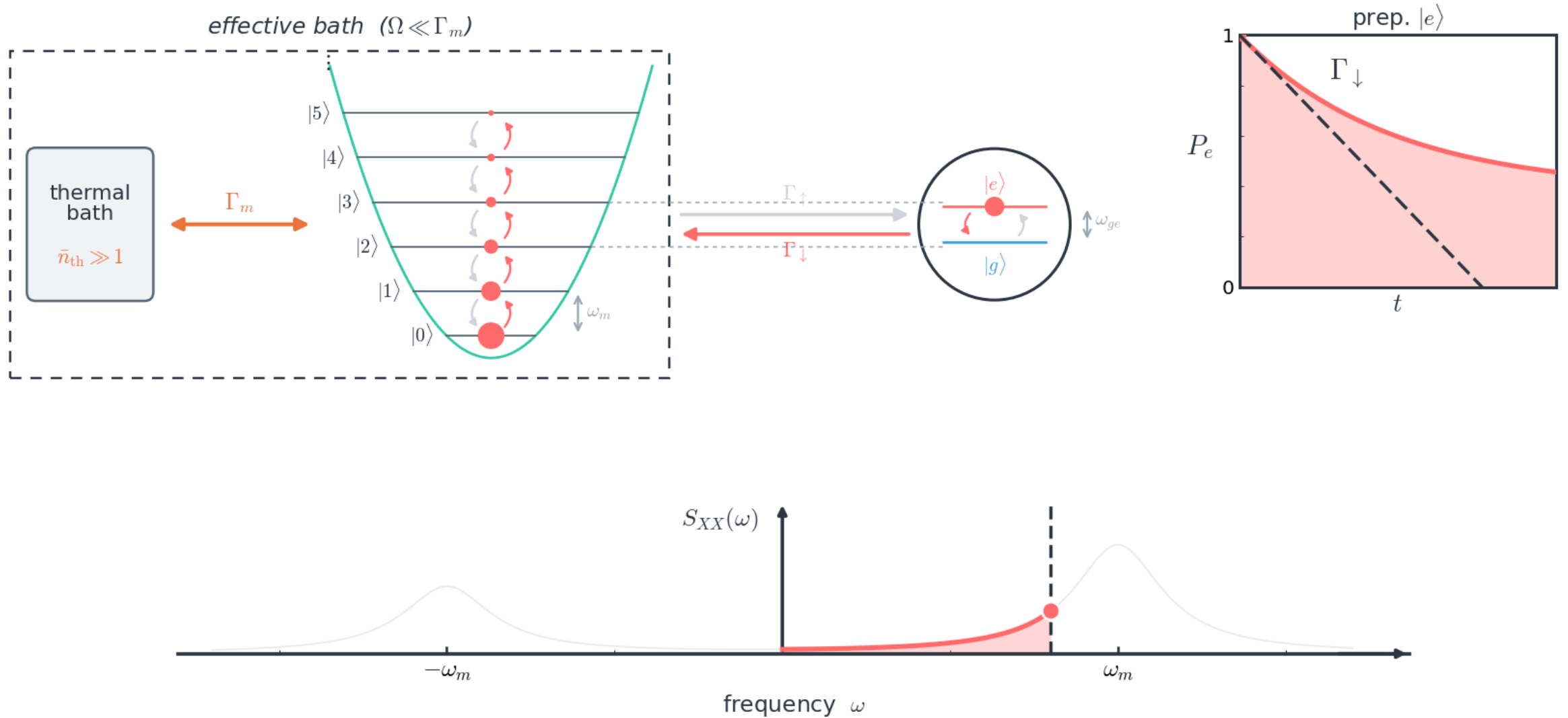
# A thought experiment ...



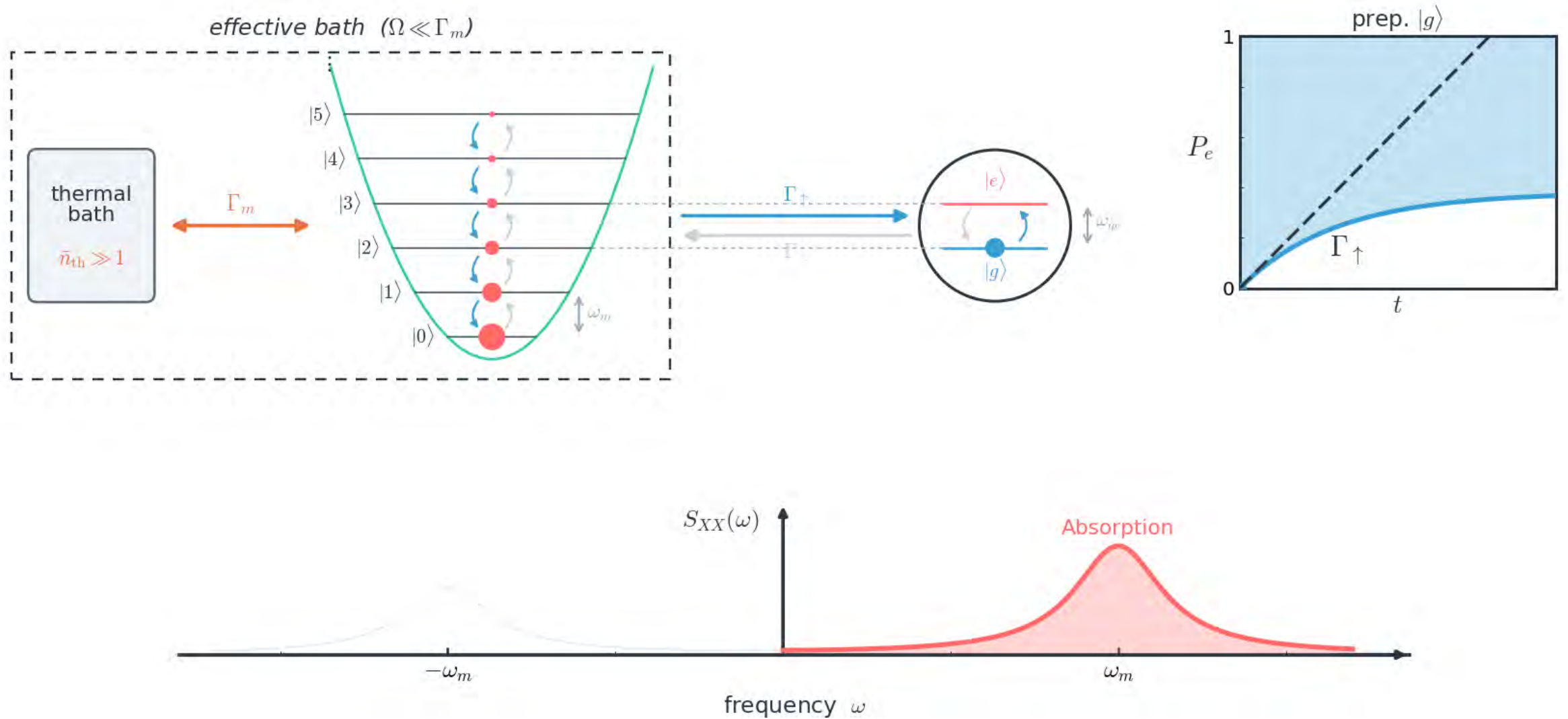
A qubit as a quantum spectrum analyzer :

- linearly couple the system observable to a qubit :  $H = \hbar \Omega \hat{X} \otimes \hat{\sigma}_x$
- in the weak-coupling regime, the system acts as an effective bath
- the rates  $\Gamma_\downarrow \propto S_{XX}(\omega_{ge})$  and  $\Gamma_\uparrow \propto S_{XX}(-\omega_{ge})$

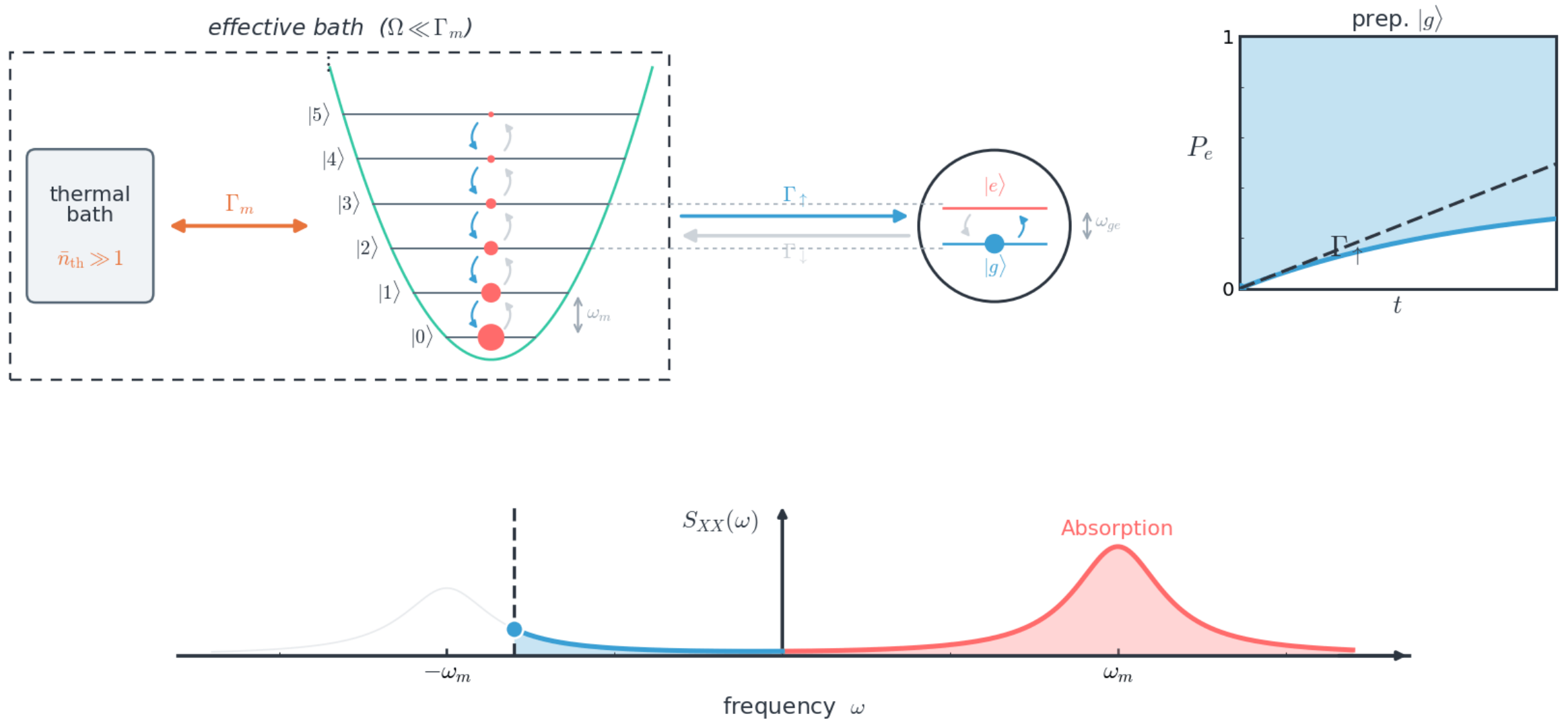
# A thought experiment ...



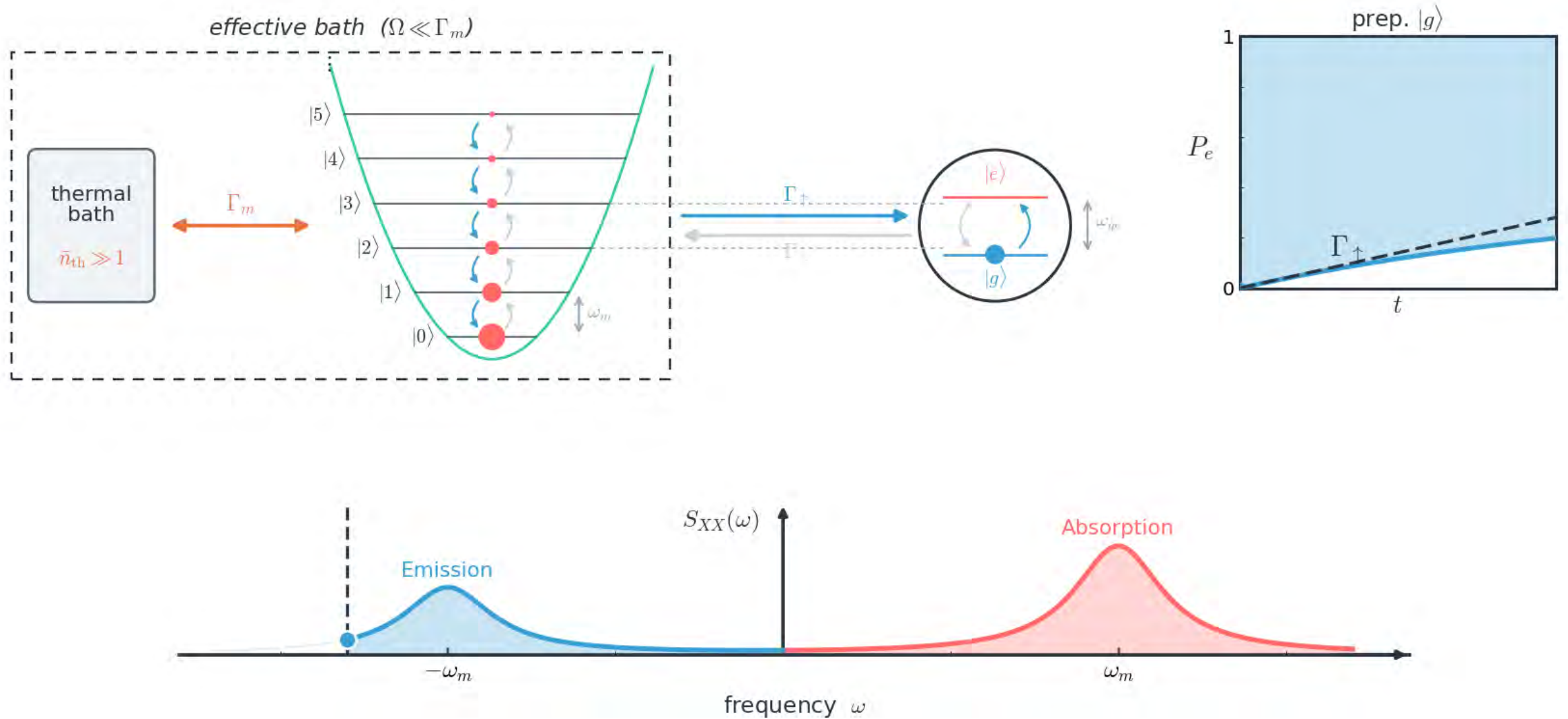
# A thought experiment ...



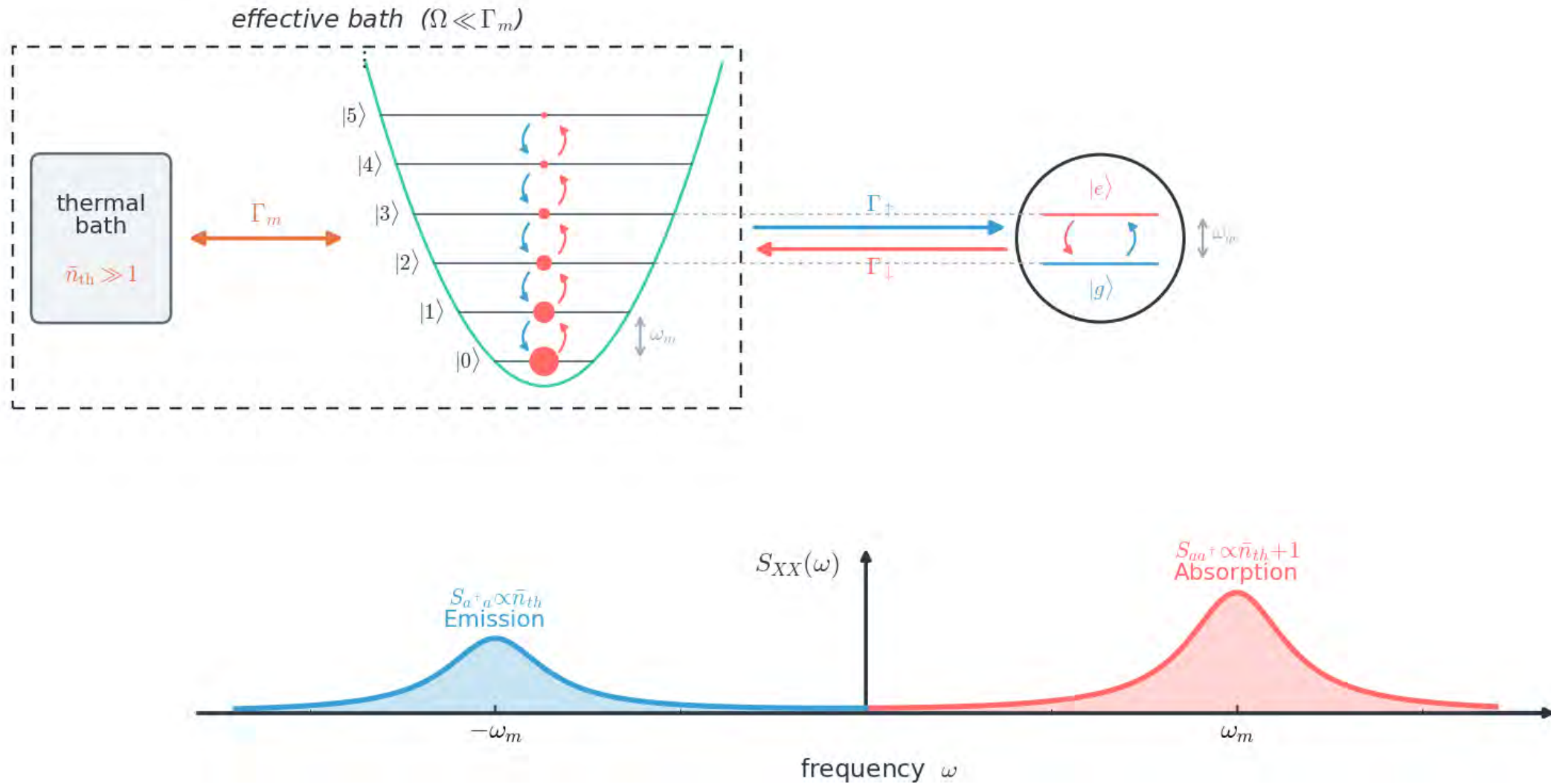
# A thought experiment ...



# A thought experiment ...

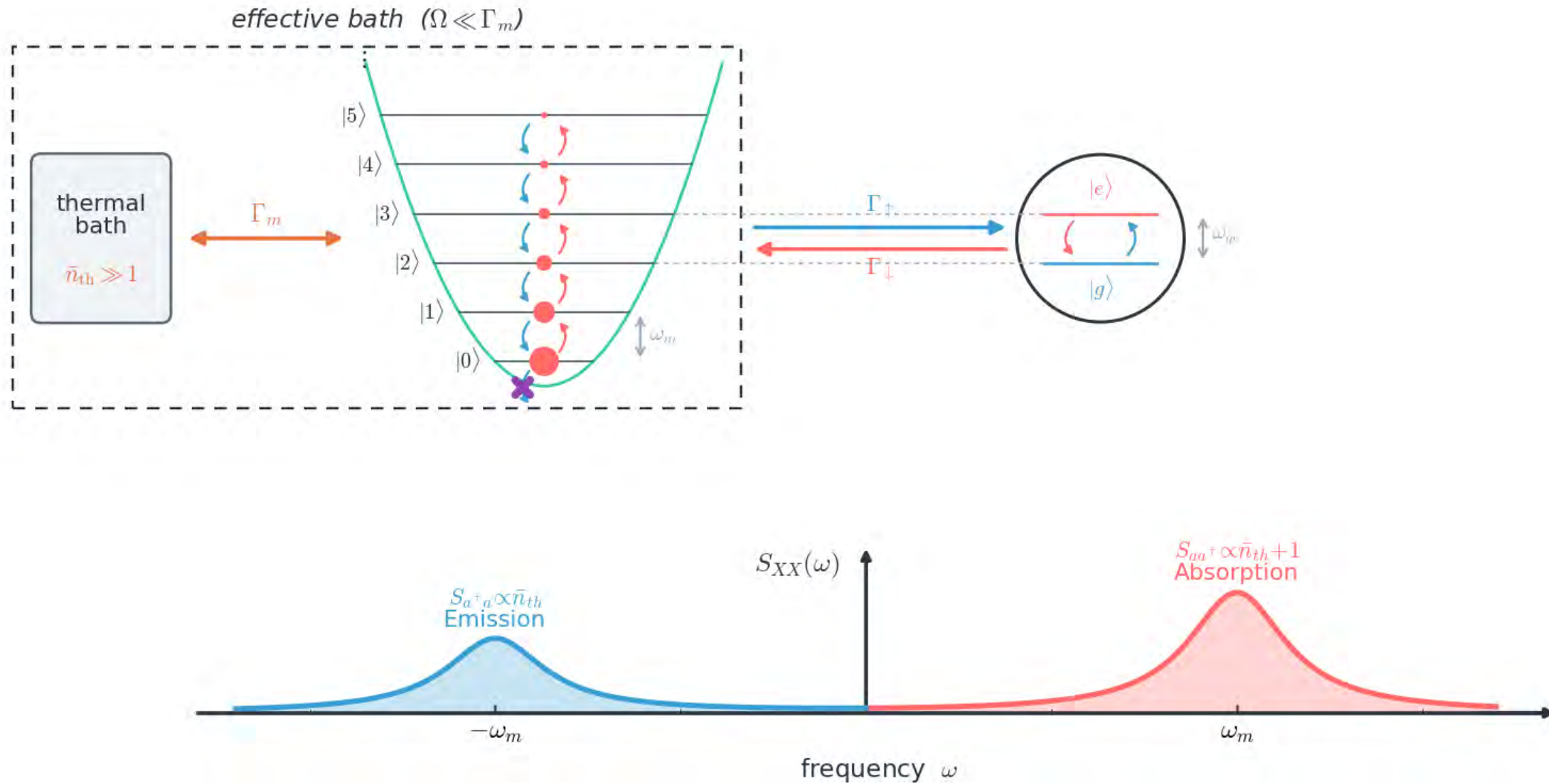


# A thought experiment ...

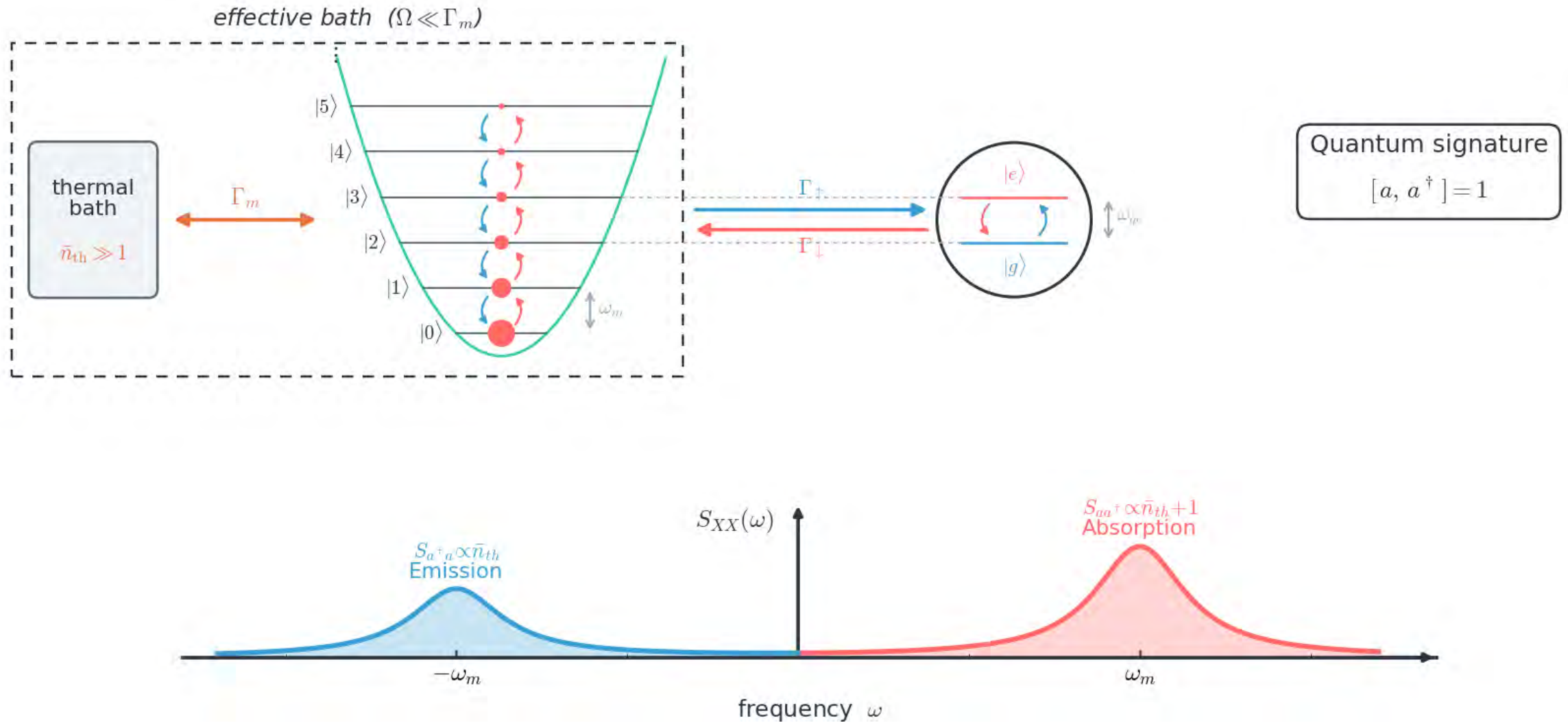




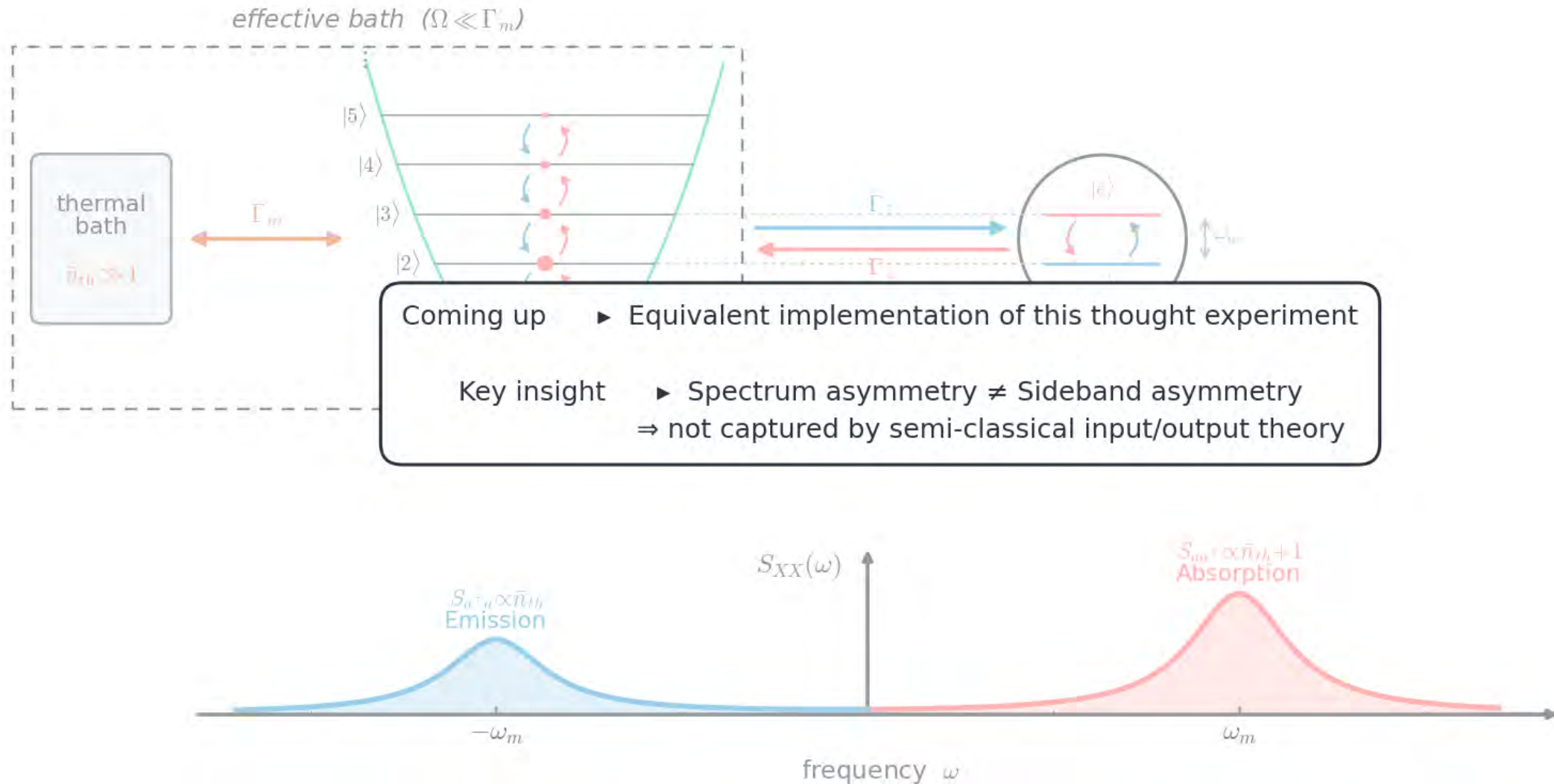
# A thought experiment ...



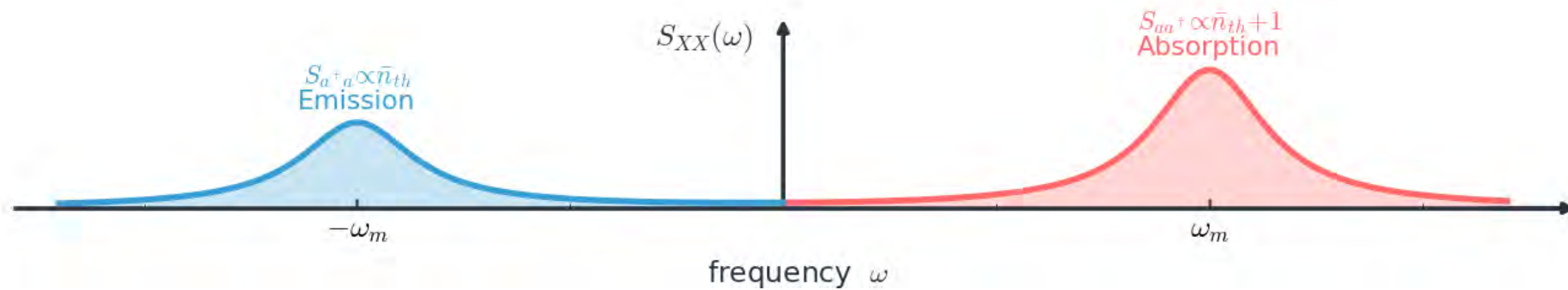
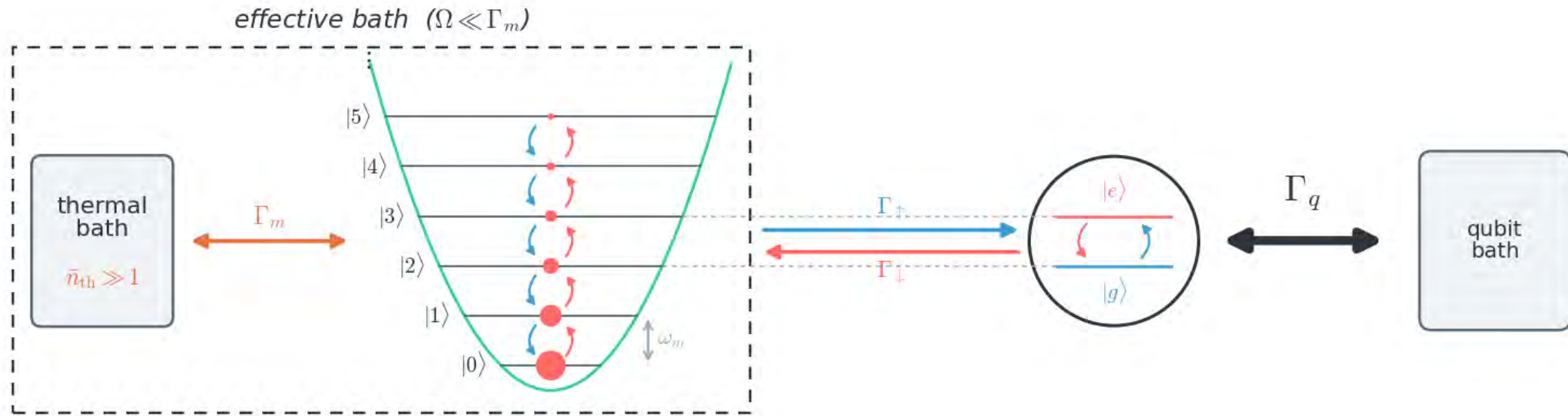
# A thought experiment ...



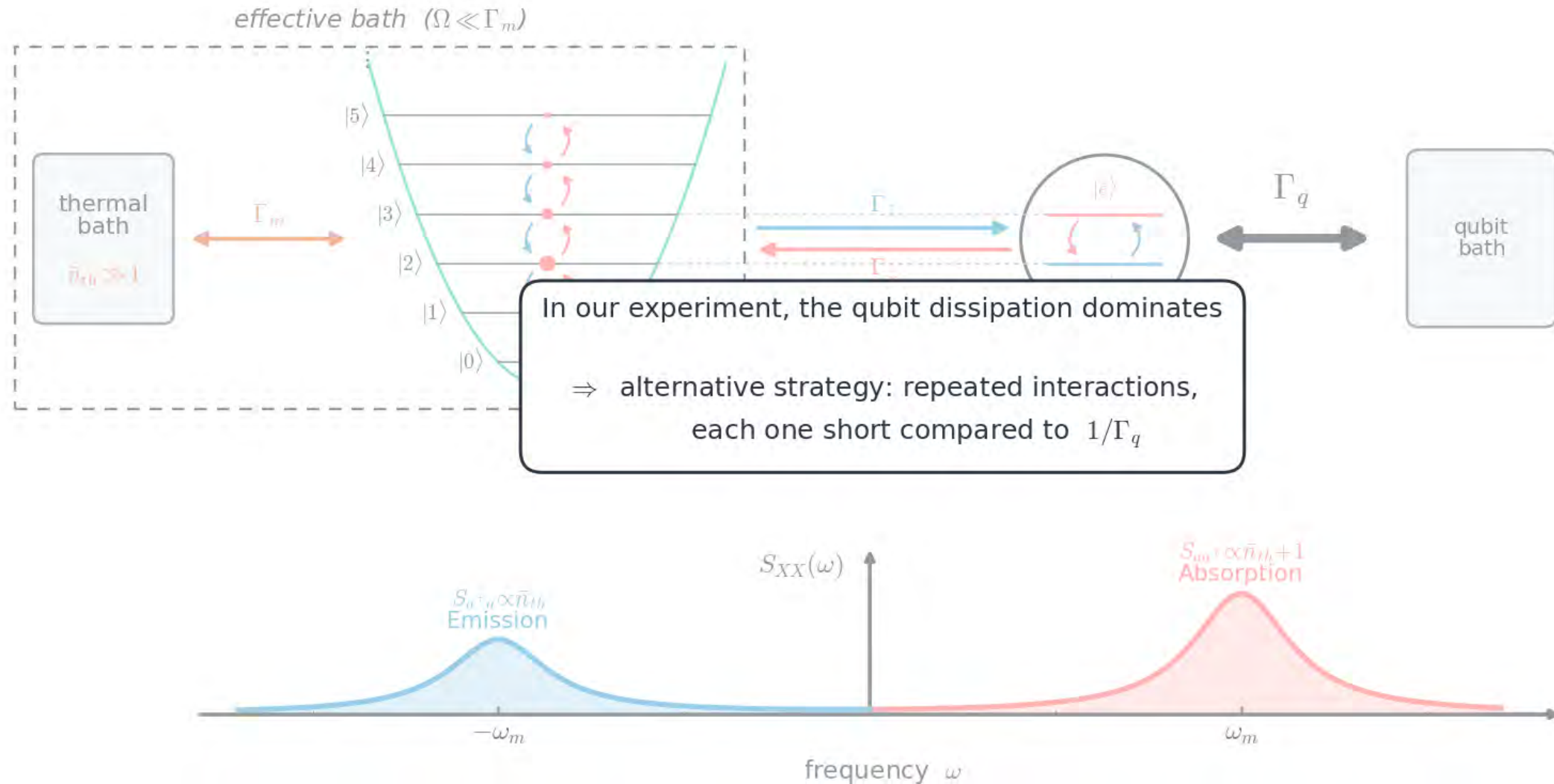
# A thought experiment ...



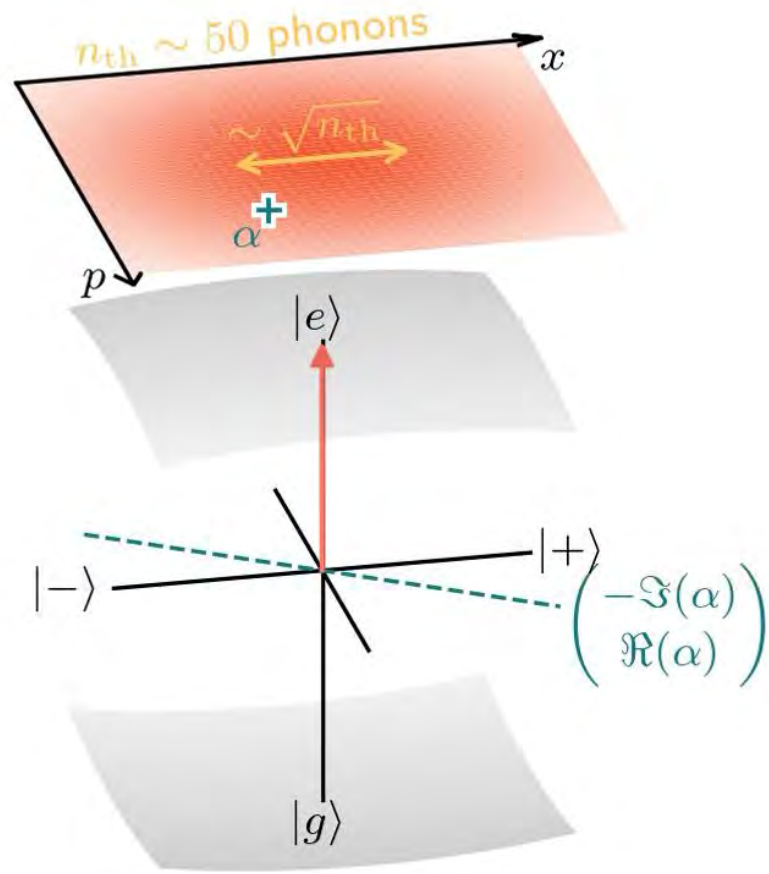
# A thought experiment ...



# A thought experiment ...



# Spectrum analyzer measurement

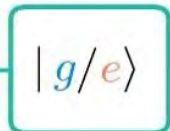


$$\theta \simeq \Omega T_{int} \sqrt{n_{th}} \approx 0.2 \text{ rad} \ll 1$$

1. Mapping:  $x \leftrightarrow \sigma_x$

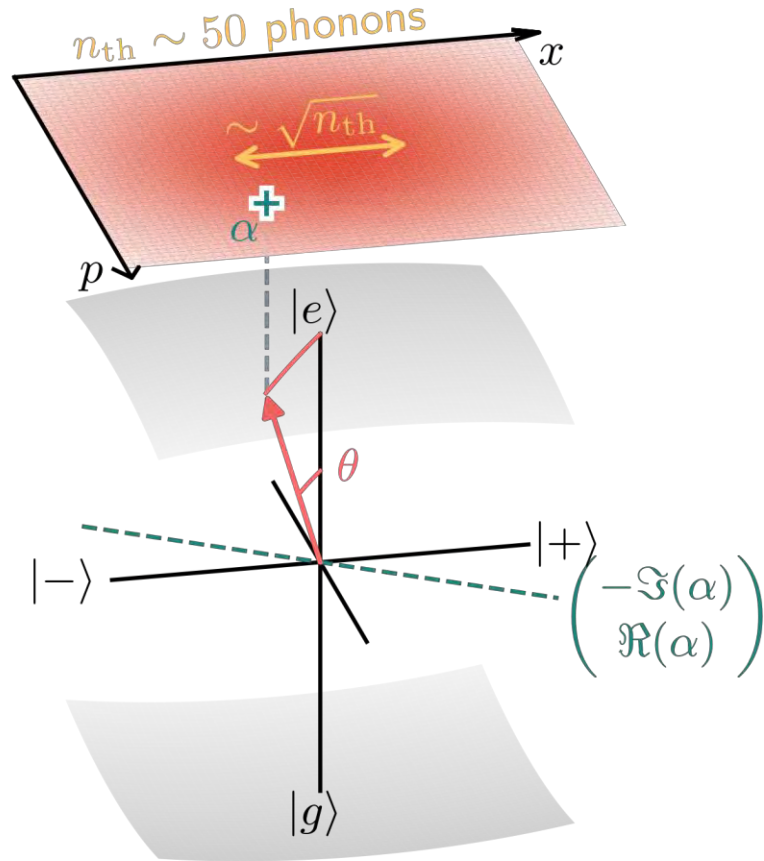
Membrane

Qubit





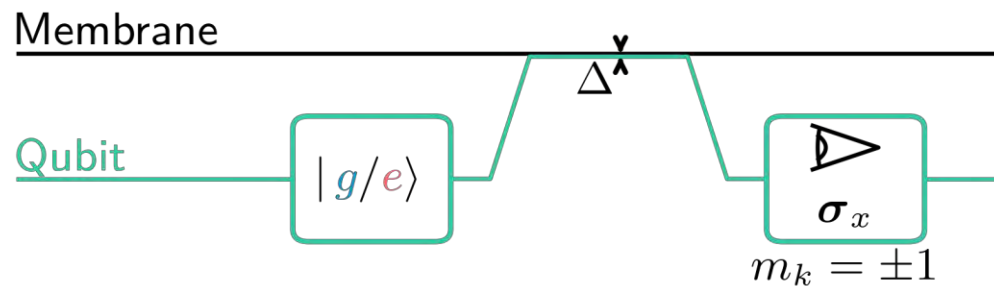
# Spectrum analyzer measurement



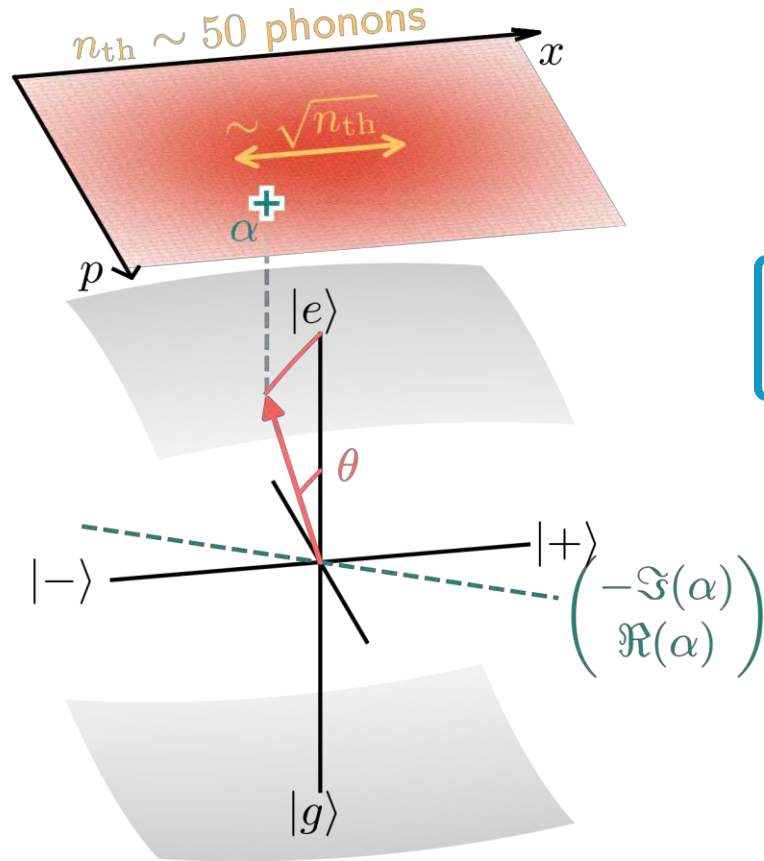
$$\theta \simeq \Omega T_{int} \sqrt{n_{th}} \approx 0.2 \text{ rad} \ll 1$$

1. Mapping:  $x \leftrightarrow \sigma_x$

2. Measure qubit  $\sigma_x$  projection



# Spectrum analyzer measurement

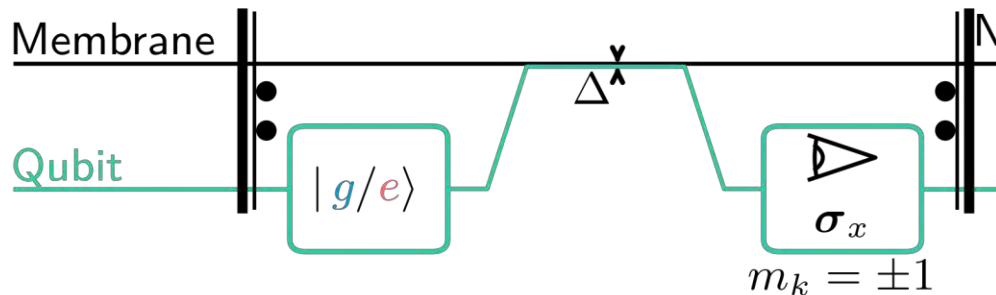


$$\theta \simeq \Omega T_{int} \sqrt{n_{th}} \approx 0.2 \text{ rad} \ll 1$$

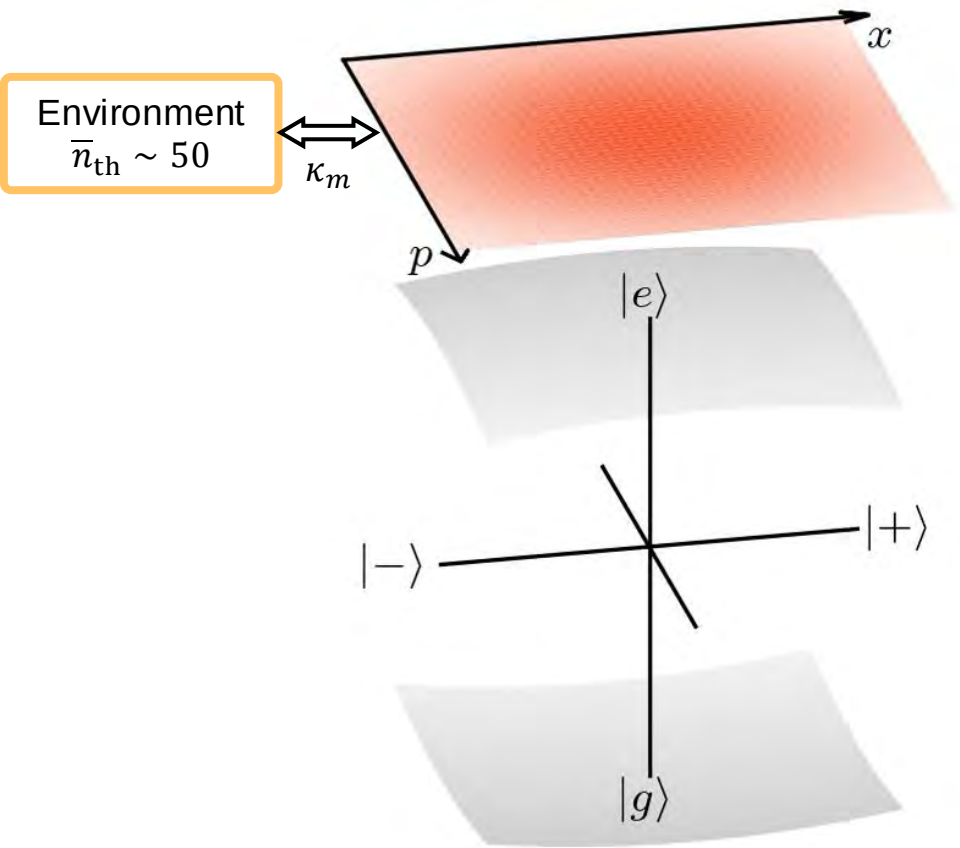
1. Mapping:  $x \leftrightarrow \sigma_x$

2. Measure qubit  $\sigma_x$  projection

Each measurements yields **one bit** carrying **weak information**; however, we can accumulate **many detections** during  $T_{1,m}$



# Spectrum analyzer measurement

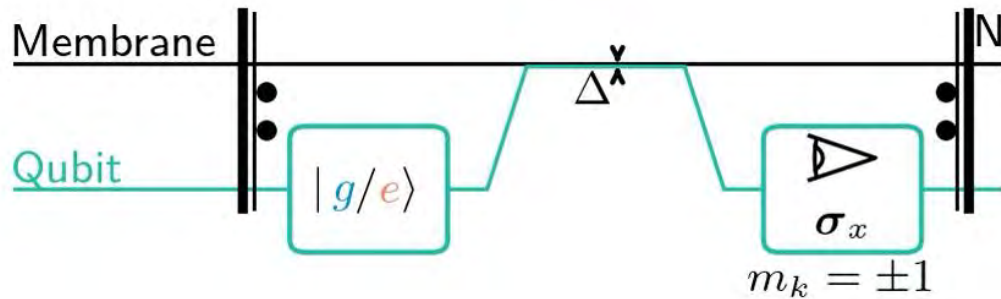
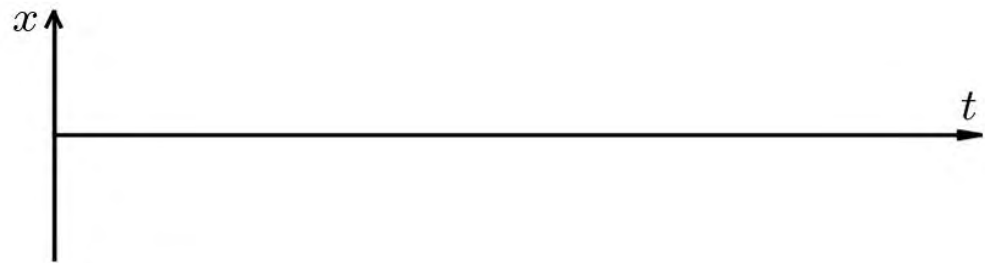


$$\theta \simeq \Omega T_{int} \sqrt{n_{th}} \approx 0.2 \text{ rad} \ll 1$$

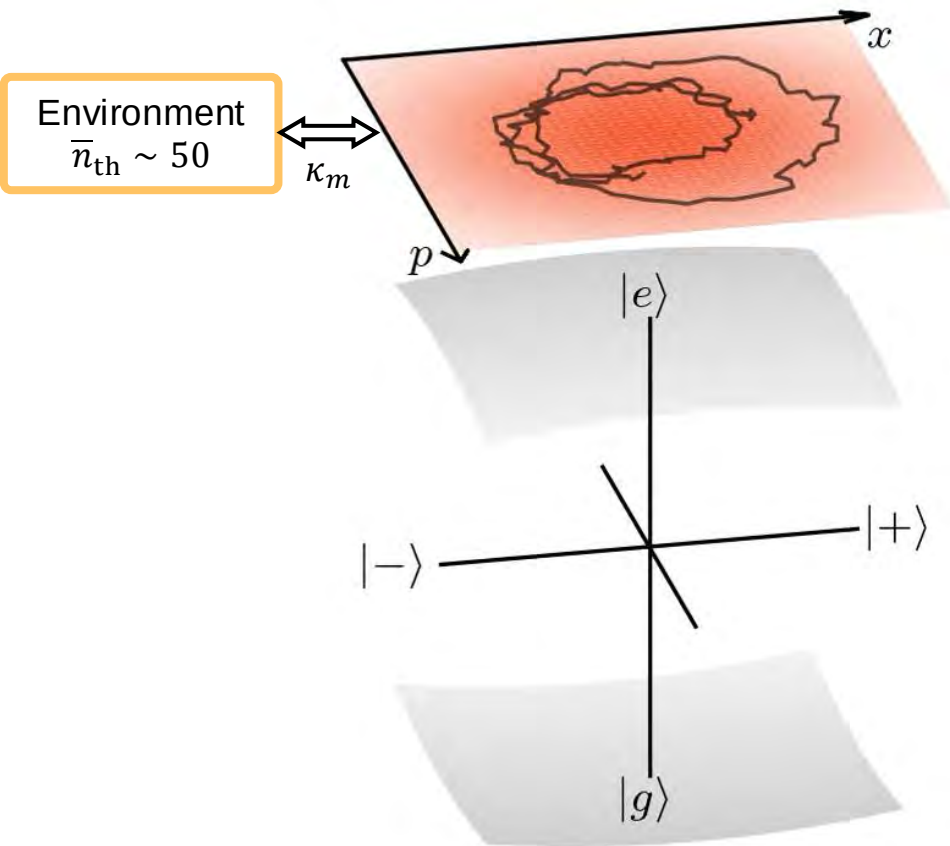
1. Mapping:  $x \leftrightarrow \sigma_x$

2. Measure qubit  $\sigma_x$  projection

Each measurements yields **one bit** carrying **weak information**; however, we can accumulate **many detections** during  $T_{1,m}$



# Spectrum analyzer measurement

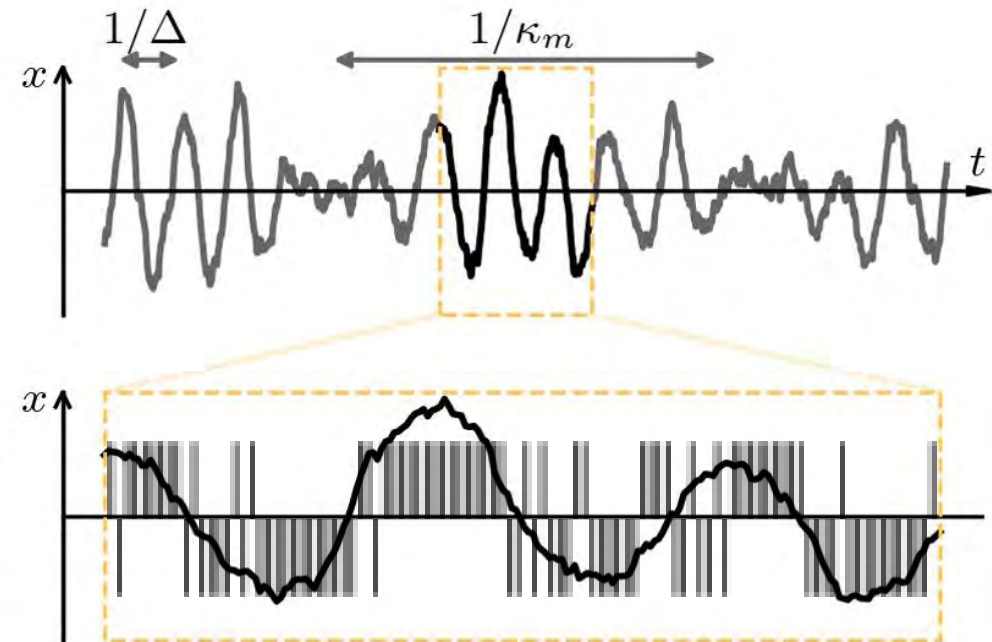
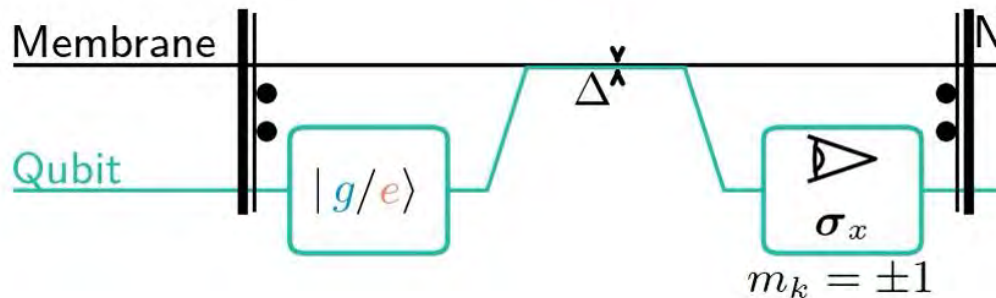


$$\theta \simeq \Omega T_{int} \sqrt{n_{th}} \approx 0.2 \text{ rad} \ll 1$$

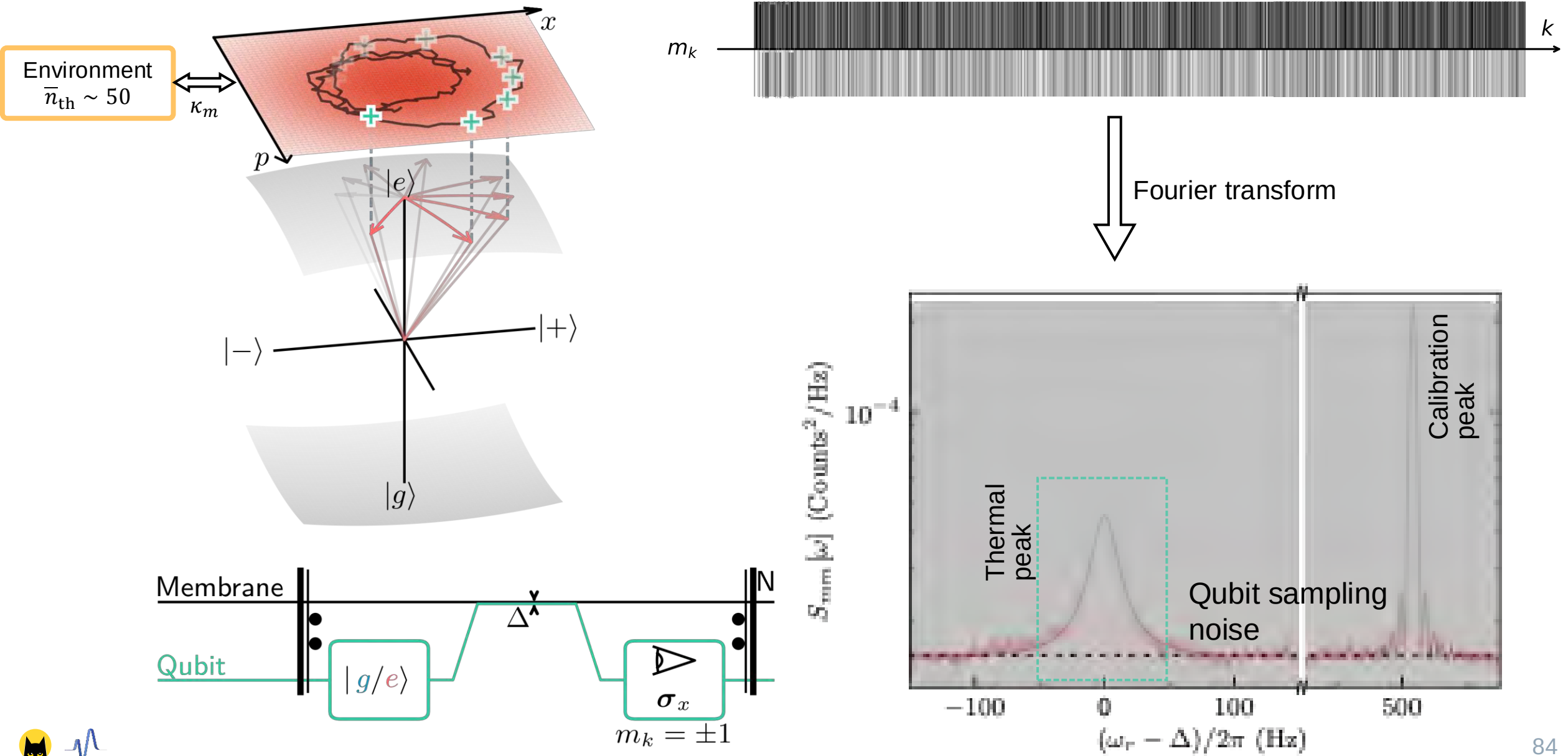
1. Mapping:  $x \leftrightarrow \sigma_x$

2. Measure qubit  $\sigma_x$  projection

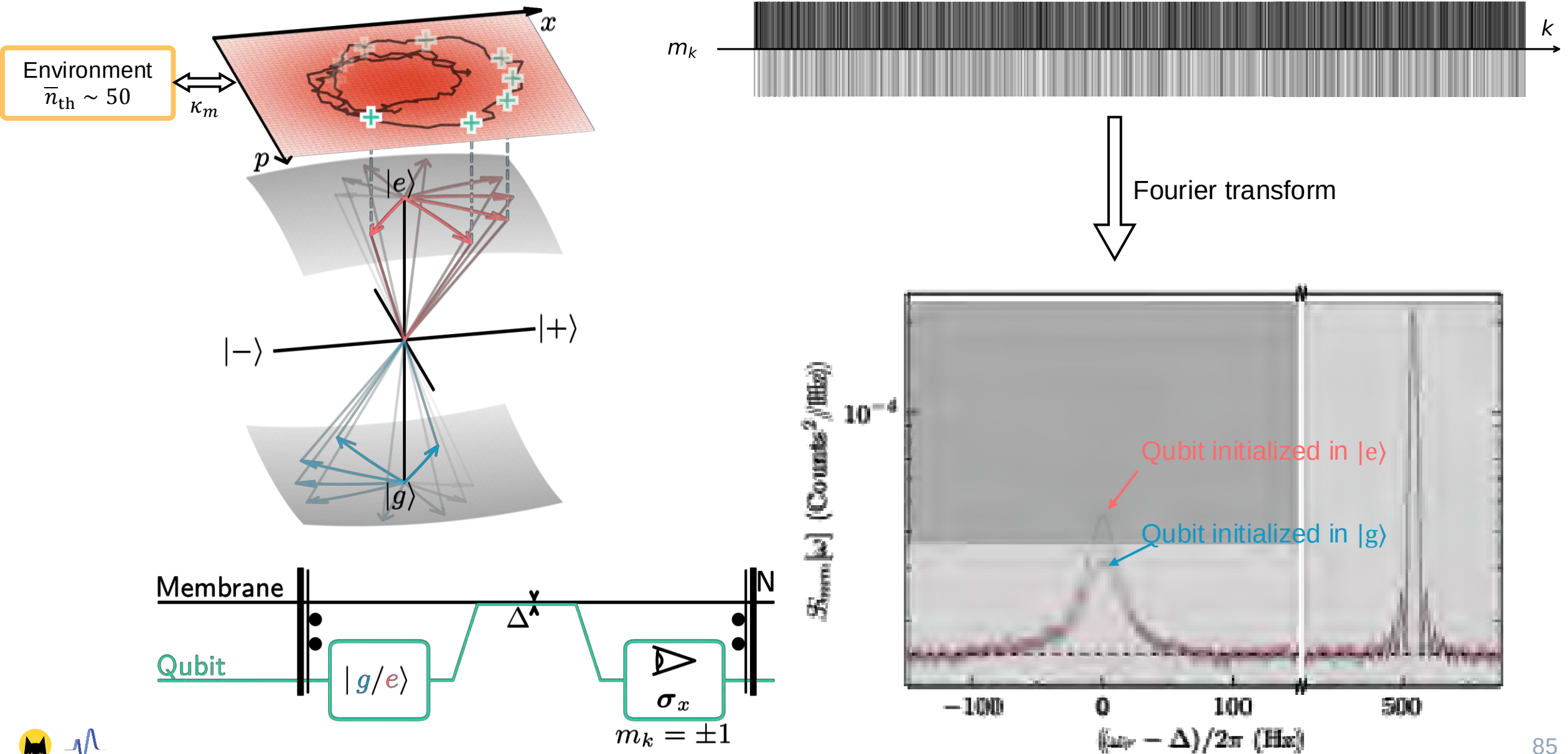
Each measurements yields **one bit** carrying **weak information**; however, we can accumulate **many detections** during  $T_{1,m}$



# Spectrum analyzer measurement

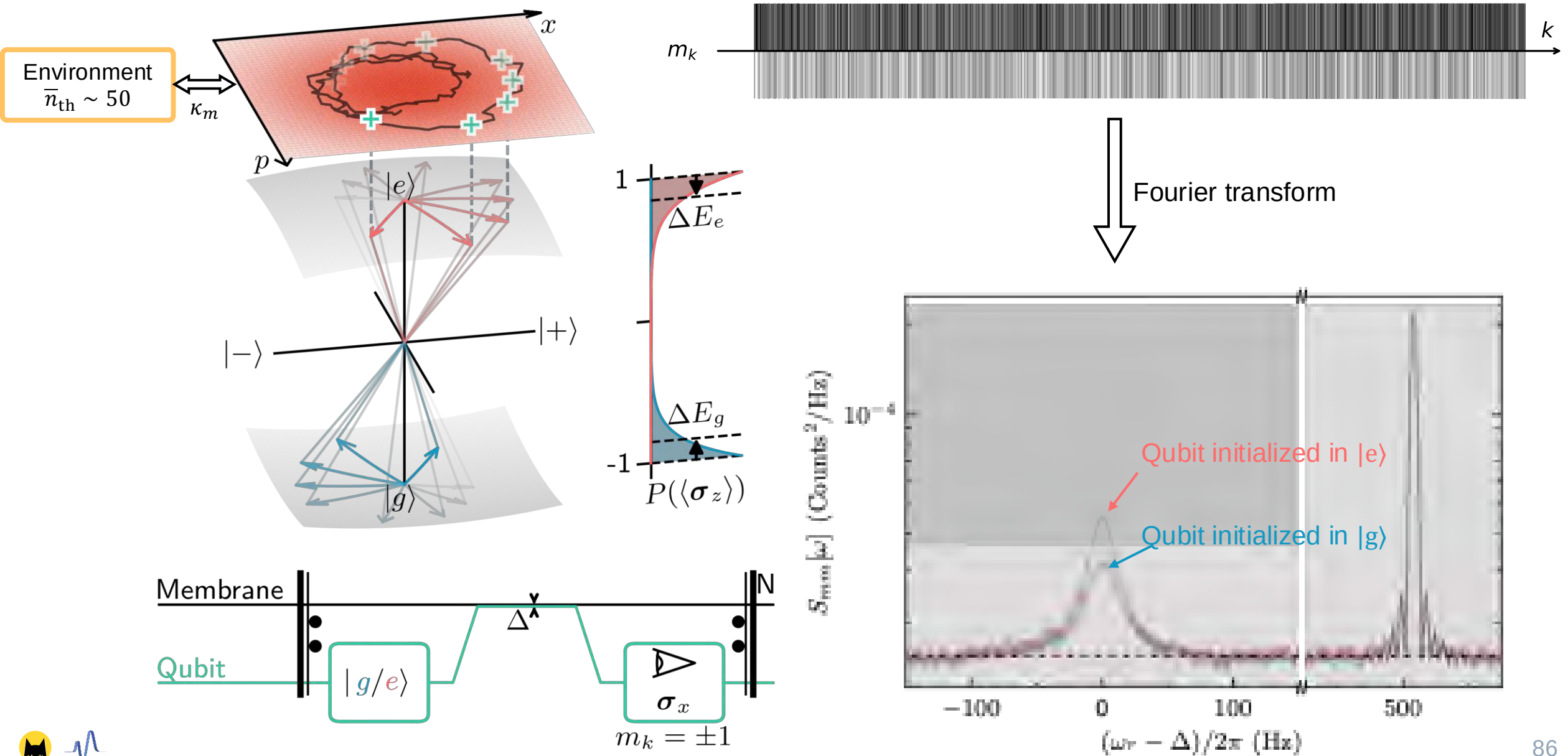


# Spectrum analyzer measurement

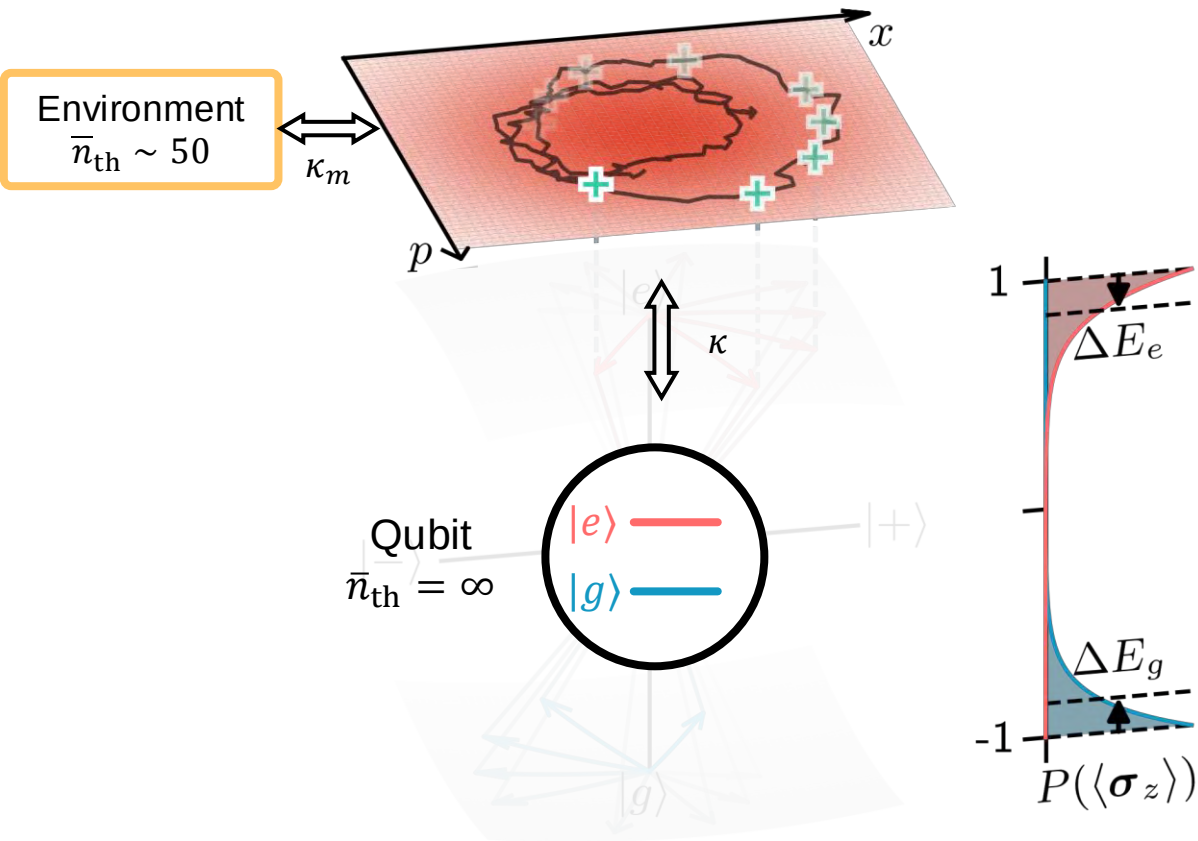




# Spectrum analyzer measurement



# Spectrum analyzer measurement



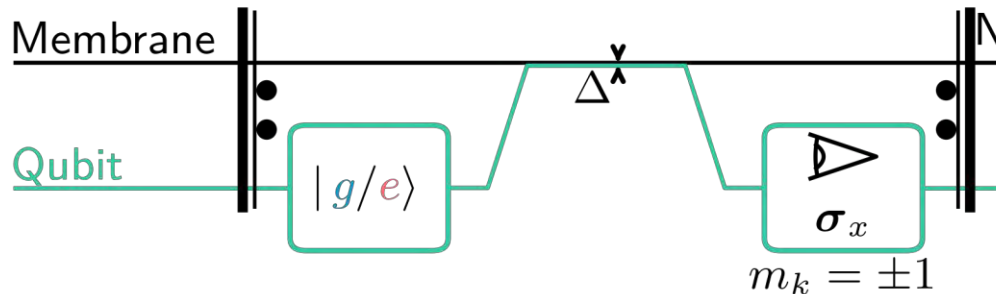
Mean energy exchange per cycle :

$$\Delta E_{g/e} = \pm n_{\text{th}} (\Omega T_{\text{int}})^2$$

Qubit induced **damping/amplification** at rate :

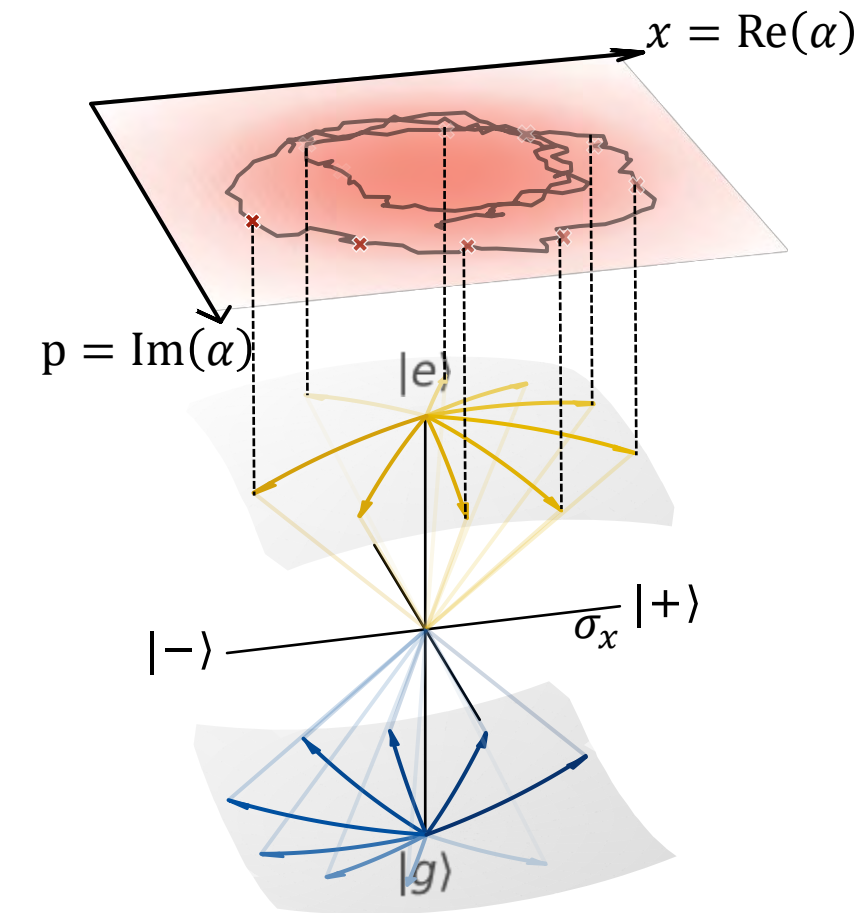
$$\kappa = \pm (\Omega T_{\text{int}})^2 / T_{\text{cycle}}$$

Similar to dynamical back action in optomechanics



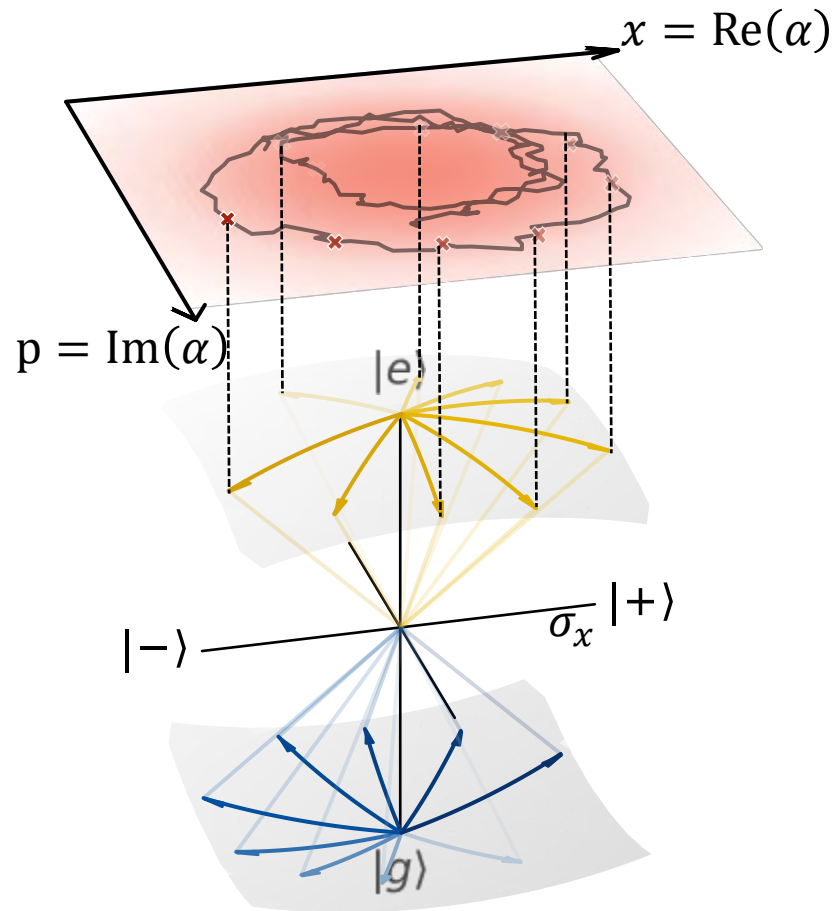
## Semi-classical model

$$H/\hbar = i \frac{\Omega}{2} (\alpha^* \sigma_- - \alpha \sigma_+)$$



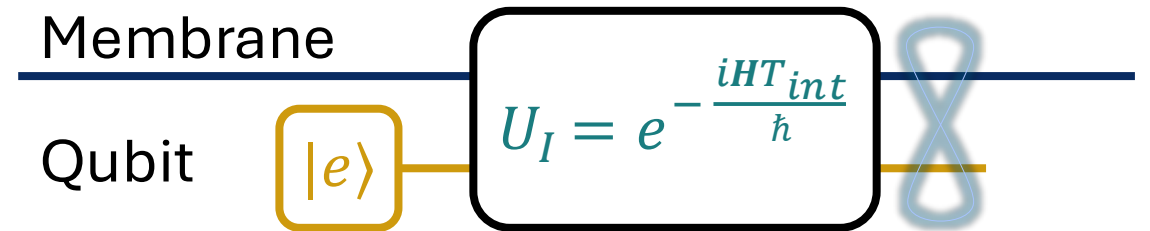
## Semi-classical model

$$H/\hbar = i\frac{\Omega}{2}(\alpha^* \sigma_- - \alpha \sigma_+)$$



## Quantum model

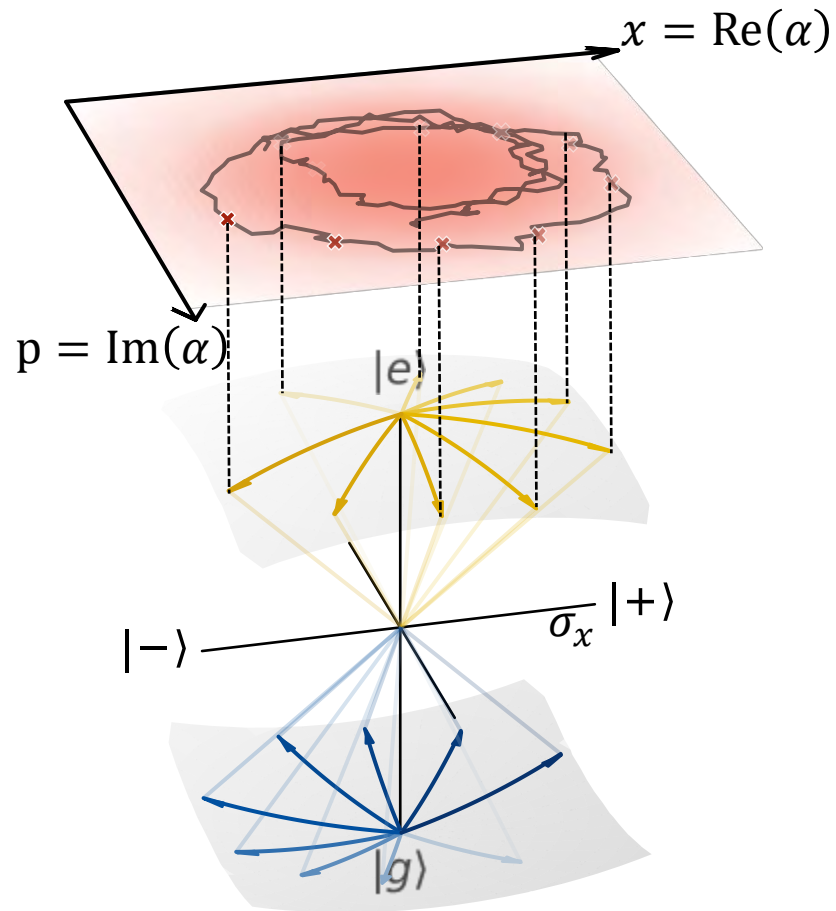
$$H/\hbar = i\frac{\Omega}{2}(a^\dagger \sigma_- - a \sigma_+)$$



- Some (small amount) of entanglement is generated by the interaction

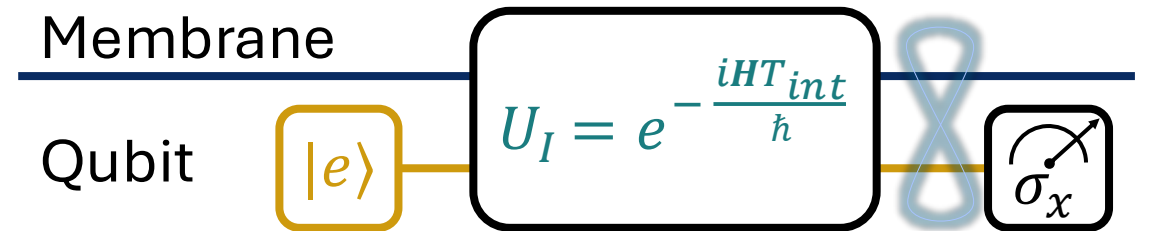
## Semi-classical model

$$H/\hbar = i \frac{\Omega}{2} (\alpha^* \sigma_- - \alpha \sigma_+)$$



## Quantum model

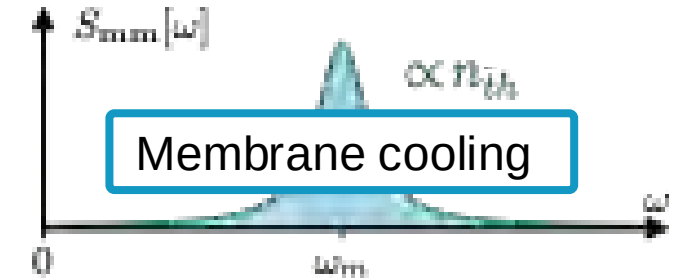
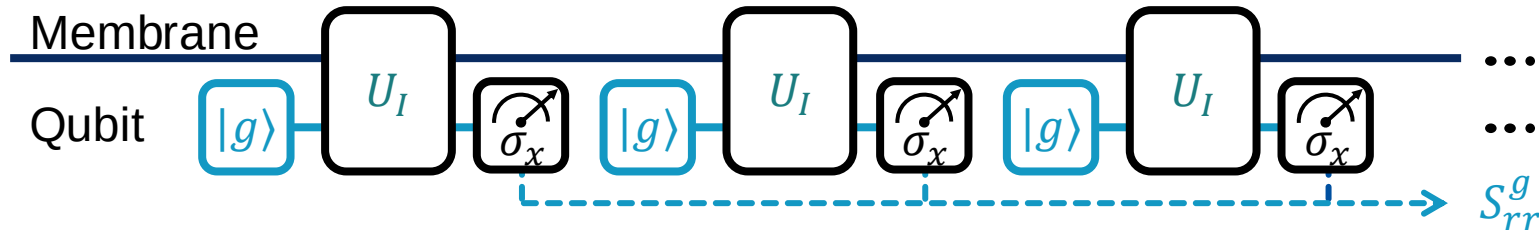
$$H/\hbar = i \frac{\Omega}{2} (a^\dagger \sigma_- - a \sigma_+)$$



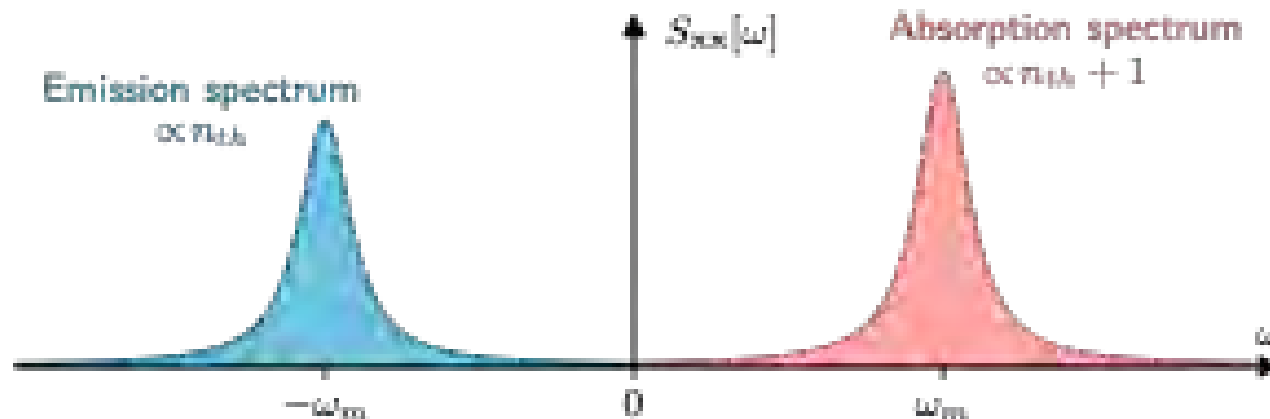
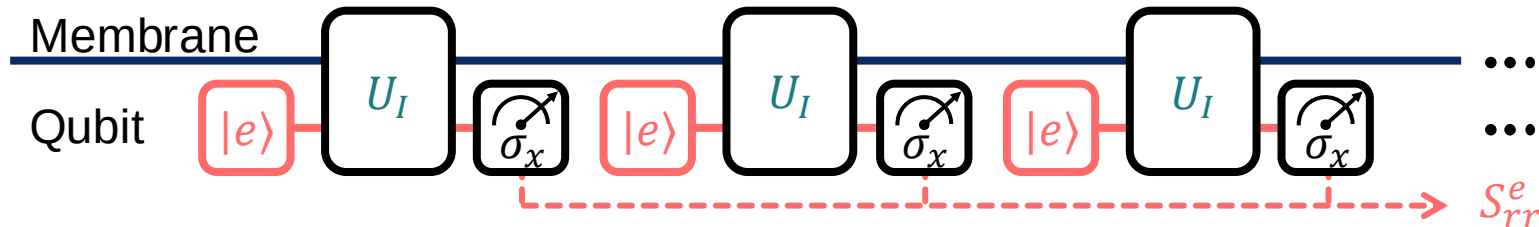
- Some (small amount) of entanglement is generated by the interaction
  - The qubit projective measurement steers the membrane state in a conditional state that depends on the measurement outcome  $|+\rangle_x$  or  $|-\rangle_x$
- Quantum backaction

# Mitigating dynamical backaction

Preparation in  $|g\rangle$ :



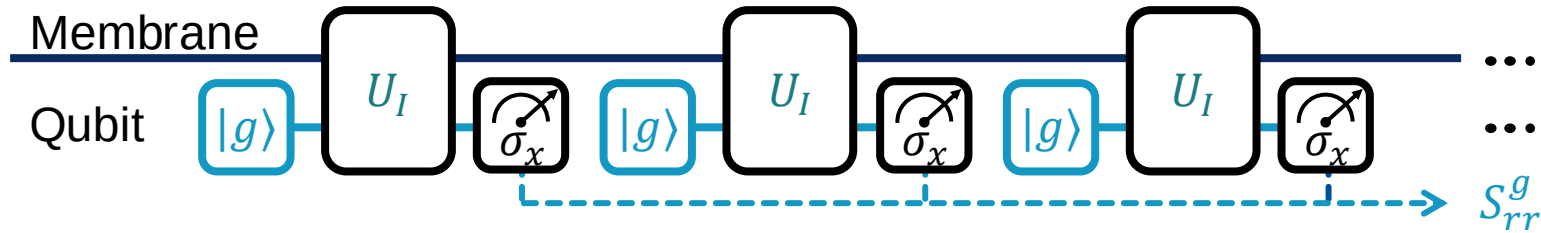
Preparation in  $|e\rangle$ :



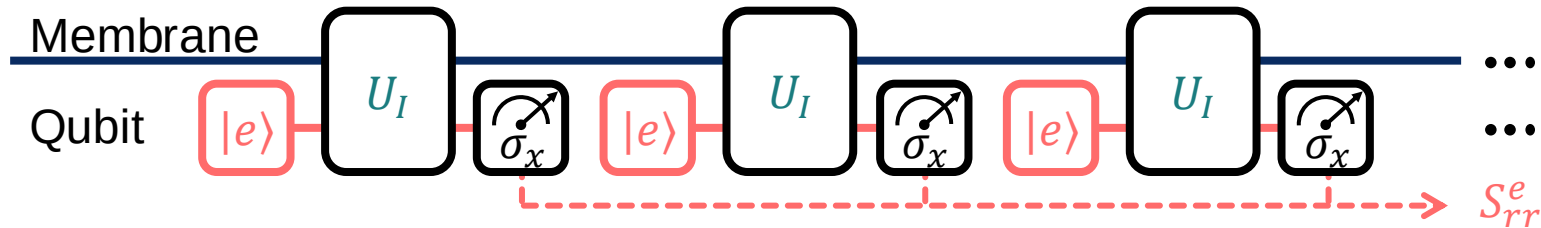


# Mitigating dynamical backaction

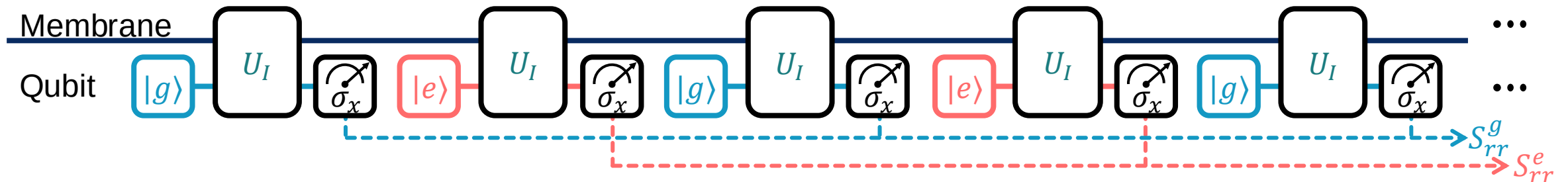
Preparation in  $|g\rangle$ :



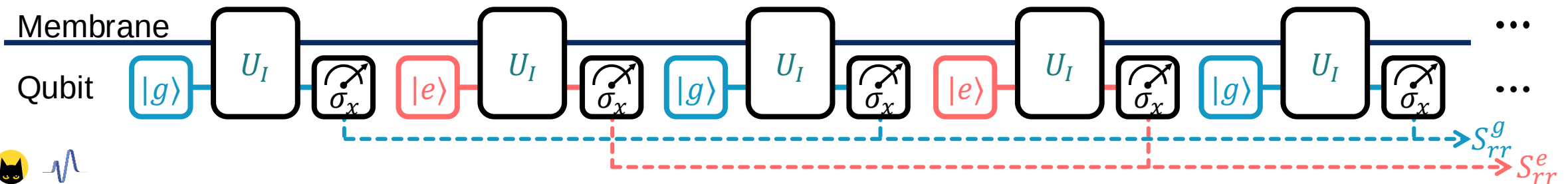
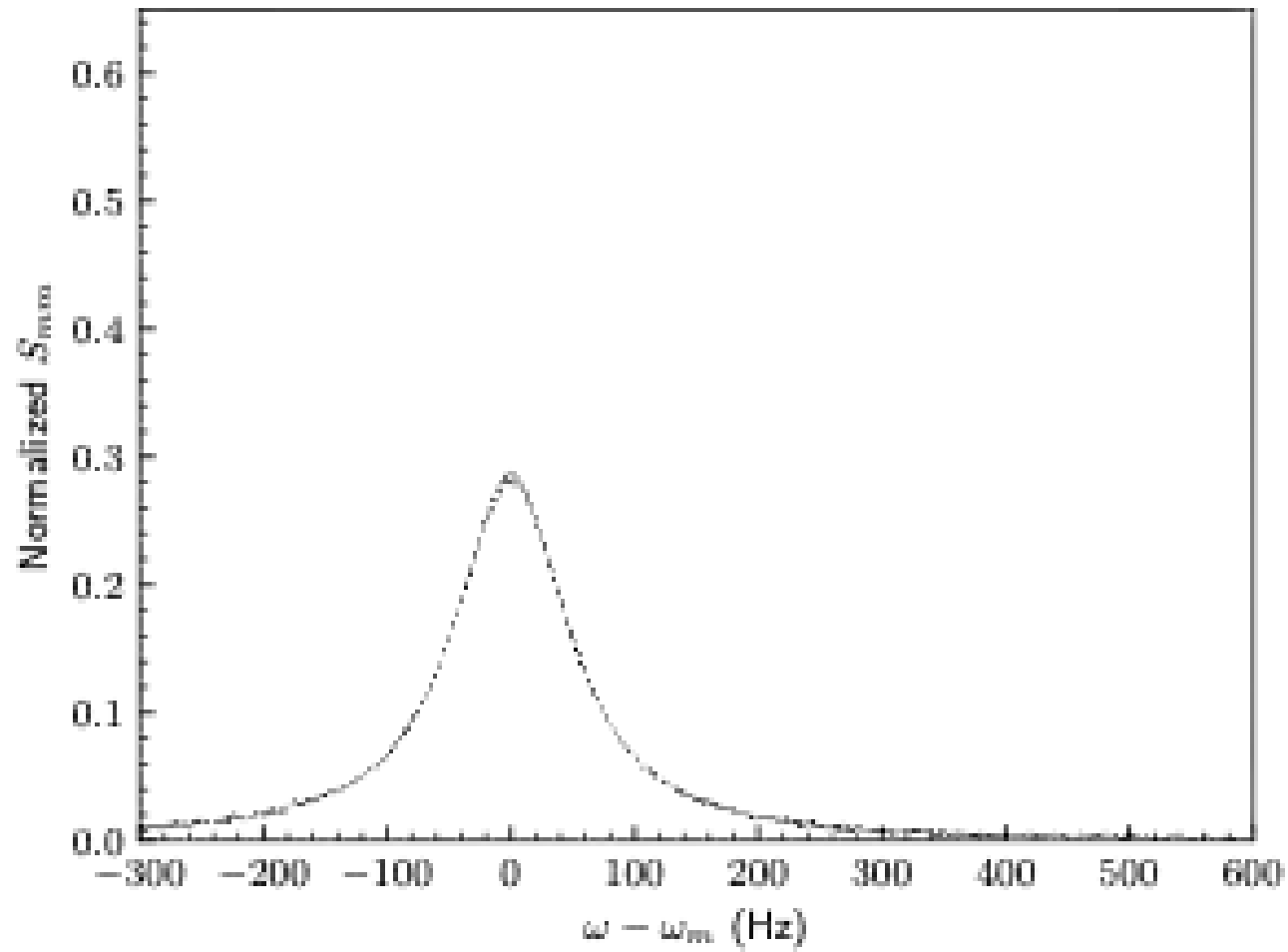
Preparation in  $|e\rangle$ :



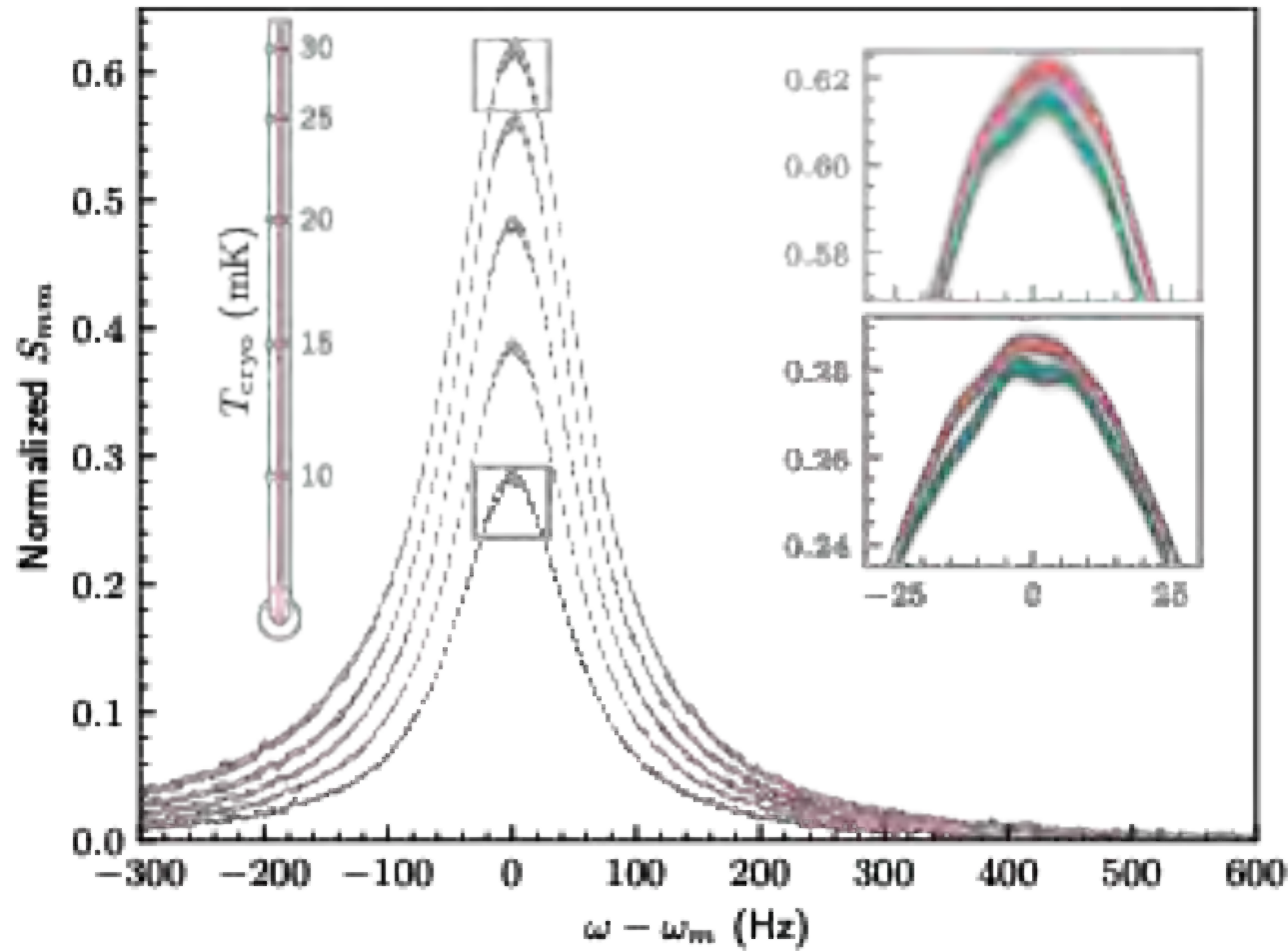
Alternating qubit state preparation:



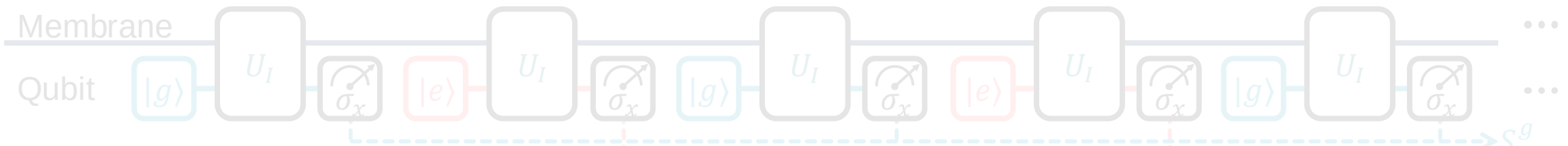
# Spectrum asymmetry



# Spectrum asymmetry

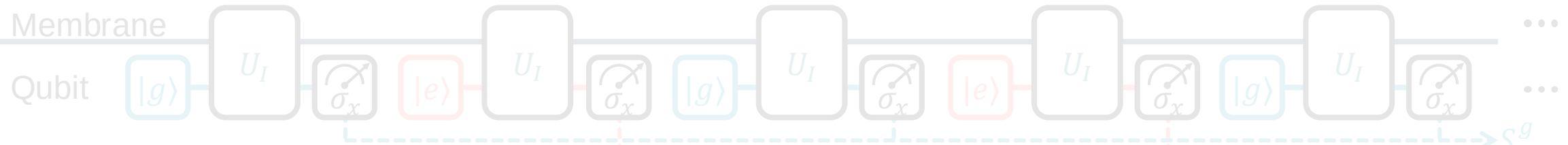
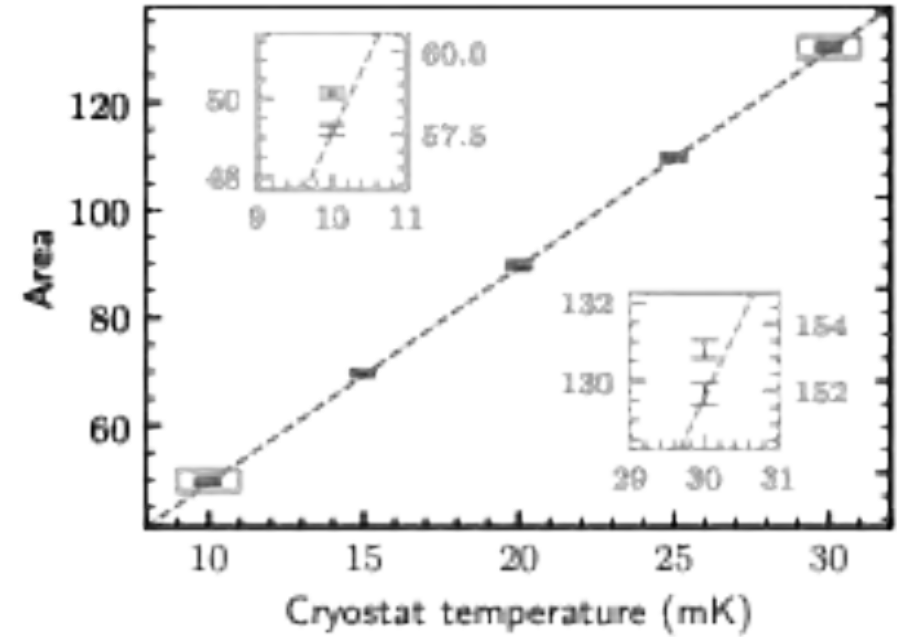
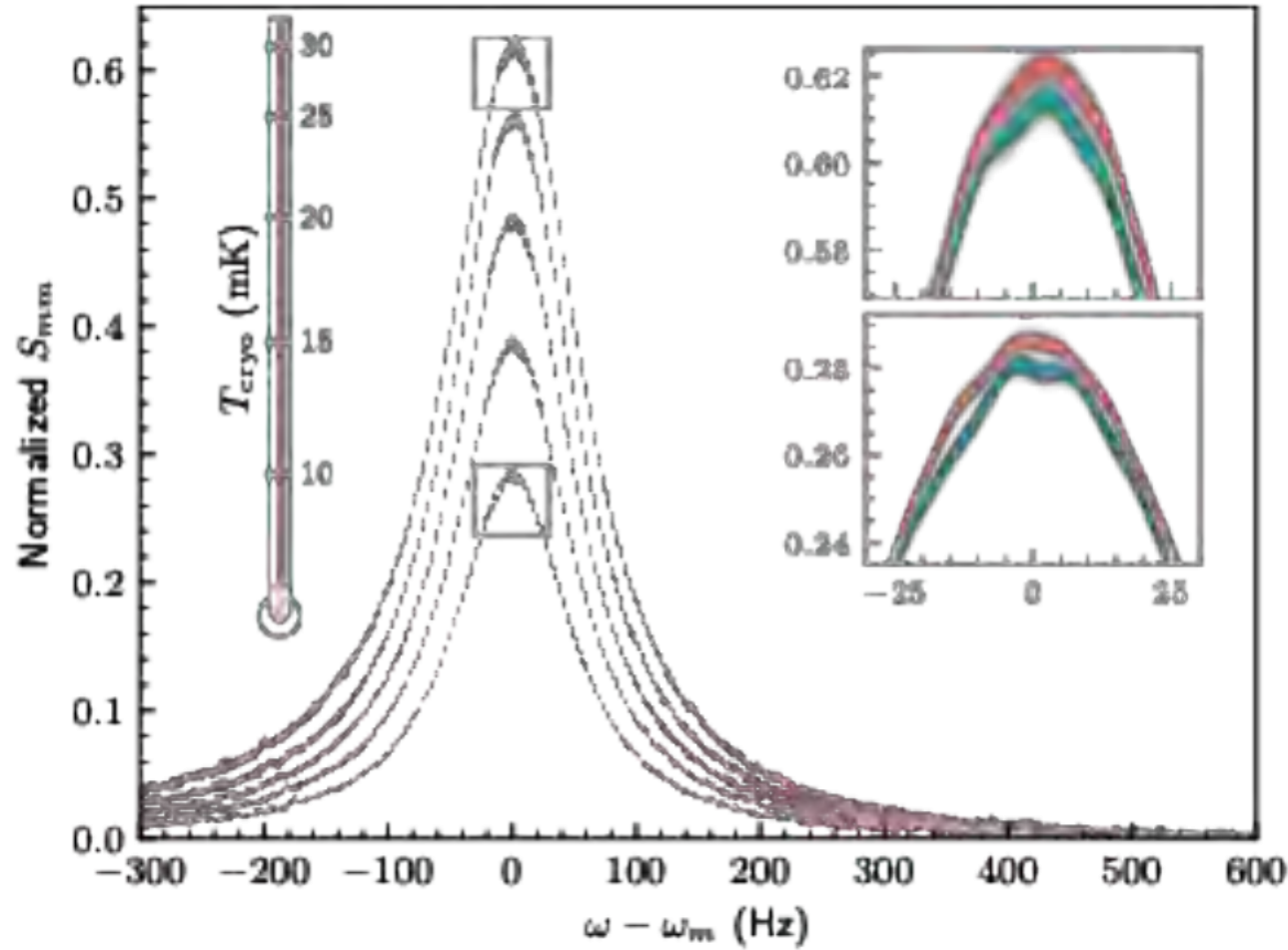


50 phonons at 10 mK  
 $\Rightarrow$  2 % difference



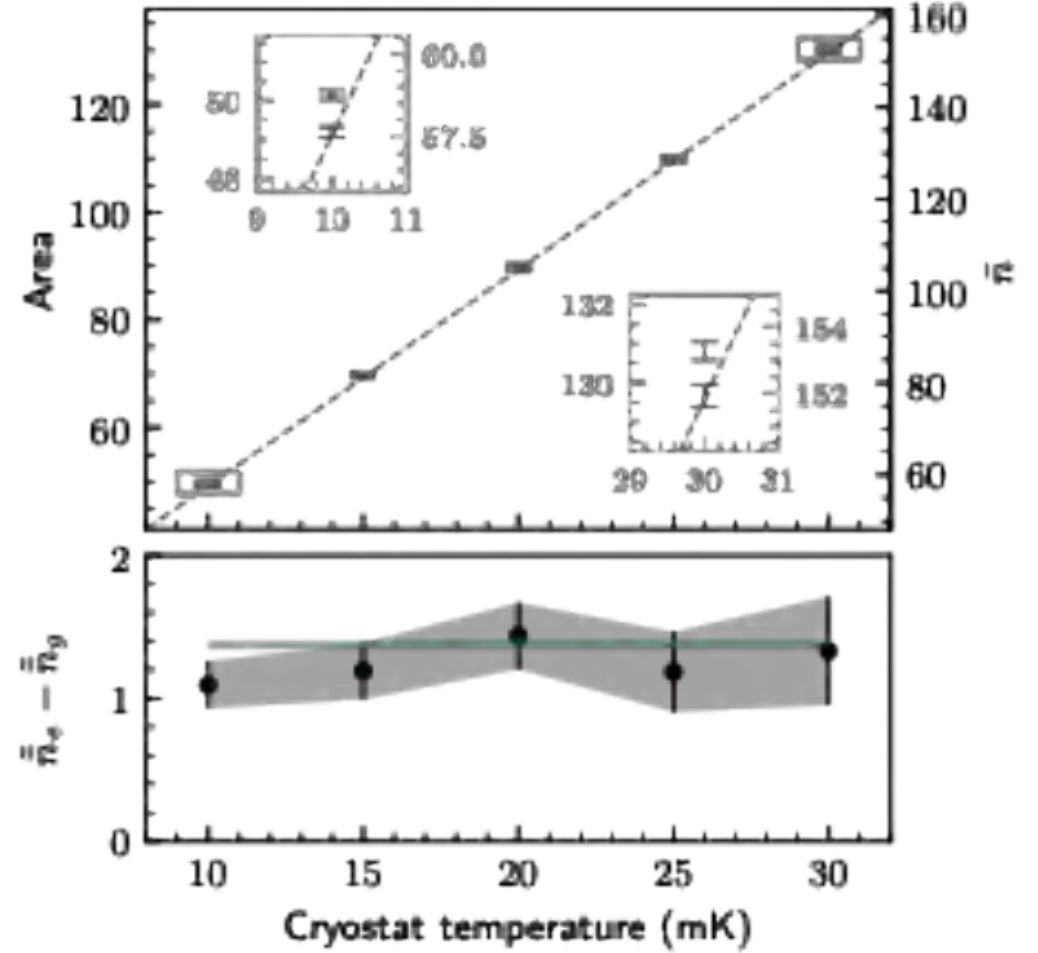
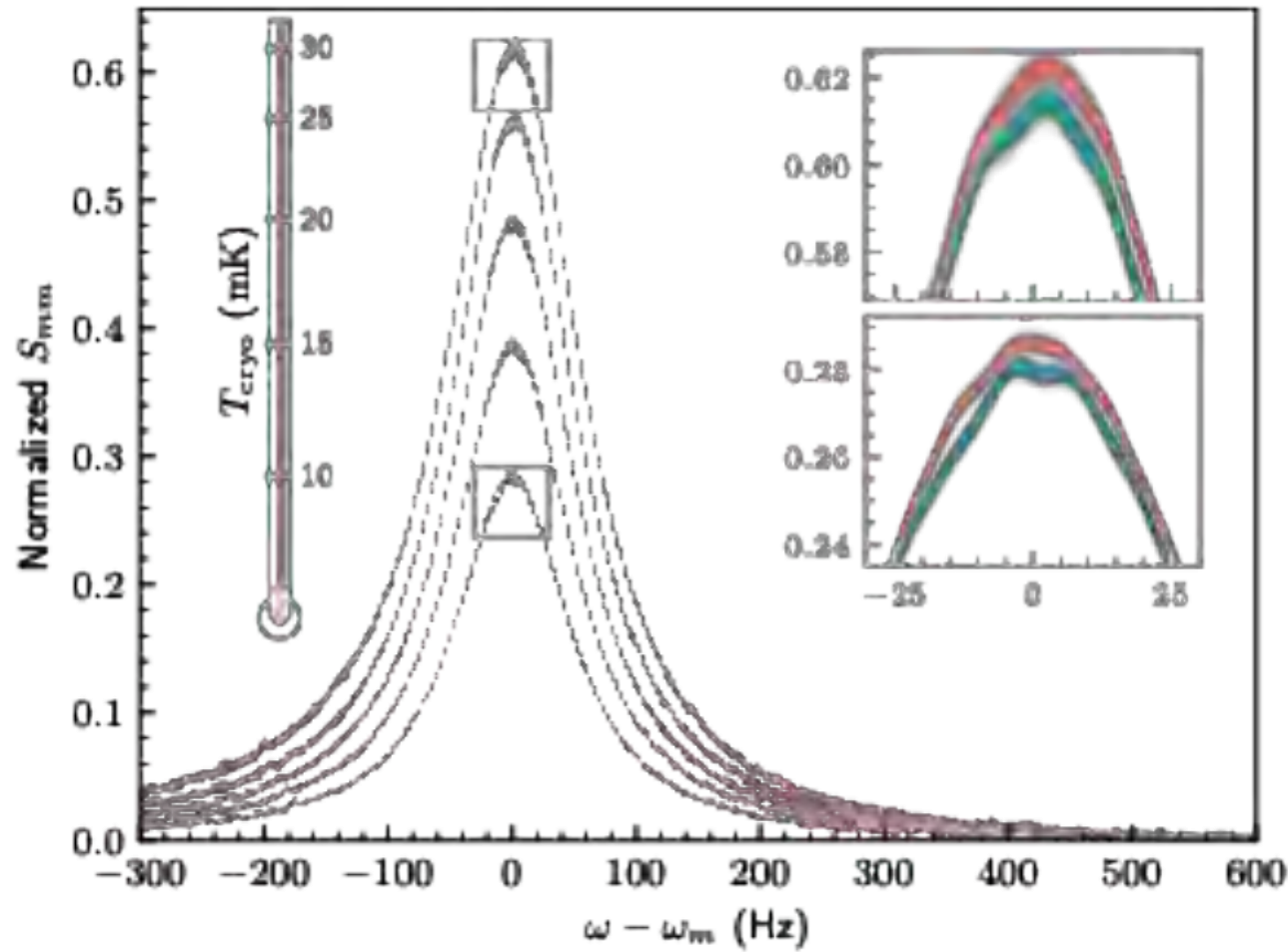
K. Gerashchenko\*, R. Rousseau\*, L. Balembois\*, *et al.*, **Probing the quantum motion of a macroscopic mechanical oscillator with a radio-frequency superconducting qubit**, arXiv:2505.21481 (2025)

# Spectrum asymmetry



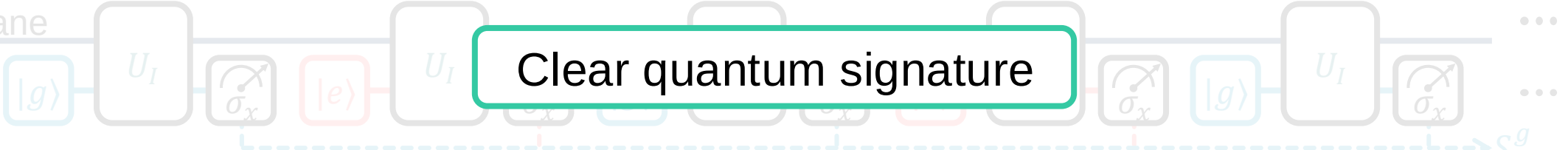
K. Gerashchenko\*, R. Rousseau\*, L. Balembois\*, *et al.*, **Probing the quantum motion of a macroscopic mechanical oscillator with a radio-frequency superconducting qubit**, arXiv:2505.21481 (2025)

# Spectrum asymmetry



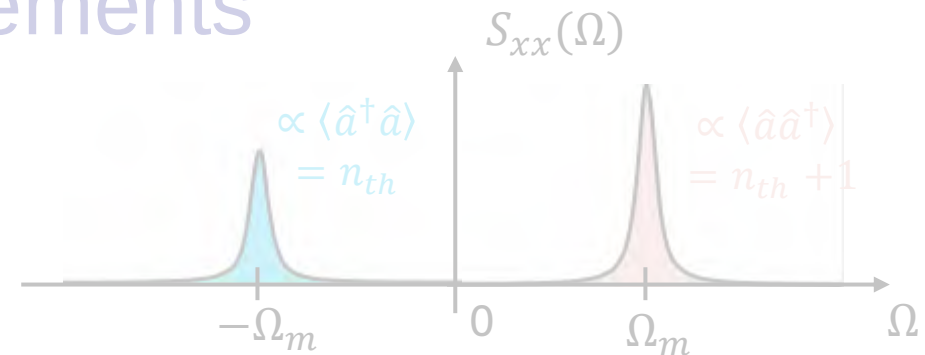
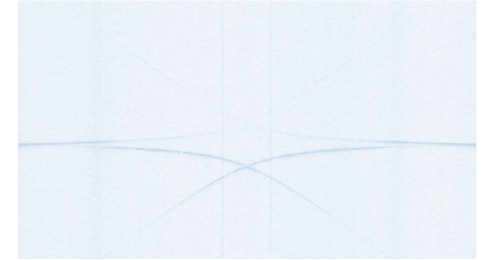
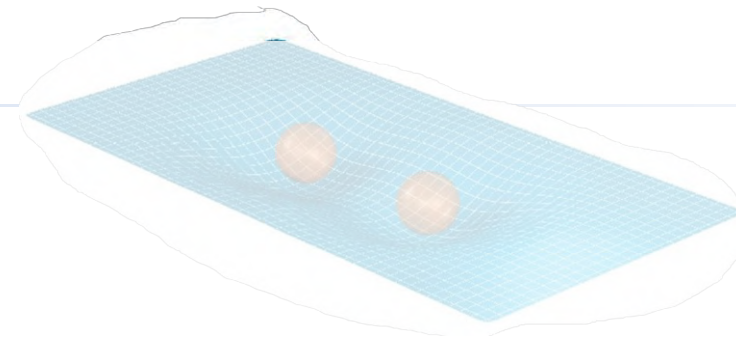
Membrane

Qubit



K. Gerashchenko\*, R. Rousseau\*, L. Balembois\*, *et al.*, Probing the quantum motion of a macroscopic mechanical oscillator with a radio-frequency superconducting qubit, arXiv:2505.21481 (2025)

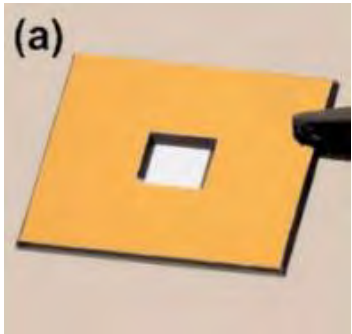
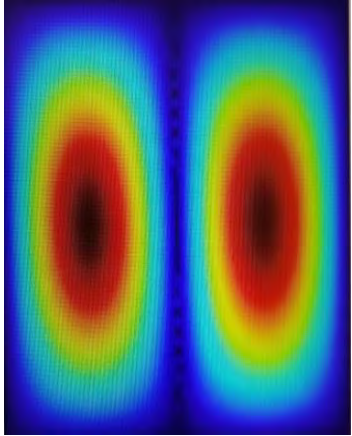
- Introduction/Motivation
- Radio-frequency fluxonium qubits
- Hybrid qubit/mechanical system: MecaFlux
- Quantum asymmetry measurements
- Perspectives





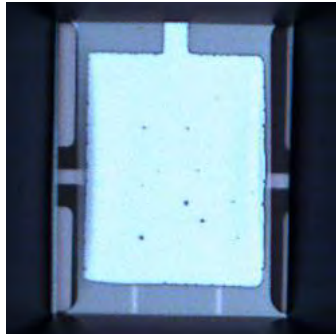
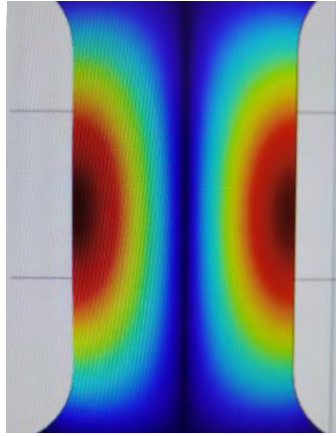
# From Rectangular to Torsional Membrane

**BEFORE**



Rectangular  
Membrane  
5 ms @ 4 MHz

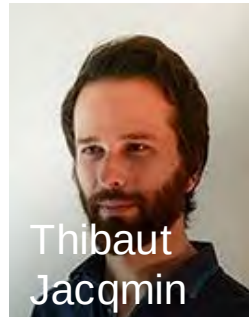
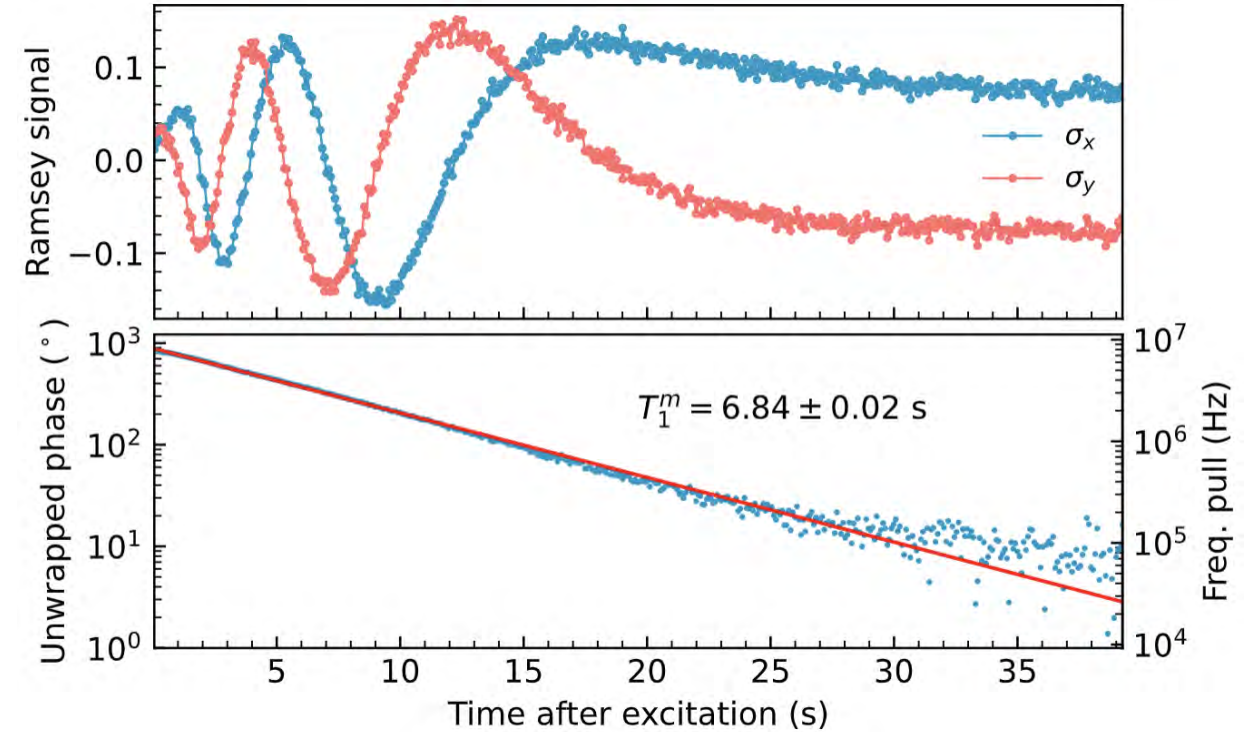
**AFTER**



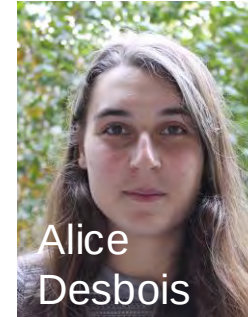
Torsional  
Membrane  
6.84 s @ 2 MHz

**1,000× improvment!**

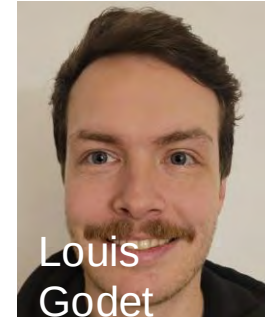
**Ringdown measurement**



Thibaut  
Jacqmin



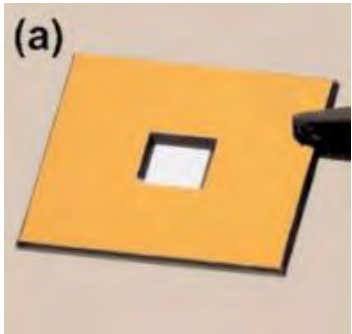
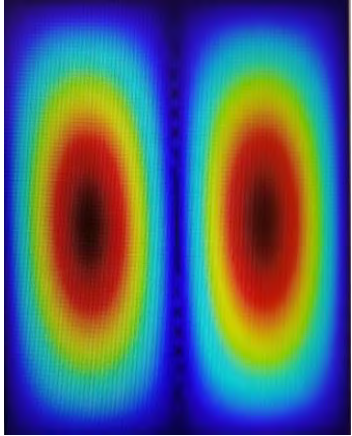
Alice  
Desbois



Louis  
Godet

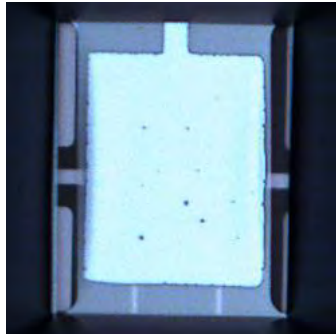
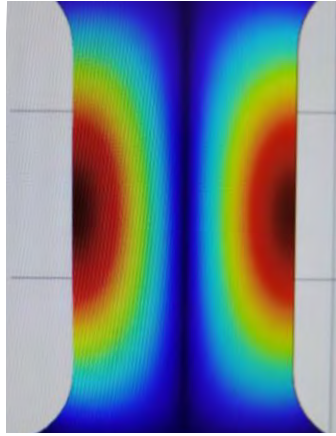
# From Rectangular to Torsional Membrane

**BEFORE**



Rectangular  
Membrane  
5 ms @ 4 MHz

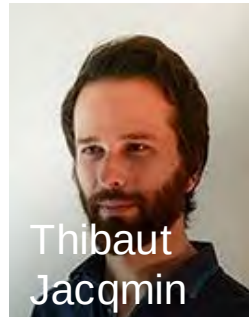
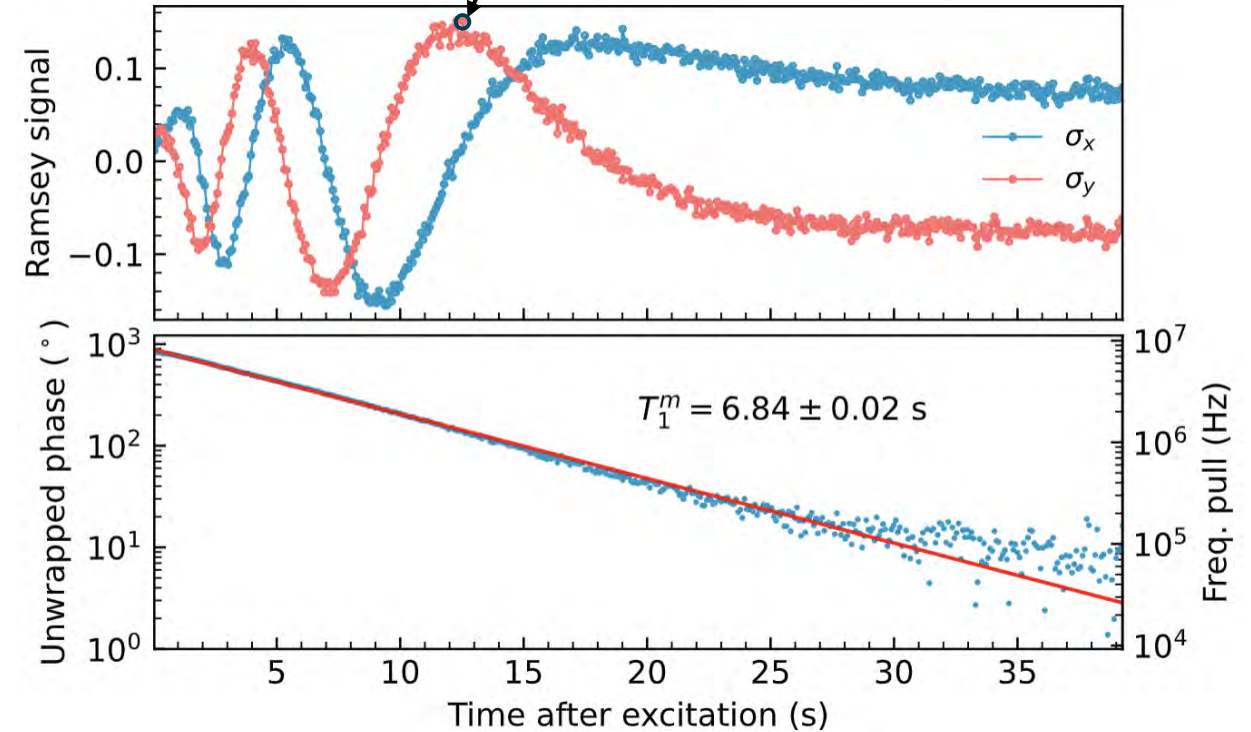
**AFTER**



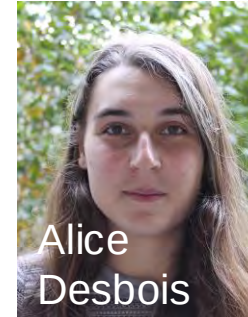
Torsional  
Membrane  
6.84 s @ 2 MHz

**1,000× improvment!**

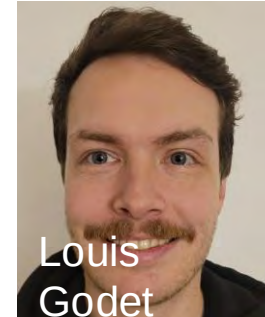
Ringdown measurement



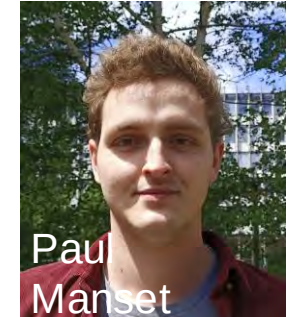
Thibaut  
Jacqmin



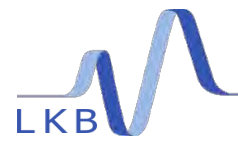
Alice  
Desbois



Louis  
Godet

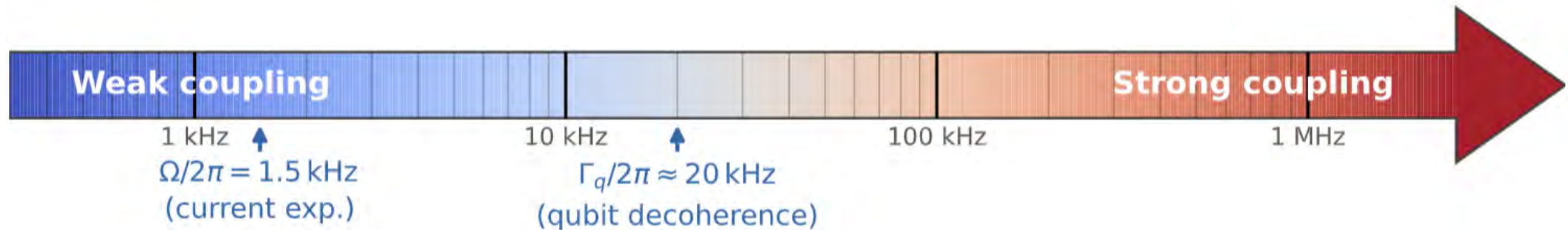


Paul  
Manset



# Perspectives: path to strong coupling

$$\text{Coupling: } \Omega \propto V_b/d^2$$

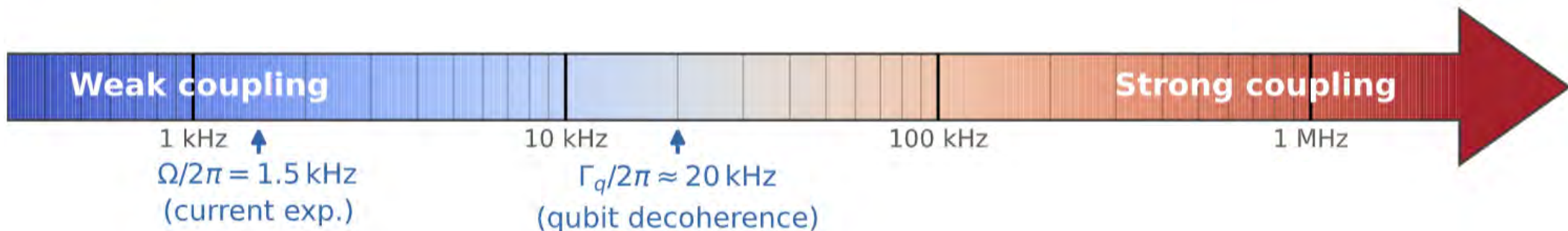
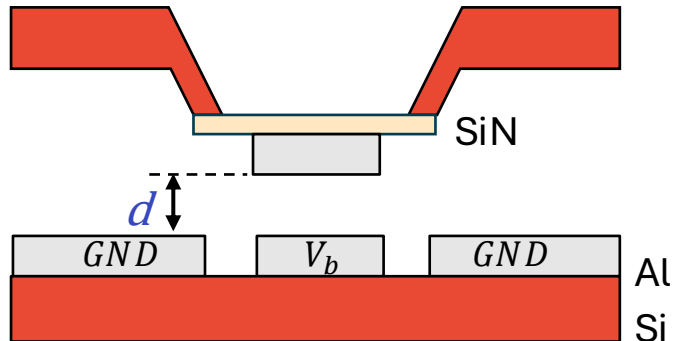




$$\text{Coupling: } \Omega \propto V_b/d^2$$

→ Reduce gap (WIP)

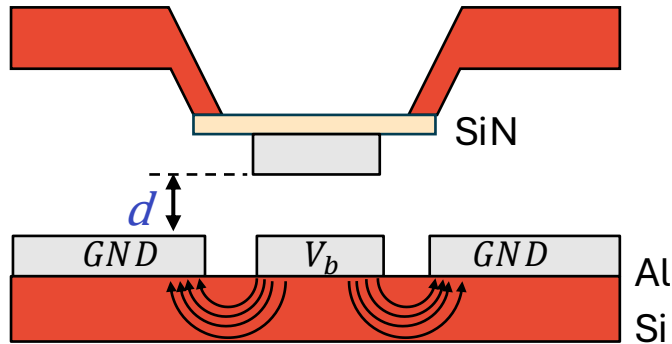
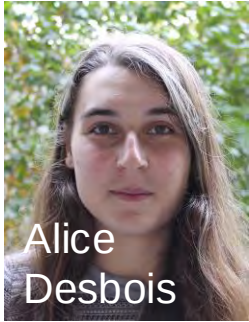
Objective:  $d = 2.5 \mu\text{m} \rightarrow 500 \text{ nm}$  ( $\Omega \times 25$ )



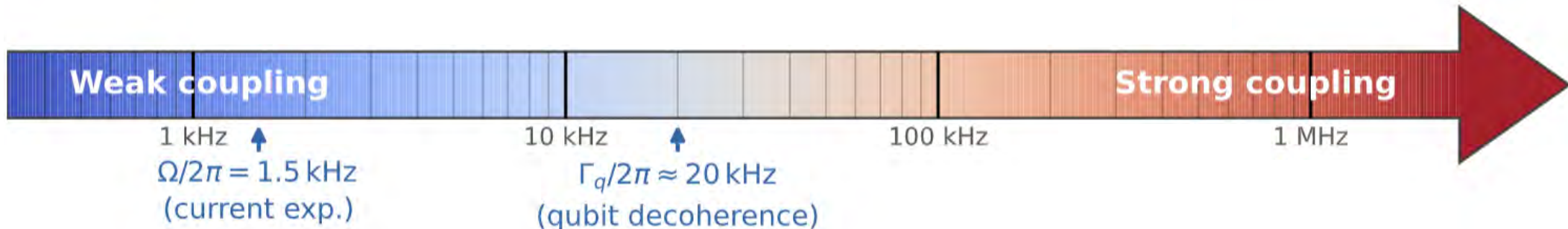
$$\text{Coupling: } \Omega \propto V_b/d^2$$

→ Reduce gap (WIP)

Objective:  $d = 2.5 \mu\text{m} \rightarrow 500 \text{ nm}$  ( $\Omega \times 25$ )



Avalanche ionization in Si  
 $|V_b| \lesssim 10 \text{ V}$



# Perspectives: path to strong coupling

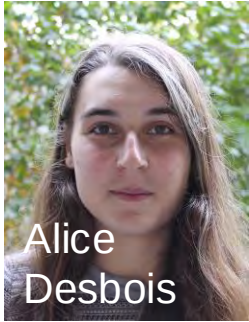
$$\text{Coupling: } \Omega \propto V_b/d^2$$

→ Reduce gap (WIP)

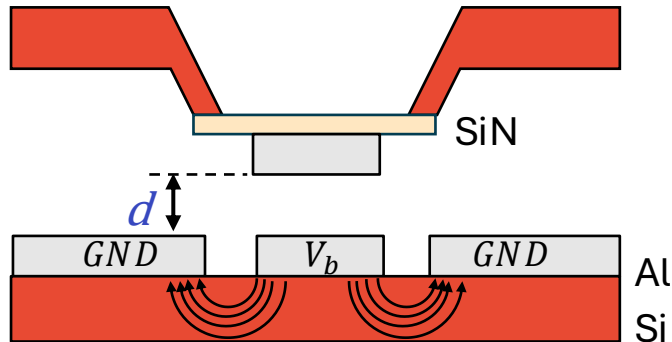
Objective:  $d = 2.5 \mu\text{m} \rightarrow 500 \text{ nm}$  ( $\Omega \times 25$ )

→ Increase bias voltage

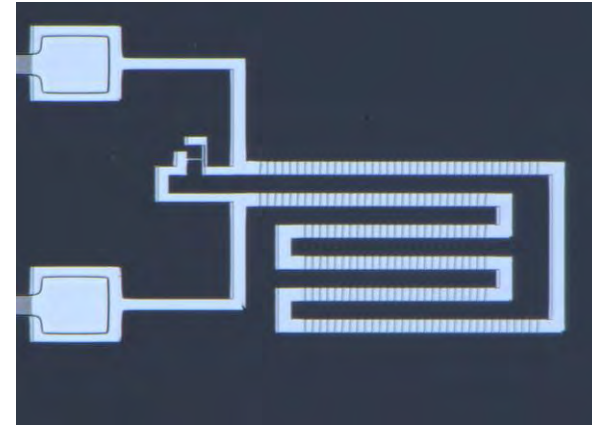
Demonstrated:  $V_b = 9 \text{ V} \rightarrow 150 \text{ V}$  ( $\Omega \times 13$ )



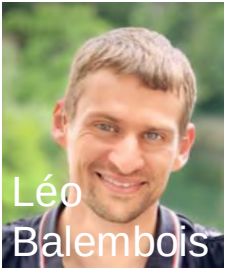
Alice Desbois



Avalanche ionization in Si



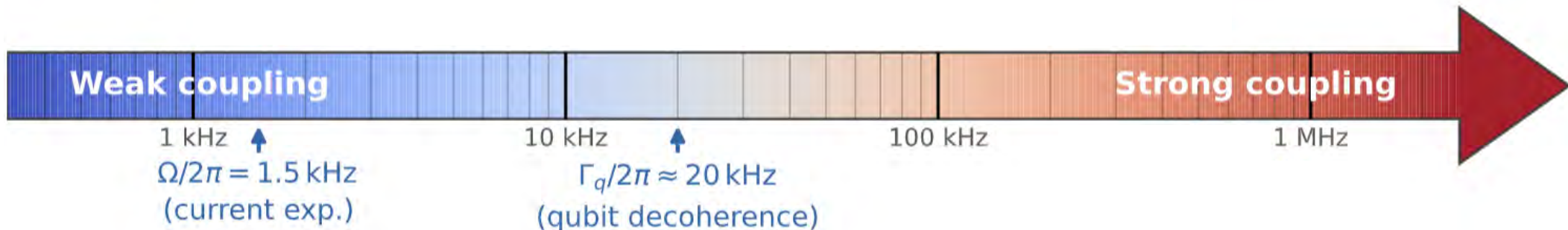
Fluxonium on Sapphire



Léo Balembois



Paul Manset





# Perspectives: path to strong coupling

$$\text{Coupling: } \Omega \propto V_b/d^2$$

→ Reduce gap (WIP)

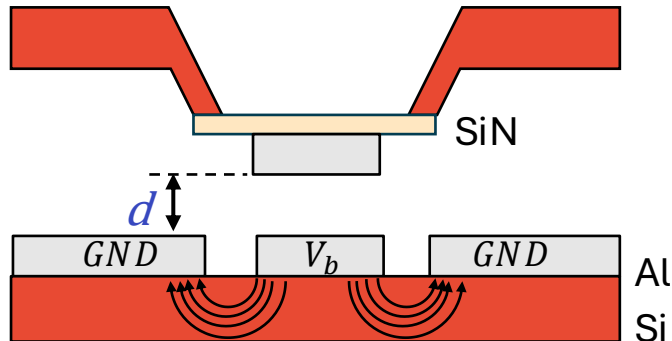
Objective:  $d = 2.5 \mu\text{m} \rightarrow 500 \text{ nm}$  ( $\Omega \times 25$ )

→ Increase bias voltage

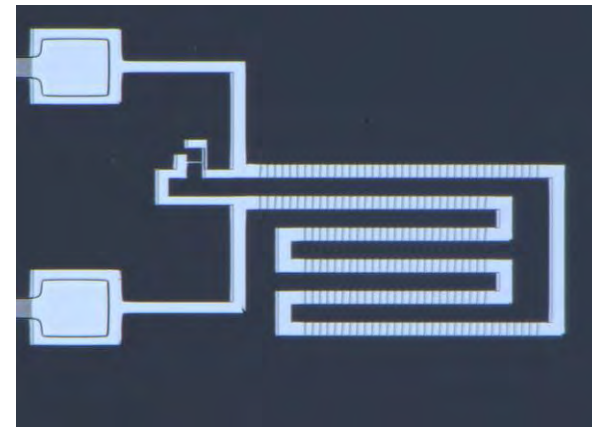
Demonstrated:  $V_b = 9 \text{ V} \rightarrow 150 \text{ V}$  ( $\Omega \times 13$ )



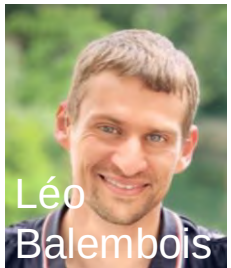
Alice Desbois



Avalanche ionization in Si



Fluxonium on Sapphire

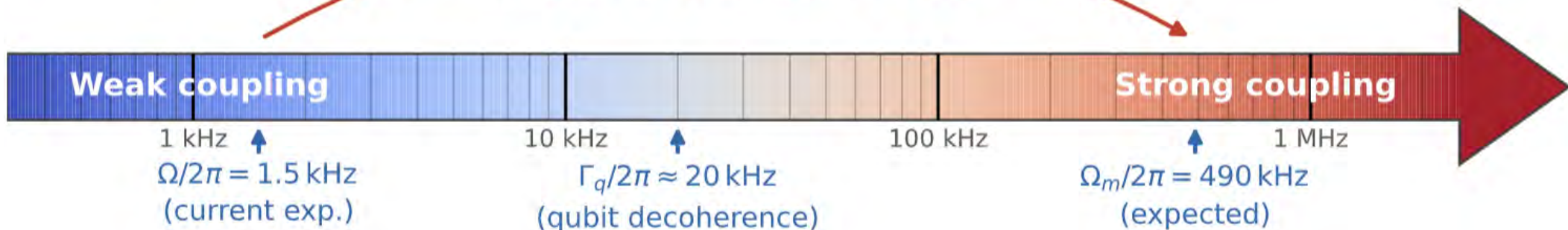


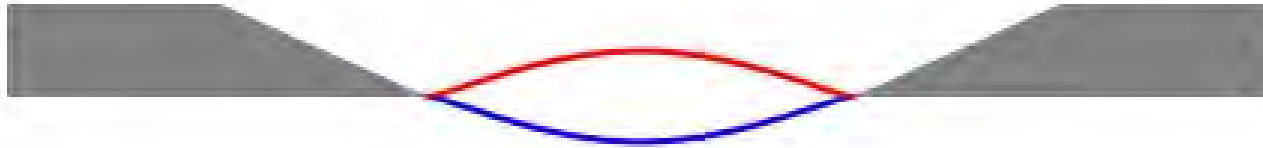
Léo Balembois



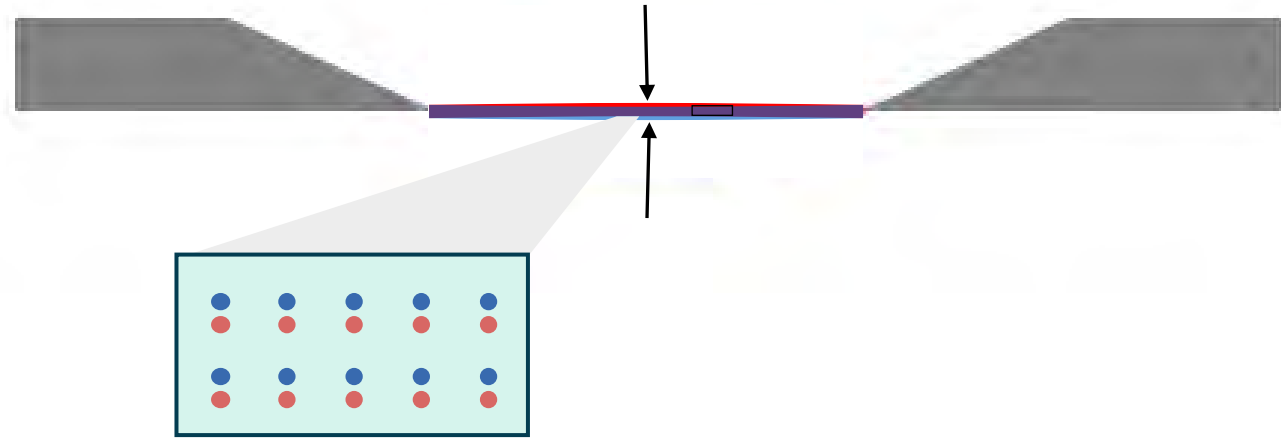
Paul Manset

$\times 25 \text{ (gap)} \times 13 \text{ (voltage)} = \times 325$





$$|\Psi\rangle = \frac{1}{\mathcal{N}} (|\alpha\rangle + |-\alpha\rangle)$$



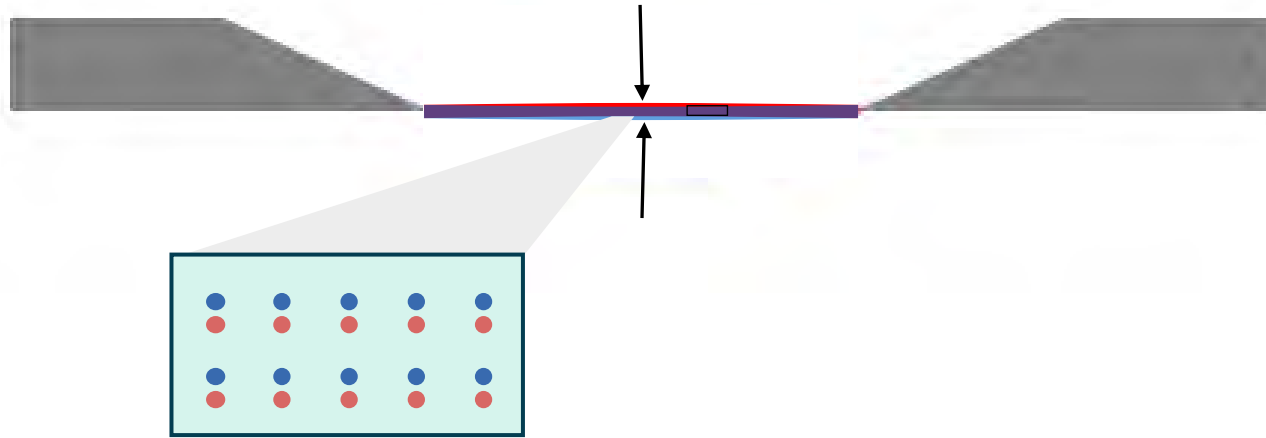
$$|\Psi\rangle = \frac{1}{\mathcal{N}} (|\alpha\rangle + |-\alpha\rangle)$$

Necessary conditions:

- spatial separation of the cat components

$$x_{zpf} |\alpha| \geq \begin{array}{c} \text{Atomic} \\ \text{nucleus} \\ \text{diameter} \end{array} \simeq 5 \text{ fm}$$

$$x_{zpf} = \sqrt{\frac{\hbar}{2m\Omega}}$$



$$|\Psi\rangle = \frac{1}{\mathcal{N}} (|\alpha\rangle + |-\alpha\rangle)$$

Necessary conditions:

- spatial separation of the cat components

$$x_{zpf} |\alpha| \geq \begin{array}{c} \text{Atomic} \\ \text{nucleus} \\ \text{diameter} \end{array} \simeq 5 \text{ fm}$$

$$x_{zpf} = \sqrt{\frac{\hbar}{2m\Omega}}$$

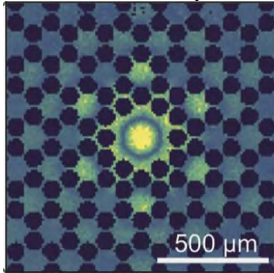
- Gravitational decoherence time

$$\tau_g \ll \begin{array}{c} \text{Decoherence time} \\ \text{due to mechanical} \\ \text{losses} \end{array}$$

$$\tau_g \approx 3 \times 10^{-15} \left( \frac{1\text{kg}}{m} \right) \text{ s}$$

# Mechanical systems in the Quantum Regime

Tsaturyan...  
Schliesser (2017)



1 MHz

2 ng

2 fm

Viennot...  
Lehnert (2018)

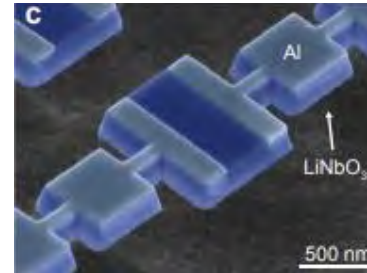


25 MHz

50 pg

4 fm

Arrangoiz-Arriola...  
Safavi-Naeini (2019)

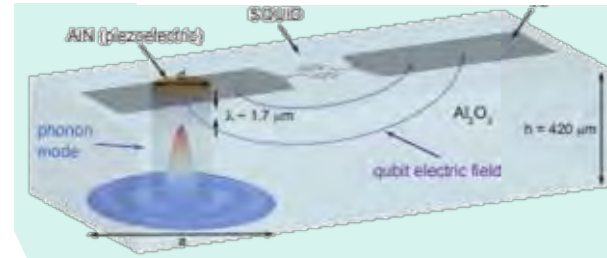


2.4 GHz

2 pg

1.3 fm

Chu...  
Schoelkopf (2017)



6.6 GHz

50 μg

150 zm

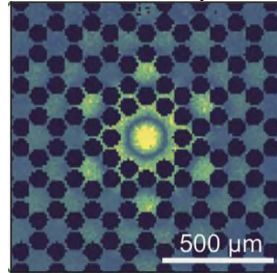
Resonator  
frequency  $\Omega_m/2\pi$

Motional  
mass  $m$

Zero-point  
fluctuations  $x_{zpf}$

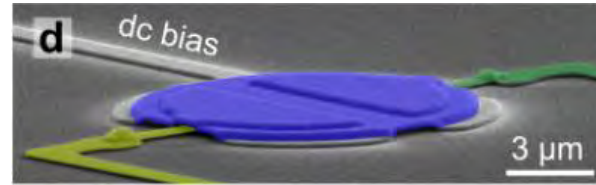
# Mechanical systems in the Quantum Regime

Tsaturyan...  
Schliesser (2017)



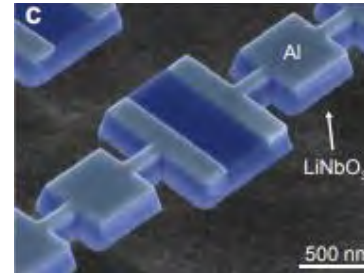
1 MHz

Viennot...  
Lehnert (2018)



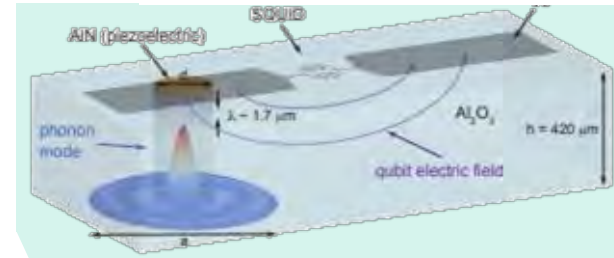
25 MHz

Arrangoiz-Arriola...  
Safavi-Naeini (2019)



2.4 GHz

Chu...  
Schoelkopf (2017)



6.6 GHz

Resonator  
frequency  $\Omega_m/2\pi$

Motional  
mass  $m$

Zero-point  
fluctuations  $x_{zpf}$

Q-factor

Decoherence time of  
a 5-fm cat  $\tau_d$

Gravitational de-  
coherence time  $\tau_g$

2 ng

50 pg

2 pg

50 μg

2 fm

4 fm

1.3 fm

150 zm

$10^9$

$10^5$

$10^4$

$10^6$

100 ms

$\geq$

1 ms

100 μs

$\leq$

60 ms

100 ns

$\ll$

1.5 s

150 fs

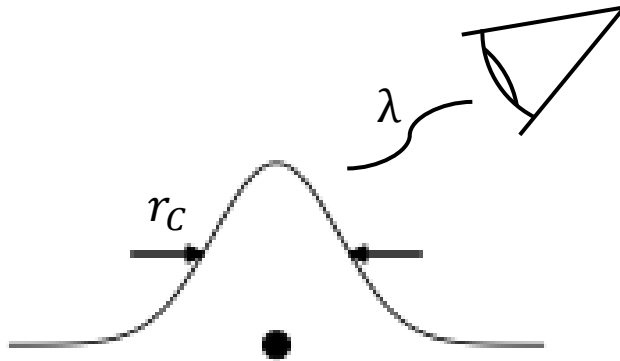
$\ll$

60 ns



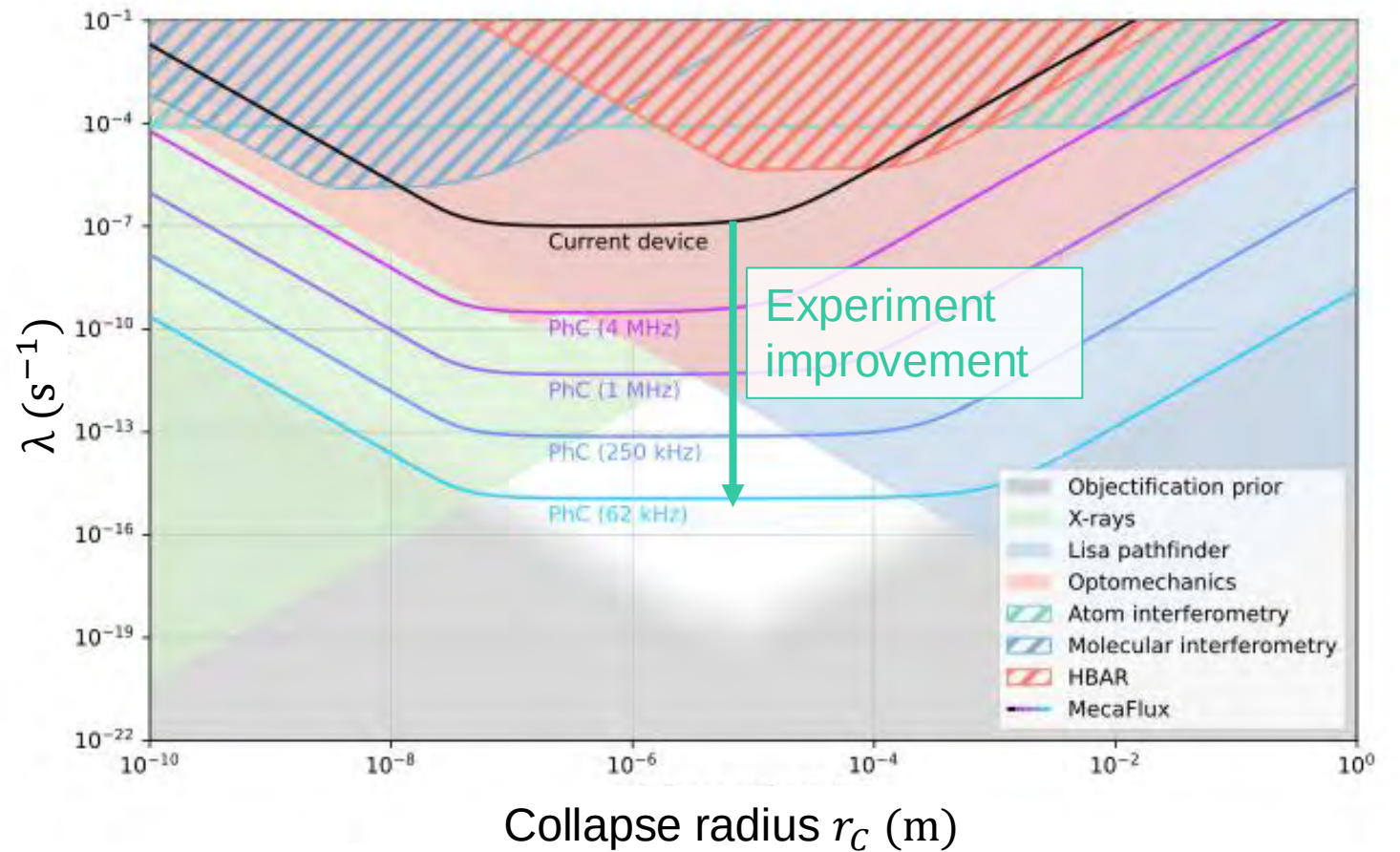
**Collapse models** resolve the **quantum measurement paradox** by introducing small modifications to the Schrodinger's equation

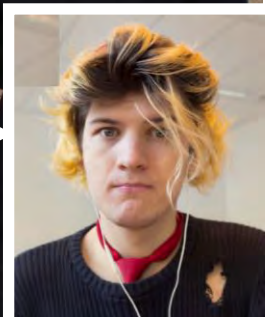
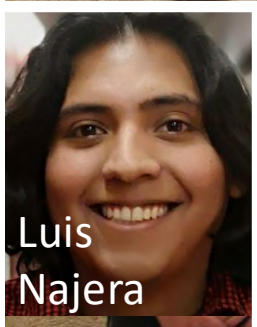
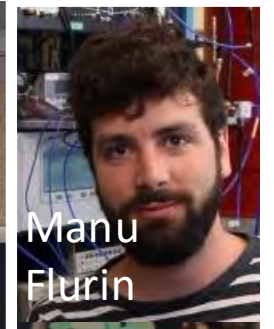
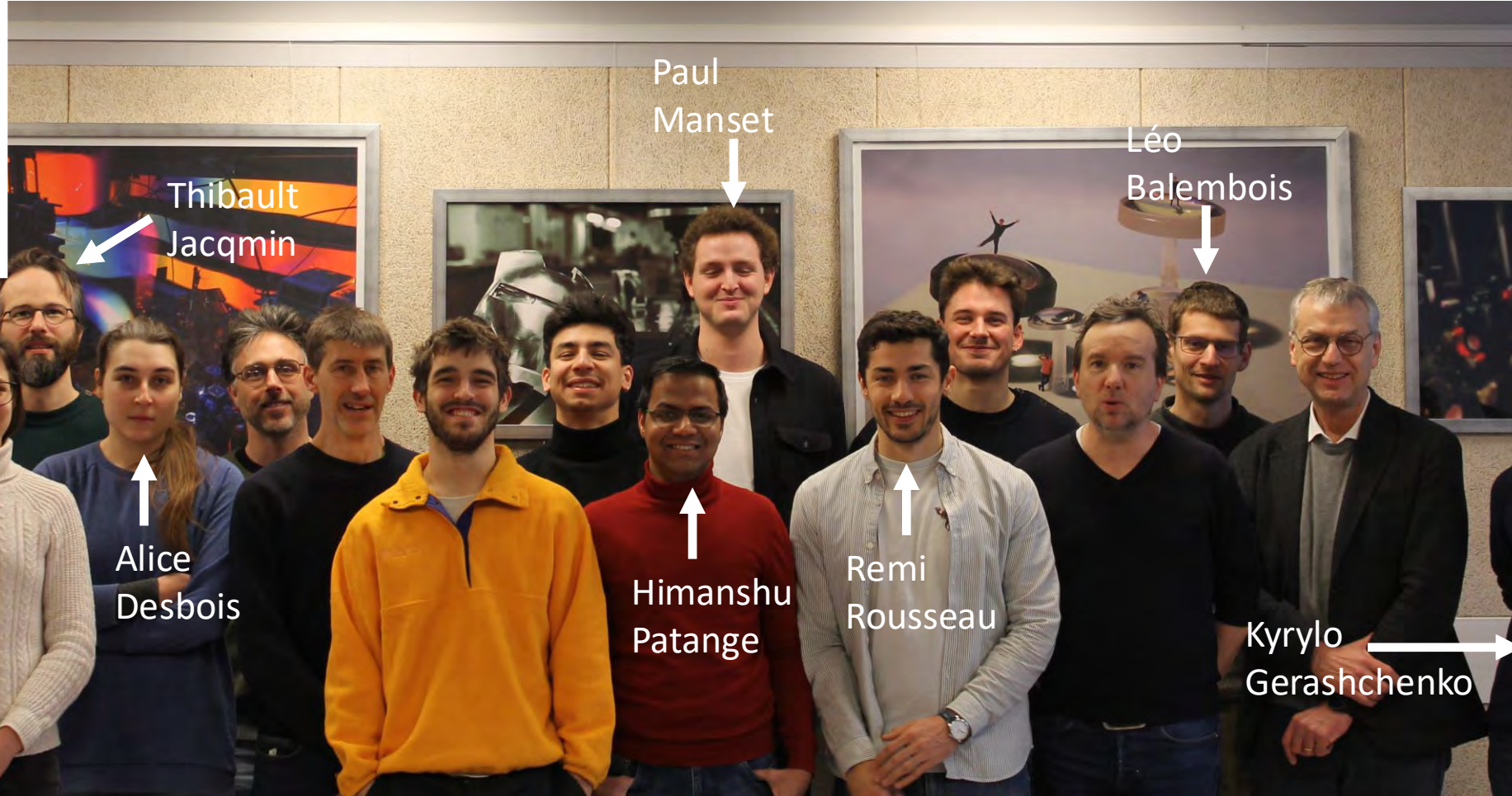
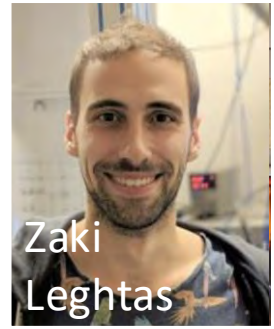
Continuous spontaneous localization model



$r_C$  : measurement radius

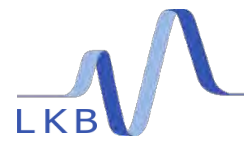
$\lambda$  : measurement rate





Thanks !!





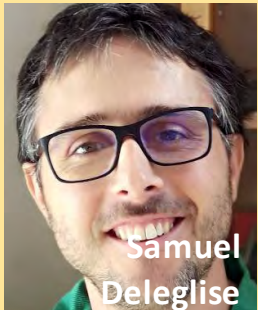
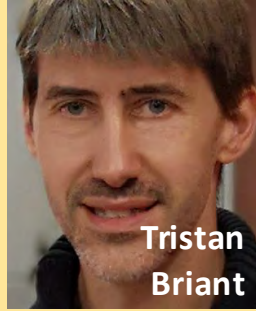
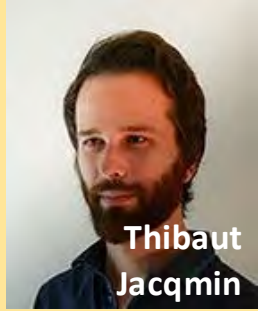
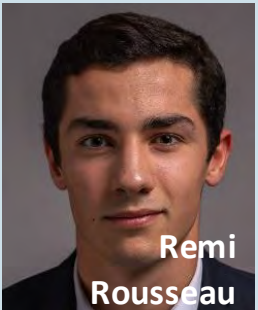
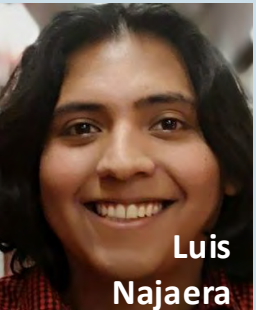
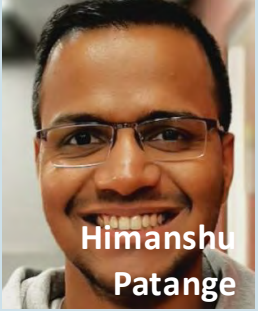
# The team @ LKB

PhD students

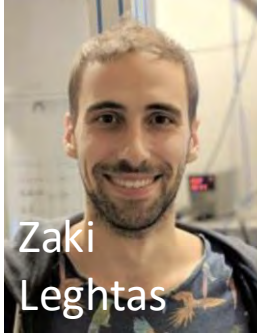
Post doc

Permanent

This project



Collab.:



Cifre  
LKB/Alice &  
Bob

Optomechanical  
thermometry

Quantum limits